

MULTIGROUP ALBEDO THEORY WITH APPLICATION TO
NEUTRONIC CALCULATIONS FOR A GAS CORE REACTOR

By

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Dedicated to My Parents
to Whom I Owe Everything.

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For nuclear reactors with effective reflectors, the balance of neutrons shows great sensitivity to the inventory of those neutrons which spend time in the reflector. It is essential, in developing a theory which is to explain neutronic phenomenon, such as the core nonleakage probabilities and the variation of the effective neutron multiplication factor, that the details of the transfer of neutrons across the core-reflector interface be included.

The main objective of this research is to develop a multigroup neutron albedo theory to neutronic calculations of thermal nuclear reactors. The neutron albedo theory method is an analytical and intuitive method that does not solve an explicit equation such as the Boltzmann transport equation. Rather, transport analysis relations are used to obtain information about the behavior of neutrons during segments of their lifetime, which are relatively easy to

visualize. The purpose is to obtain a variety of descriptive nonleakage probabilities and the effective and infinite neutron multiplication factors. Thus, the neutron albedo theory methodology is very different from the conventional methodology used in the XSDRNPM computer code.

The main parameters of the four-group albedo theory approach are investigated with the aid of calculational examples. The three reactor system models chosen are representative of one-dimensional, spherical geometry, and thermal reactor systems which exhibit transport effects to a significant degree. The computations are also performed on the systems with the XSDRNPM computer code to serve as a reference.

The most significant result is the ability of the four-group albedo method to compute the core nonleakage probabilities and the effective neutron multiplication factor with a reasonably fast and low-priced method that is clearly more understandable in physical terms than conventional neutron transport methods. This can be accomplished without calculating the core neutron fluxes of the reactor systems investigated. The most significant application value of the method is in the ability to calculate and understand reactivity coefficients.

The agreement of the theoretical calculations confirmed the validity of the multigroup neutron albedo theory for the nonleakage probability and effective neutron multiplication factor calculations for reactor system of the types treated.

CHAPTER I INTRODUCTION

Background: Two-Group Albedo Theory

In 1958, a two-group albedo theory [1,2] was developed by Alan M. Jacobs for the calculation of reactivity temperature coefficients of nuclear reactors characterized by the Pennsylvania State University research reactor (PSR). The PSR is of the swimming pool variety with highly enriched fuel elements and natural water as moderator-coolant and reflector.

The agreement of the experimental measurements and theoretical calculations [1] confirmed the validity of this two-group albedo theory for temperature coefficient calculations for reactor systems of the type treated.

Motivation for the Work: Generalization to Multigroup Theory

For nuclear reactors with effective reflectors the balance of neutrons shows great sensitivity to the inventory of those neutrons which spend time in the reflector. It is essential, in developing a theory which is to explain neutronic phenomenon, such as the core nonleakage probabilities and the variation of the effective neutron multiplication factor, that the details of the transfer of neutron across the core-reflector interface be included.

The generalization of the two-group albedo theory to a multigroup albedo theory is required if the method of analysis is to be useful in systems that require a more detailed energy spectrum treatment.

Particular Application to a Gas Core Reactor

Gas core reactors have undergone theoretical and experimental studies at the University of Florida. A significant number of neutronic analyses have been conducted. The results of the researchers [3] indicated that gas core reactor concepts have some distinct advantages for power generation compared with the other nuclear power. It also has been demonstrated that typical neutronic calculations using multidimensional codes exhibit high computational cost. Because of the extremely important role of the reflector-moderator in gas core concepts, it is thought to be useful to apply the multigroup albedo approach to such systems. The gas core examples studied at the University of Florida provide excellent test cases for the value of this approach.

Dissertation Objectives

The main objective of this research is to develop a multigroup neutron albedo theory to neutronic calculations of thermal nuclear reactors. One major application is in the determination of the core nonleakage probabilities and

the variation of the effective neutron multiplication factor with a reasonably quick and inexpensive method.

Thus, the primary objective of this research can be summarized as follows:

- (1) Develop a multigroup neutron albedo theory, and
- (2) Compare the results of the calculations for the model problems based on this multigroup neutron albedo theory with the theoretical results obtained when performed with the XSDRNPM computer code [4].

Dissertation Organization

In Chapter II, the systematic derivation of the multigroup neutron albedo theory is presented.

The evaluation techniques for the core and reflector albedos are discussed in Chapter III.

In Chapter IV, the evaluation techniques for partial principal nonleakage probabilities, P_{o_1} , are derived and described.

In Chapters V, VI, and VII, the model problems are introduced and results obtained from static one-dimensional neutronic calculations.

The conclusions obtained from this research and some suggestions for future work are included in Chapter VIII.

The appendices provide detailed support for the material discussed in this work as well as detailed specification of some parameter calculations which are too cumbersome to include in the text body.

CHAPTER II

DERIVATION OF THE MULTIGROUP NEUTRON ALBEDO THEORY

Formulation of the Effective Neutron Multiplication Factor

All real reactors are to some degree nonuniform. The nonuniformity may consist of a reflector, i.e., a region of low absorption material around the core which serves to reflect neutrons back into the multiplying core-region, or it may consist of a breeding blanket which catches the escaping neutrons for the purpose of transmutation. For nuclear reactors with effective reflectors, the absorption rate of neutrons by the nuclei of the core shows enormous sensitivity to the inventory of those escaping (at least once) neutrons which spend time in the reflector.

To facilitate the present theory, it is convenient to express the effective neutron multiplication factor, k_{eff} , as the product of the infinite medium multiplication factor, k_{∞} , and the sum of two nonleakage probabilities, P_o , the total principal nonleakage probability, and P , the total secondary nonleakage probability. The first of these nonleakage probabilities, P_o , is defined as the fraction of neutrons which, having been produced by fission in the core, never traverse the core boundary. The second probability, P , is defined as the fraction of neutrons which, though

produced by fission in the core, spend time in the reflector but are eventually absorbed by nuclei of the core. Thus,

$$\begin{aligned}k_{\text{eff}} &= k_{\infty} (P_o + P) \\P_o &= P_o (C, R) \\P &= P(P_o, \text{core and reflector albedos})\end{aligned}\tag{2-1}$$

In Equation 2-1, C refers to core geometry and group constants, and R refers to reflector geometry and group constants.

Suppose there are m neutron energy groups indexed $1, 2, \dots, i, \dots, m$, where 1 is the group of highest energy and m represents the thermal group. One can define k_{eff} as follows

$$k_{\text{eff}} = \sum_{i=1}^m P_{\lambda_1}^f \frac{v_1 \Sigma_{f_1}}{\Sigma_{a_1}^f} \tag{2-2}$$

where

$P_{\lambda_1}^f$ is the fraction of fission neutrons in group i which are absorbed by fuel nuclei,

v_1 is the average number of fission neutrons released in a fission reaction induced by a neutron in group i ,

Σ_{f_1} is the macroscopic fission cross section for group i , and

$\Sigma_{a_1}^f$ is the macroscopic absorption cross section of the fuel for group i .

Albedos

Neutrons undergo reflection and absorption when they are incident on a scattering/absorbing medium. The reflection coefficient, or albedo, is defined as the current of the reflected neutrons divided by the incident neutron current, i.e.,

$$\text{albedo} = \frac{J_{\text{reflected}}}{J_{\text{incident}}} . \quad (2-3)$$

If the medium is of finite thickness, neutrons also are transmitted. The transmission coefficient, or transmittance, is similarly defined as the fraction of incident current that is transmitted through the medium

$$\text{transmittance} = \frac{J_{\text{transmitted}}}{J_{\text{incident}}} . \quad (2-4)$$

For the purpose of simplifying the development, it is assumed here that there is but one thermal group so that neutron upscattering can be neglected. The reflector-core neutron transfer is described by means of m -group albedos of which there are $\frac{m(m+1)}{2}$ types under the neutron energy transfer assumptions of multigroup diffusion theory, specifically,

- (1) The probability that a neutron from group 1 be reflected with group 1 energy, here denoted by ${}_1\alpha_1$ for core and ${}_1\beta_1$ for reflector.
- (2) The probability that a neutron from group 1 be reflected with group 2 energy, here denoted by

${}_1\alpha_2$ for core and ${}_1\beta_2$ for reflector.

...

...

- (i) The probability that a neutron from group i be reflected with group i' energy, here denoted by ${}_i\alpha_{i'}$ for core and ${}_i\beta_{i'}$ for reflector.

...

...

In matrix form, one can write the core (α) and reflector (β) albedos as follows

$$[\alpha] = \begin{bmatrix} {}_1\alpha_1 & {}_1\alpha_2 & {}_1\alpha_3 & \dots & {}_1\alpha_1 & \dots & {}_1\alpha_m \\ 0 & {}_2\alpha_2 & {}_2\alpha_3 & \dots & {}_2\alpha_1 & \dots & {}_2\alpha_m \\ 0 & 0 & {}_3\alpha_3 & \dots & {}_3\alpha_1 & \dots & {}_3\alpha_m \\ \vdots & \vdots & \vdots & & \vdots & & \vdots \\ 0 & 0 & 0 & \dots & {}_i\alpha_1 & \dots & {}_i\alpha_m \\ \vdots & \vdots & \vdots & & \vdots & & \vdots \\ 0 & 0 & 0 & \dots & 0 & \dots & {}_m\alpha_m \end{bmatrix} \quad (2-5)$$

$$[\beta] = \begin{bmatrix} {}_1\beta_1 & {}_1\beta_2 & {}_1\beta_3 & \dots & {}_1\beta_1 & \dots & {}_1\beta_m \\ 0 & {}_2\beta_2 & {}_2\beta_3 & \dots & {}_2\beta_1 & \dots & {}_2\beta_m \\ 0 & 0 & {}_3\beta_3 & \dots & {}_3\beta_1 & \dots & {}_3\beta_m \\ \vdots & \vdots & \vdots & & \vdots & & \vdots \\ 0 & 0 & 0 & \dots & {}_i\beta_1 & \dots & {}_i\beta_m \\ \vdots & \vdots & \vdots & & \vdots & & \vdots \\ 0 & 0 & 0 & \dots & 0 & \dots & {}_m\beta_m \end{bmatrix} \quad (2-6)$$

The neglect of upscattering leads to an upper triangular form for the albedo matrices.

Total Principal Nonleakage Probability, P_o

It is convenient to decompose P_o into its energy group contributions. Hence,

$$P_o = P_o (P_{o1}, P_{o2}, \dots, P_{o1}, \dots, P_{om}) = 1 - \sum_{i=1}^m {}^2S_i \quad (2-7)$$

where P_{oi} is the probability that group i neutrons never spend time in the reflector and 2S_i (henceforth termed the "source of initial leakage for group i neutrons") is defined as the fraction of group i neutrons that leaks from the core for the first time.

Source of Initial Leakage, 2S_i

The sources of initial leakage with the simplification of only two energy groups (i.e., $m = 2$) can be discussed with the aid of Figure 2-1.

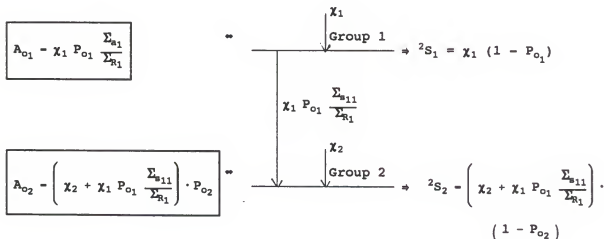


Figure 2-1. Schematic drawing showing the strategy for evaluation of the source of initial leakage under the two-group albedo approach.

In Figure 2-1,

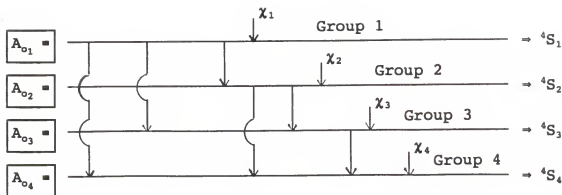
χ_i = the probability that a fission neutron will be born with an energy in group i . Note, $\chi_1 + \chi_2 = 1$.

A_{o_i} = the fraction of neutrons which, having been produced by fission in the core, never traverse the core boundary and are absorbed as group i neutron by nuclei of the core.

Of course, $P_o = A_{o_1} + A_{o_2}$.

$\Sigma_{R_i} - \Sigma_{t_i} - \Sigma_{s_{ii}}$ = the group i macroscopic removal cross section for the core, Σ_{t_i} is the group i macroscopic total cross section for the core, and $\Sigma_{s_{ii}}$ is the macroscopic group-transfer cross section from group i to group i for the core.

A straightforward method to arrive at expressions for the multigroup source of initial leakage equations is to apply the concept of neutron balance to each energy group. An algorithm for calculating mS_i values is created employing the known source of initial leakage equations for group i of the $(m-1)$ group albedo theory, ${}^{m-1}S_i$. Details of this idea are in Appendix A. For a small number of energy groups, $m \leq 10$, a heuristic schematic drawing is far more understandable and immediately useful. As an example, the 4S_i and A_{o_i} equations are given in Figure 2-2 using the four-group albedo method. These results are then used in subsequent four-group calculations.



where

$$^4S_1 = \chi_1 (1 - P_{o1}) \quad , \quad A_{o1} = \chi_1 P_{o1} \frac{\Sigma_{s1}}{\Sigma_{R1}} \quad ,$$

$$^4S_2 = \left(\chi_2 + \chi_1 P_{o1} \frac{\Sigma_{s12}}{\Sigma_{R1}} \right) (1 - P_{o2}) \quad ,$$

$$A_{o2} = \left(\chi_2 + \chi_1 P_{o1} \frac{\Sigma_{s12}}{\Sigma_{R1}} \right) P_{o2} \frac{\Sigma_{s2}}{\Sigma_{R2}} \quad ,$$

$$^4S_3 = \left[\chi_3 + \chi_1 P_{o1} \frac{\Sigma_{s13}}{\Sigma_{R1}} + \left(\chi_2 + \chi_1 P_{o1} \frac{\Sigma_{s12}}{\Sigma_{R1}} \right) P_{o2} \frac{\Sigma_{s23}}{\Sigma_{R2}} \right] (1 - P_{o3}) \quad ,$$

$$A_{o3} = \left[\chi_3 + \chi_1 P_{o1} \frac{\Sigma_{s13}}{\Sigma_{R1}} + \left(\chi_2 + \chi_1 P_{o1} \frac{\Sigma_{s12}}{\Sigma_{R1}} \right) P_{o2} \frac{\Sigma_{s23}}{\Sigma_{R2}} \right] P_{o3} \frac{\Sigma_{s3}}{\Sigma_{R3}} \quad ,$$

$$^4S_4 = \left[\chi_4 + \chi_1 P_{o1} \frac{\Sigma_{s14}}{\Sigma_{R1}} + \frac{^4S_2}{1 - P_{o2}} P_{o2} \frac{\Sigma_{s24}}{\Sigma_{R2}} + \frac{^4S_3}{1 - P_{o3}} P_{o3} \frac{\Sigma_{s34}}{\Sigma_{R3}} \right] (1 - P_{o4})$$

$$A_{o4} = \frac{^4S_4}{1 - P_{o4}} \quad , \quad \text{and}$$

$$P_o = \sum_{i=1}^4 A_{oi}$$

Figure 2-2. Schematic drawing showing the evaluation of the source of initial leakage considering the four-group albedo approach.

Figure 2-2 shows one fission neutron, i.e., $\chi_1 + \chi_2 + \chi_3 + \chi_4 = 1$, cascading downward in energy from group to group as it is moderated by scattering collisions, a fraction being absorbed by the nuclei of the core, and a fraction leaving the core region for the first time.

The Ping-Pong Decision Process

Consider now the fraction of neutrons which do spend time in the reflector, but are absorbed eventually in the core. This total secondary nonleakage probability, P , is the sum of m separate processes, specifically,

- (m) The probability that a neutron produced by fission in the core, leaks as thermal ($i = m$), and is absorbed in the core, here denoted by mP_m .

...

...

- (i) The probability that a neutron produced by fission in the core, leaks as a group i neutron, and is absorbed in the core, here denoted by mP_i .

...

...

- (1) The probability that a neutron produced by fission in the core, leaks as a group 1 neutron, and is absorbed in the core, here denoted by mP_1 .

In these terms,

$$P = \sum_{i=1}^m {}^mP_i \quad (2-8)$$

where the mP_i are termed "the effective partial secondary nonleakage probability."

It is convenient to decompose mP_i into its energy group contributions. Hence

$${}^mP_i = {}^mS_i \sum_{i'=1}^n {}^iQ_{i'} \quad (2-9)$$

where

${}^iQ_{i'}$, termed the fractional secondary nonleakage probability, is the probability that if one neutron leaks as a group i neutron from the core, and returns as a group i' neutron, it will be absorbed in the core. The partial secondary nonleakage probability, mQ_i , is defined by

$${}^mQ_i = \sum_{i'=1}^n {}^iQ_{i'} \quad (2-10)$$

Algorithm for mP_i ($m \geq 3$)

The following algorithm for calculating mQ_i is developed in a manner such that assumed known partial secondary nonleakage probabilities of $m-1$ group albedo theory imply results for m group theory. The starting expressions are from the two group albedo theory [2] and given by

$${}_2Q_1 = \frac{(1 - {}_1\alpha_1) {}_1\beta_1}{1 - {}_1\alpha_1 {}_1\beta_1} - \frac{{}_1\beta_1 {}_1\alpha_2}{1 - {}_1\alpha_1 {}_1\beta_1} \cdot \frac{(1 - {}_2\beta_2)}{1 - {}_2\alpha_2 {}_2\beta_2} +$$

$$\frac{{}_1\beta_2}{1 - {}_1\alpha_1 {}_1\beta_1} \cdot \frac{(1 - {}_2\alpha_2)}{1 - {}_2\alpha_2 {}_2\beta_2} \quad \text{and} \quad (2-11)$$

$${}^2Q_2 = \frac{(1 - {}^2\alpha_2) {}^2\beta_2}{1 - {}^2\alpha_2 {}^2\beta_2}.$$

Derivation of Equation 2-11 constitutes Appendix B.

First step: Evaluation of mP_m

Using the assumed known expression of ${}^{m-1}Q_{m-1}$, replace the subscript (m-1) by the subscript m; thus,

$${}^mP_m = {}^mS_m {}^mQ_m = {}^mS_m {}^{m-1}Q_{m-1} \quad (m-1 \rightarrow m) \quad (2-12)$$

or

$${}^mP_m = {}^mS_m \frac{(1 - {}^m\alpha_m) {}^m\beta_m}{1 - {}^m\alpha_m {}^m\beta_m}.$$

Second step: Evaluation of mP_i ($2 \leq i \leq m-1$)

$${}^mP_i = {}^mS_i {}^mQ_i = {}^mS_i {}^{m-1}Q_{i-1} \begin{pmatrix} i-1 \rightarrow i \\ \dots & \dots \\ \dots & \dots \\ m-1 \rightarrow m \end{pmatrix}. \quad (2-13)$$

Last Step: Evaluation of mP_1

For the evaluation of mP_1 , one can consider m separate processes given by

$${}^mP_1 = {}^mS_1 [{}^mQ_m + {}^1Q_{m-1} + \dots + {}^1Q_2 + {}^1Q_1]. \quad (2-14)$$

Keeping in mind that mQ_i is the fraction that is absorbed in the core if one neutron leaks as a group i neutron from the core, one can define $\bar{{}^mQ}_i$ as the fraction that is absorbed in

the reflector if one neutron leaks as a group i neutron from the reflector.

For computer implementation one can write

$$\bar{Q}_1 = Q_1 \begin{pmatrix} \alpha \rightarrow \beta \\ \beta \rightarrow \alpha \end{pmatrix}. \quad (2-15)$$

First process: Evaluation of Q_m

$$Q_m = \frac{1\beta_m}{1 - 1\alpha_1 1\beta_1} \left[1 - {}^{m-1}Q_{m-1} (m-1 \rightarrow m) \right]. \quad (2-16)$$

Second process: Evaluation of Q_{m-1}

$$Q_{m-1} = \frac{1\beta_{m-1}}{1 - 1\alpha_1 1\beta_1} \left[1 - {}^{m-2}Q_{m-2} \begin{pmatrix} m-2 \rightarrow m-1 \\ m-1 \rightarrow m \end{pmatrix} \right]. \quad (2-17)$$

(m-1)th process: Evaluation of Q_2

$$Q_2 = \frac{1\beta_2}{1 - 1\alpha_1 1\beta_1} \left[1 - {}^{m-1}Q_1 \begin{pmatrix} 1 \rightarrow 2 \\ 2 \rightarrow 3 \\ \dots \dots \dots \\ m-1 \rightarrow m \end{pmatrix} \right]. \quad (2-18)$$

(m)th process: Evaluation of Q_1

$$Q_1 = \frac{(1 - 1\alpha_1) 1\beta_1}{1 - 1\alpha_1 1\beta_1} - \frac{1\beta_1}{1 - 1\alpha_1 1\beta_1} \sum_{i=2}^m 1\alpha_i [1 - Q_i]. \quad (2-19)$$

Multigroup Neutron Albedo Approach
Flow Diagram

As the multigroup albedo approach becomes more complex, a flow diagram is helpful in planning and structuring a program for solving a calculational example. The basic features of the multigroup neutron albedo strategy are shown in the simplified flow diagram of Figure 2-3. The actual programs are contained in Appendix G.

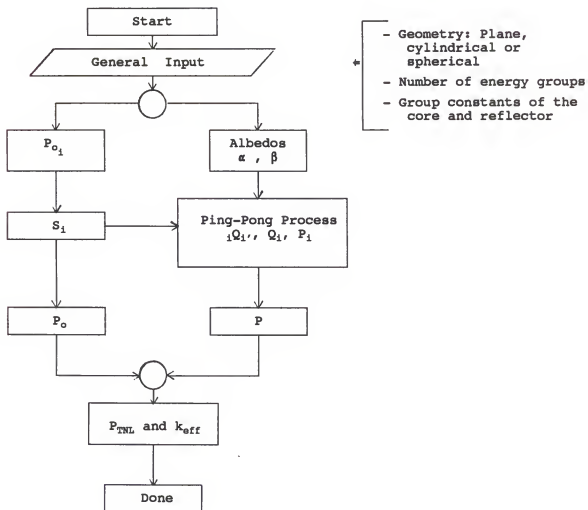


Figure 2-3. Simplified flow diagram of multigroup neutron albedo strategy.

CHAPTER III EVALUATION TECHNIQUES FOR THE CORE (α) AND REFLECTOR (β) ALBEDOS

Introduction

Neutrons undergo reflection when they are incident on a scattering/absorbing medium. The reflection coefficient, or albedo, is defined as the current of the reflected neutrons divided by the incident neutron current. If the medium is of finite thickness, neutrons are also transmitted. The transmission coefficient, or transmittance, is similarly defined as the fraction of incident current that is transmitted through the medium. Here the evaluation techniques for the albedo/transmittance values are confined to two different approaches: (1) low-order approximations to solution of the Boltzmann transport equation such as diffusion theory, and (2) first interaction free-flight transport methods.

Albedo Calculation Using the Neutron Diffusion Approximation

The diffusion approximation to the Boltzmann transport equation is sufficiently accurate for the calculation of neutron densities in systems which are homogeneous, or nearly homogeneous, and which are so large that the radius of curvature of their boundaries is almost everywhere large

compared to the mean free path of the neutrons. These conditions are satisfied in large solid or liquid reactors of the usual, rather regular, shapes.

In order to demonstrate the considerations involved, the case of a spherical core with an infinite reflector and example albedo calculations for ${}_1\alpha_4$ and ${}_2\beta_3$ are here outlined.

Determination of ${}_1\alpha_4$

The governing equations and boundary conditions for the determination of ${}_1\alpha_4$ are (see Figure 3-1)

$$\begin{aligned}
 & -D_1 \nabla^2 \phi_1 + \Sigma_{R1} \phi_1 = 0 \\
 & -\Sigma_{S12} \phi_1 - D_2 \nabla^2 \phi_2 + \Sigma_{R2} \phi_2 = 0 \\
 & -\Sigma_{S13} \phi_1 - \Sigma_{S23} \phi_2 - D_3 \nabla^2 \phi_3 + \Sigma_{R3} \phi_3 = 0 \\
 & -\Sigma_{S14} \phi_1 - \Sigma_{S24} \phi_2 - \Sigma_{S34} \phi_3 - D_4 \nabla^2 \phi_4 + \Sigma_{R4} \phi_4 = 0
 \end{aligned} \tag{3-1}$$

At $r = 0$, ϕ_1 , ϕ_2 , ϕ_3 , and ϕ_4 are finite.

At $r=R$, $J_{-1-1} = 1$, $J_{-1-2,3,4} = 0$

where J_i are the partial current densities and all group parameters (D_i , Σ_{Ri} , and Σ_{Sii}) refer to the homogeneous

core. At $r = R$, ${}_1\alpha_4 = \frac{J_{-1-4}}{J_{-1-1}} = \frac{\phi_4}{4} - \frac{D_4}{2} \frac{d}{dr} \phi_4$. The results

of this calculation as well as all other four-group albedos used in this work are presented in Appendices C and D.

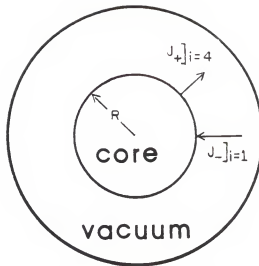


Figure 3-1. Schematic drawing showing the evaluation of $1\alpha_4$.

Determination of $2\beta_3$

The governing equations and boundary conditions for determination of $2\beta_3$ are (see Figure 3-2)

$$- D_2 \nabla^2 \phi_2 + \Sigma_{R_2} \phi_2 = 0$$

(3-2)

$$- \Sigma_{S_{23}} \phi_2 - D_3 \nabla^2 \phi_3 + \Sigma_{R_3} \phi_3 = 0$$

As $r \rightarrow \infty$, ϕ_2 and ϕ_3 are finite.

At $r=R$, $J_{+}]_{i=2} = 1$, $J_{+}]_{i=3} = 0$

where J_i are the partial current densities and all group parameters (D_i , Σ_{R_i} , $\Sigma_{S_{ii}}$) refer to the homogeneous

reflector. At $r = R$, $2\beta_3 = \frac{J_{-}]_{i=3}}{J_{+}]_{i=2}} = \frac{\phi_3}{4} + \frac{D_3}{2} \frac{d}{dr} \phi_3$.

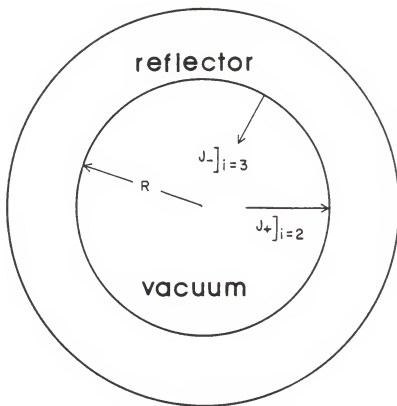


Figure 3-2. Schematic drawing showing the evaluation of ${}_2\beta_3$.

First-Interaction Transmission Coefficients

If the radius of curvature of boundaries is small compared to the mean free path of neutrons such as is the case for gas core reactors, core albedo considerations are changed to core transmittance considerations. Transmission coefficients can be evaluated following a suggestion by Tsoulfanidis [5] as follows: Consider a point isotropic monoenergetic source to be on the surface of a spherical core of diameter $2R$ (see Figure 3-3). For neutrons emitted from the reflector at an angle ω , measured from the axis or diameter of the spherical core, the probability of interaction in the core is $1 - \exp[-\Sigma_t \times 2R \times \cos \omega]$. The probability of emission between angles ω and $\omega + d\omega$ is $\frac{\sin \omega}{2} d\omega$. Σ_t is the core total macroscopic cross section. The core transmittance, γ , is given by

$$\gamma = \frac{\int_{\omega=0}^{\pi/2} \exp[-\Sigma_t \times 2R \times \cos \omega] \frac{\sin \omega}{2} d\omega}{\int_{\omega=0}^{\pi/2} \frac{\sin \omega}{2} d\omega} \quad (3-3)$$

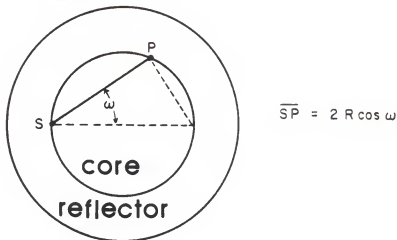


Figure 3-3. Schematic drawing for evaluation of the core transmission coefficients due to a point isotropic neutron source on a spherical core.

Other Approaches for
Albedo/Transmittance Evaluation

Here are presented additional options for albedo and transmittance calculation. Analytical studies of the albedo equations in the theory of the diffusion and slowing down of neutrons were carried out by Orlov [6] in 1961. When the invariant imbedding theory [7, 8] is applied to the Boltzmann equation of neutron transport some of quantities of interest become the reflection function and the transmission function. The mathematical theory of the invariant imbedding approach has been studied by Wing [9] especially for the reflection probability.

Finally, in this work, the core and reflector albedo values are obtained employing the multigroup neutron diffusion theory and the core transmittance results are evaluated using the first-interaction, free-flight transport method.

CHAPTER IV
EVALUATION TECHNIQUES FOR THE PARTIAL PRINCIPAL
NONLEAKAGE PROBABILITY, P_{o_i}

Introduction

The initial nonleakage probability, P_{o_i} , is the probability that neutrons from group i never spend time in the reflector. The analytical expression of P_{o_i} in terms of the core and reflector parameters involve analytical complications. Here, two evaluation approaches for the initial nonleakage probabilities are suggested: (1) a bare core approximation, and (2) a consistent multigroup diffusion theory method.

Basic Ideas

As an introductory, heuristic example, consider the case of one-group analysis in an infinite slab core reflected by infinitely thick media. The one-group neutron flux in the core varies as $\cos(Bx)$, where B^2 is the buckling of the core material. Therefore, $v\Sigma_f \cos(Bx)$ is the source function which should be placed in a convolution with a nonreentrant flux kernel in order to determine the neutron flux which has never been in the reflector. The

kernel, $K(x)$, for this component of the flux satisfies the differential equation

$$D \frac{d^2 K}{dx^2} - \Sigma_a K = 0 \quad (4-1)$$

and has the boundary conditions that $K(x)$ be continuous and that the resultant reentrant current of neutrons to the core at the core-reflector interfaces be zero. The second boundary condition is approximated by requiring $K(x)$ to vanish at the extrapolated boundaries of the core, i.e., $K(x = \tilde{a}) = 0$ where $\tilde{a} = a + 2.1312 D$ and a is the position of the core-reflector interface (see Figure 4-1).

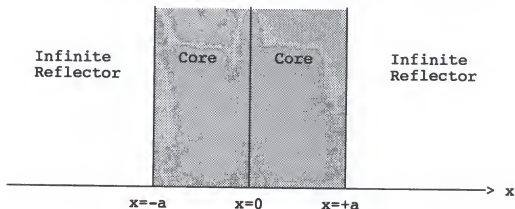


Figure 4-1. Reflected slab reactor.

With a unit strength source plane at $x = x_0$, and the boundary conditions

$$\left[D \frac{dK_-}{dx} - D \frac{dK_+}{dx} \right]_{x=x_0} = 1,$$

$$[K_+]_{x=x_0} = [K_-]_{x=x_0}, \text{ and} \quad (4-2)$$

$$[K_+]_{x=-\bar{a}} = 0 = [K_-]_{x=\bar{a}}$$

the solution for $K(x)$ is

$$K_- = K(x) = \frac{\sinh [c_1 (\bar{a} + x)]}{c_1 D \sinh [c_1 (\bar{a} + x_0)] [\coth [c_1 (\bar{a} - x_0)] + \coth [c_1 (\bar{a} + x_0)]]}$$

$$\text{for } -a \leq x < x_0 \quad (4-3)$$

$$K_+ = K(x) = \frac{\sinh [c_1 (\bar{a} - x)]}{c_1 D \sinh [c_1 (\bar{a} - x_0)] [\coth [c_1 (\bar{a} - x_0)] + \coth [c_1 (\bar{a} + x_0)]]}$$

$$\text{for } x_0 < x \leq a$$

where

$$c_1 = \sqrt{\frac{\Sigma_a}{D}}.$$

Thus, the solution for the flux component which has never been in the reflector, denoted here by ϕ_0 , is

$$\phi_0(x) = \int_{x_0-a}^a v \Sigma_t \cos(Bx_0) K(x, x_0) dx_0. \quad (4-4)$$

Under the one-group assumption, the initial nonleakage probability, P_{o1} , is given in terms of ϕ_0 , as

$$P_{o1} = \left[1 + \frac{J_{\phi_0}(x=a)}{\Sigma_a \int_{x=0}^a \phi_0 dx} \right]^{-1} \quad (4-5)$$

Notice that it is the ratio of the outgoing neutron current at the core boundary to the absorption rate of the flux which has not traversed the core boundary which is of importance.

This one-group treatment is extended to a multigroup evaluation of P_{o1} later in this chapter.

Bare Core Approximation

A simplification suggested by Jacobs [1] is to use for P_{o1} the initial nonleakage probabilities which would be applicable to a bare core reactor with the dimensions of the core of the reflected reactor. One can write

$$P_{o1} = \left[1 + (B')^2 \frac{D_1}{\Sigma_{R1}} \right]^{-1} \quad (4-6)$$

where $(B')^2$ is the geometrical buckling of a bare core of the same dimensions as the reflected core, Σ_{R1} are the core macroscopic removal cross sections, and D_1 are the core diffusion coefficients.

Consistent Multigroup Diffusion Theory Method

The initial nonleakage probabilities, P_{o1} , are given by

$$P_{o1} = \left[1 + B_1^2 \frac{D_1}{\Sigma_{R1}} \right]^{-1} \quad (4-7)$$

under the multigroup diffusion approximation. In Equation (4-7),

$$B_1^2 = - \frac{\int_S \nabla \phi_{o1} \cdot dS}{\int_V \phi_{o1} dV}$$

$$\Sigma_{R1} = \Sigma_{t1} - \Sigma_{s11}$$

Σ_{R1} are the core macroscopic removal cross sections, Σ_{t1} are the core total macroscopic cross sections, Σ_{s11} are the core self-group-transfer cross sections, D_1 are the core diffusion coefficients, ϕ_{o1} is the component of the i group neutron flux which has not been in the reflector, and the

integrations are over the volume and the surface of the reactor core.

In spherical geometry (i.e., a spherical core in an infinite reflector) the ϕ_{o_i} are given by

$$\phi_{o_i}(r) = \phi_i(r) - \frac{A_i}{r} \sinh \left(\sqrt{\frac{\Sigma_{R_i}}{D_i}} r \right) \quad (4-8)$$

and subject to the boundary conditions at $r = R$, $J_{o_i} = 0$.

The i group neutron flux, $\phi_i(r)$, includes all neutrons whether or not they have spent time in the reflector.

Because of the algebraically complicated developments involved, solution of $\phi_{o_i}(r)$ and P_{o_i} is relegated to Appendix E. The results listed there are used in the remainder of this work.

CHAPTER V
STATIC, ONE-DIMENSIONAL NEUTRONIC ANALYSIS OF A LARGE
OPTICAL-PATH-LENGTH CORE REACTOR (TWO-REGION
SOLID CORE REACTOR)

Introduction

In this section, the main parameters of the four-group albedo theory are investigated with the aid of a calculational example. A solid core reactor with a water moderator/reflector is studied as a first model case. This problem is selected because both S_n transport theory and diffusion theory approximations can be successfully applied and give essentially the same results [10].

Core Atomic Densities

As a first example problem consider a large optical-path-length spherical core reactor surrounded by an infinite water reflector where it is assumed that the core materials are a homogeneous mixture. In Table 5-1 are the core atomic densities employed.

TABLE 5-1. CORE ATOMIC DENSITIES OF THE SOLID CORE REACTOR PROBLEM

Material	Atomic Densities (atoms/(barn x cm))
Uranium - 235	.12200 E-03
Uranium - 238	.59700 E-02
Oxygen - 16	.34420 E-01
Natural Chromium	.93460 E-03
Manganese - 55	.94200 E-04
Natural Iron	.33470 E-02
Natural Nickel	.47110 E-03
Hydrogen - 1	.44470 E-01

Core and Reflector Group Constants

For this particular problem, NITAWL [11] is used to generate a 123 fine group neutron cross-section library. The 123 fine group neutron cross-sections (30 thermal groups) are collapsed down to a 26 energy-group structure (6 thermal groups) and then to a 4 energy-group structure (one thermal group) with XSDRNPM. Tables 5-2 and 5-3 show the core and reflector group constants, respectively.

TABLE 5-2. CORE GROUP CONSTANTS OF THE SOLID CORE REACTOR PROBLEM FROM THE XSDRNPM CODE

	Four Energy-Group Structure			
	i = 1	i = 2	i = 3	i = 4
D	.17607 E+01	.80339	.47001	.19923
$\Sigma_{n,2n}$.13954 E-03	.00000	.00000	.00000
χ	.74415	.25565	.20189 E-03	.12480 E-08
Σ_a	.33928 E-02	.18935 E-02	.17635 E-01	.57172 E-01
$\nu\Sigma_f$.72250 E-02	.51635 E-02	.59613 E-02	.66730 E-01
Σ_t	.27502	.71374	.13068 E+01	.22259 E+01
i - i'	$\Sigma_{s,ii'}$			
	i' = 2	i' = 3	i' = 4	
1	.89651 E-01	.46418 E-03	.15529 E-06	
2	-	.95330 E-01	.31330 E-04	
3	-	-	.98090 E-01	

- (D in cm and Σ in cm^{-1})
- Radius of the spherical core = 64.0 cm.

From this table, the ratio of core diameter to the total mean free paths in the solid core is found to be very large (~285.0 for thermal neutrons), thus, this solid core has indeed very large optical dimensions.

TABLE 5-3. REFLECTOR GROUP CONSTANTS OF THE SOLID CORE REACTOR PROBLEM FROM THE XSDRNP CODE

	Four Energy-Group Structure			
	i = 1	i = 2	i = 3	i = 4
D	.18109 E+01	.78453	.50770	.14915
Σ_a	.31290 E-03	.95302 E-05	.57242 E-03	.15539 E-01
Σ_c	.27073	.85161	.15510 E+01	.30612 E+01
i \rightarrow i'	$\Sigma_{s_{ii'}}$			
	i' = 2	i' = 3	i' = 4	
1	.11270	.69381 E-03	.23278 E-06	
2	-	.14163	.46992 E-04	
3	-	-	.14601	

- (D in cm and Σ in cm^{-1})

Core (α) and Reflector (β) Albedos

The diffusion theory approximation within the solid core reactor is valid. Thus, Tables 5-4 and 5-5 show the core and reflector albedo values obtained using the expressions in Appendices C and D, respectively.

TABLE 5-4. CORE ALBEDOS OF THE SOLID CORE REACTOR PROBLEM USING THE FOUR-GROUP ALBEDO THEORY

i \rightarrow i'	$\alpha_{ii'}$			
	i' = 1	i' = 2	i' = 3	i' = 4
1	.13824	.45412	.95613 E-01	.10531
2	-	.30383	.20285	.87856 E-01
3	-	-	.37764	.21696
4	-	-	-	.65669

- Radius of the spherical core = 64.0 cm.

TABLE 5-5. REFLECTOR ALBEDOS OF THE SOLID CORE REACTOR PROBLEM USING THE FOUR-GROUP ALBEDO THEORY

$i \rightarrow i'$	${}_1\beta_{i'}$			
	$i' = 1$	$i' = 2$	$i' = 3$	$i' = 4$
1	.18267 E-01	.20087	.90822 E-01	.14104
2	-	.18251	.22292	.21026
3	-	-	.28086	.39499
4	-	-	-	.81662

- Radius of the spherical core = 64.0 cm.
- Infinite water reflector.

From this table, the probability that a neutron from group 1 be reflected by the reflector is found to be small (${}_1\beta_1 + {}_1\beta_2 + {}_1\beta_3 + {}_1\beta_4 = .45$). Thus, 55% of the incident neutrons from group 1 are absorbed in the reflector.

Partial Principal Nonleakage Probabilities, P_{o_i}

The evaluation of the partial principal nonleakage probabilities constitute the key item for a multigroup albedo theory technique. The diffusion theory approximation within the solid core reactor is valid. In Table 5-6 are the partial principal nonleakage probability, P_{o_i} , and source of initial leakage, S_i , values obtained using the consistent multigroup diffusion theory method (which should be valid). Values are obtained using the expressions summarized in Appendix E.

TABLE 5-6. PARTIAL PRINCIPAL NONLEAKAGE PROBABILITIES, P_{o_i} , AND SOURCES OF INITIAL LEAKAGE, S_i , OF THE SOLID CORE PROBLEM USING THE FOUR-GROUP ALBEDO THEORY METHOD

	Four-Group Albedo Theory Method			
	$i = 1$	$i = 2$	$i = 3$	$i = 4$
P_{o_i}	.94055	.97792	.98077	.96782
S_i	.44242 E-01	.20447 E-01	.17131 E-01	.23846 E-01

- Radius of the spherical core = 64.0 cm.
- Infinite water reflector.

From this table, the partial principal nonleakage probabilities, P_{o_i} , are very large and the sources of initial leakage, S_i , are very small. Thus, this solid core reactor problem shows large total principal nonleakage probability, P_o , i.e., $P_o = 1 - (S_1 + S_2 + S_3 + S_4) = .89433$. In other words, 89.43% of the neutrons which are produced by fission in the core never traverse the core boundary.

The Ping-Pong Decision Process

Considering now the fraction of neutrons which do spend time in the reflector, in Table 5-7 are the fractional secondary nonleakage probability, ${}_1Q_i$, values obtained using the ping-pong decision process developed in Chapter II. The program ALB1 used for these numerical evaluations is in Appendix G.

TABLE 5-7. FRACTIONAL SECONDARY NONLEAKAGE PROBABILITIES, ${}_1Q_{i'}$, OF THE SOLID CORE PROBLEM USING THE PING-PONG DECISION PROCESS

$i \rightarrow i'$	${}_1Q_{i'}$			
	$i' = 1$	$i' = 2$	$i' = 3$	$i' = 4$
1	.95088 E-02	.13974	.65900 E-01	.10468
2	-	.10804	.17081	.16480
3	-	-	.16858	.32712
4	-	-	-	.60456

- Radius of the spherical core = 64.0 cm.
- Infinite water reflector.

From Tables 5-6 and 5-7, the total secondary nonleakage probability, P , is given by

$$P = \sum_{i=1}^4 \left[S_i \sum_{i'=1}^4 {}_1Q_{i'} \right] = .46129 \text{ E-01} .$$

Thus, 4.61% of the neutrons which are produced by fission in the core and spend time in the reflector are eventually absorbed by nuclei of the core.

Dissection of P_o and P

It is interesting to know (1) the fraction of neutrons produced by fission in the core, which never traverse the core boundary and are absorbed as group i , here denoted by $A_{o,i}$, and (2) the fraction of neutrons which are produced by fission in the core, spend time in the reflector and are eventually absorbed by nuclei of the core as group i , here denoted by A_i .

In Table 5-8, are the probabilities, A_{o_1} and A_1 , where

$$P_o = \sum_{i=1}^4 A_{o_i} \text{ and } P = \sum_{i=1}^4 A_i .$$

The numerical values are found using the program ALB1 described in Appendix G.

TABLE 5-8. NONLEAKAGE PROBABILITIES, A_{o_1} AND A_1 , OF THE SOLID CORE REACTOR PROBLEM USING THE FOUR-GROUP ALBEDO THEORY METHOD

	Four-Group Albedo Theory Method			
	i = 1	i = 2	i = 3	i = 4
A_{o_1}	.26402 E-01	.17628 E-01	.13317	.71714
A_1	.63172 E-05	.10933 E-03	.18880 E-02	.44126 E-01

From Table 5-8, the total nonleakage probability, P_{TNL} , is found to be large, approximately 94%, i.e.,

$$P_{TNL} = \sum_{i=1}^4 (A_{o_i} + A_i) = .94046 .$$

About 94% of the neutrons produced by fission in the core are absorbed by nuclei of the core.

Effective Neutron Multiplication Factor, k_{eff}

The neutrons produced by the (n,2n) interaction in the solid core are accounted for in the k_{eff} evaluation. Thus, k_{eff} is given by

$$k_{eff} = (A_{o_1} + A_1) \frac{v_1 \Sigma_{f_1} + 2 \Sigma_{n,2n}}{\Sigma_{a_1}} + \sum_{i=2}^4 (A_{o_i} + A_i) \frac{v_1 \Sigma_{f_i}}{\Sigma_{a_1}} , \quad (5-1)$$

i.e., $k_{eff} = .99512$.

Comparison: XSDRNPM Versus Albedo Method Results

A one-dimensional static neutronic analysis has been performed on the solid core reactor problem with the XSDRNPM computer code. The $(A_{o_1} + A_1)$ results performed with XSDRNPM and albedo method are compared with one another in Table 5-9.

TABLE 5-9. COMPARISON OF $(A_{o_1} + A_1)$ RESULTS OF THE SOLID CORE REACTOR PROBLEM USING THE XSDRNPM COMPUTER CODE AND THE ALBEDO METHOD

	$(A_{o_1} + A_1)$			
	i = 1	i = 2	i = 3	i = 4
XSDRNPM	.23226 E-01	.17755 E-01	.13548	.76360
Albedo	.26409 E-01	.17738 E-01	.13505	.76126

- Radius of the spherical core = 64.0 cm
- Infinite water reflector for the albedo results
- Reflector 100 cm thick for XSDRNPM, S_4 P_3 , calculation
- XSDRNPM parameters: Four-group, S_4 P_3 , number of spacial intervals = 60, overall convergence = 1.0 E-04
- XSDRNPM calculation on an IBM 3090-600J

Note that in Table 5-9 the values of $(A_{o_1} + A_1)$ are compared because the XSDRNPM computer code does not give A_{o_1} and A_1 values separably as does the albedo method.

In Table 5-10 are the total nonleakage probability, P_{TNL} , and effective neutron multiplication factor, k_{eff} , for the system as obtained from XSDRNPM, four-group, S_4 P_3 calculation and the four-group albedo method.

TABLE 5-10. TOTAL NONLEAKAGE PROBABILITY, P_{TNL} , AND EFFECTIVE NEUTRON MULTIPLICATION FACTOR, k_{eff} , VALUES OF THE SOLID CORE REACTOR PROBLEM USING THE XSDRNPM COMPUTER CODE AND THE ALBEDO METHOD

	XSDRNPM	Albedo Method
P_o	-	.89433
P	-	.46129 E-01
$P_{TNL} = P_o + P$.94256	.94046
k_{eff}	.99667	.99512
CPU Time	8.0 sec	< 1 sec

- Radius of the spherical core = 64.0 cm
- Infinite water reflector for the albedo results
- Reflector 100 cm thick for XSDRNPM calculation
- XSDRNPM parameters: Four-group, S_4 P_3 , number of spatial intervals = 60, overall convergence = 1.0 E-04
- XSDRNPM calculation on an IBM 3090-600J

In order to avoid relying on a mainframe system, a personal computer version of the four-group albedo approach program is created and written in FORTRAN-77 (see Appendix G). Table 5-10 shows that, for the examined configuration, the albedo theory approximation is found to underpredict k_{eff} by about 0.2% relative to the S_4 P_3 result.

For various core density factors, the k_{eff} results performed with XSDRNPM are compared with albedo method in Table 5-11. These density factors are used to effect a density variation in a mixture. Zero for a density factor affords a convenient way for mocking a void region.

TABLE 5-11. COMPARISON OF k_{eff} RESULTS OF THE SOLID CORE REACTOR PROBLEM USING VARIOUS CORE DENSITY FACTORS (D.F.) FROM THE XSDRNPM CODE AND THE FOUR-GROUP ALBEDO METHOD

D.F.	k_{eff}		k_{eff} Relative Error	CPU Time	
	XSDRNPM	Albedo		XSDRNPM	Albedo
1.2	1.0133	1.0120	.128 E-02	9.7 sec	< 1 sec
1.1	1.0059	1.0044	.149 E-02	8.8 sec	< 1 sec
1.0	.99667	.99512	.155 E-02	8.0 sec	< 1 sec

- Radius of the spherical core = 64.0 cm
- Infinite water reflector for the albedo results
- Reflector 100 cm thick for XSDRNPM calculation
- XSDRNPM parameters: Four-group, S_4 P_3 , number of spacial intervals = 60, overall convergence = 1.0 E-04
- XSDRNPM calculation on an IBM 3090-600J

$$- k_{eff} \text{ relative error} = \frac{k_{eff} (\text{XSDRNPM}) - k_{eff} (\text{albedo})}{k_{eff} (\text{XSDRNPM})}$$

Table 5-11 shows that for the density factors used, the albedo theory approximation is found to underpredict k_{eff} by about 0.2% relative to the S_4 P_3 results. An examination of the calculation times, in this table, shows that albedo method is much faster than the S_4 P_3 approximation.

The main reason for the disagreements between results obtained from XSDRNPM and albedo calculations is due to the

methodology. XSDRNPM solves the one-dimension Boltzmann transport equation using a discrete ordinates (S_n) approximation method while the albedo theory approach is an analytical and intuitive method of determining the nonleakage probabilities which also involves easily identified approximations. The four-group albedo method does lead to a very satisfactory result with an analytic and easily understandable solution. Calculations of the sensitivity of these results to variations in the relevant approximations is relegated to the next chapters which deal with gas core reactors. It is in the application to gas core reactors that the albedo method yields calculational approaches which could have major significance.

CHAPTER VI
STATIC, ONE-DIMENSIONAL NEUTRONIC ANALYSIS OF A SMALL
OPTICAL-PATH-LENGTH CORE REACTOR
(TWO-REGION GAS CORE REACTOR)

Introduction

Multigroup neutron albedo methods are developed in Chapters II, III, and IV, and then applied to a solid core reactor in Chapter V. The multigroup albedo technique is a generalization of the method used by Jacobs [2] to successfully predict and explain the origin of positive temperature coefficient effects in "swimming pool" research reactors, which have large optical dimensions. The ideas employed in the multigroup neutron albedo approach are actually more useful and applicable to gaseous cores than to the original "solid core" application.

A significant number of studies have been performed on many different gas core reactor concepts over the past 30 years. These studies have focused primarily on externally-moderated, gaseous core reactor designs and have generally found standard diffusion theory to be invalid in the gas core region since the gas core dimensions rarely exceed a mean free path [10].

Core and Reflector Atomic Densities

Consider a small optical-path-length core reactor surrounded by an infinite beryllium oxide reflector where it is assumed that the materials are a homogeneous mixture in the spherical core. In Table 6-1 are the core and reflector atomic densities employed.

TABLE 6-1. CORE AND REFLECTOR ATOMIC DENSITIES OF THE GAS CORE REACTOR PROBLEM

Material	Atomic Densities (atoms/(barn x cm))
Uranium - 235	.10400 E-04
Uranium - 238	.18350 E-05
Fluorine - 19	.48930 E-04
Be in BeO	.72230 E-01
Oxygen - 16	.72230 E-01

Core and Reflector Group Constants

For this problem, NITAWL [11] is used to generate a 123 fine group neutron cross-section library. The 123 fine group neutron cross-sections (30 thermal groups) are collapsed down to a 26 energy-group structure (6 thermal groups) and then to a 4 energy-group structure (one thermal group) with XSDRNPM. In Tables 6-2 and 6-3 are the core and reflector group constants, respectively.

TABLE 6-2. CORE GROUP CONSTANTS OF THE GAS CORE REACTOR
PROBLEM FROM THE XSDRNPM CODE

	Four Energy-Group Structure			
	i = 1	i = 2	i = 3	i = 4
2	.19855 E+04	.86368 E+03	.44596 E+03	.94931 E+02
$\Sigma_{n, 2n}$.11286 E-06	.00000	.00000	.00000
χ	.74351	.25629	.20253 E-03	.12520 E-08
Σ_a	.15107 E-04	.23846 E-04	.39035 E-03	.31475 E-02
$v\Sigma_f$.36552 E-04	.43809 E-04	.53687 E-03	.64360 E-02
Σ_t	.22733 E-03	.43198 E-03	.39035 E-03	.35183 E-02
i-i'	$\Sigma_{s_{ii'}}$			
	i' = 1	i' = 2	i' = 3	i' = 4
1	.18190 E-03	.30356 E-04	.27863 E-08	.16460 E-11
2	-	.40549 E-03	.24982 E-05	.12203 E-10
3	-	-	.36193 E-03	.21137 E-05
4	-	-	-	.37083 E-03

- D in cm and Σ in cm^{-1}

- Radius of the spherical core = 70.0 cm

From Table 6-2, the ratio of core diameter to the total mean free paths in the gas core is found to be very small (~ 0.5 for thermal neutrons), thus, this gas core has small optical dimensions.

TABLE 6-3. REFLECTOR GROUP CONSTANTS OF THE GAS CORE REACTOR PROBLEM FROM THE XSDRNPM CODE

	Four Energy-Group Structure			
	i = 1	i = 2	i = 3	i = 4
D	.11681 E+01	.54420	.51029	.37525
$\Sigma_{n, 2n}$.58851 E-02	.00000	.00000	.00000
Σ_a	.29134 E-02	.80982 E-05	.21532 E-04	.38769 E-03
Σ_t	.37032	.66685	.70088	.99156
i-i'	$\Sigma_{s_{ii'}}$			
	i' = 2	i' = 3	i' = 4	
1	.62765 E-01	.15147 E-04	.17131 E-11	
2	-	.21315 E-01	.00000	
3	-	-	.14595 E-01	

- D in cm and Σ in cm^{-1}

From Table 6-3, the ratio of $\Sigma_{n, 2n}/(\Sigma_a + \Sigma_{n, 2n})$ for group 1 neutrons in the reflector is found to be very large (~0.7). This observation indicates that for the evaluation of the reflector albedos the neutron production via the $\text{Be}^9 (n, 2n) \text{Be}^8$ reaction must be considered.

Reflector (β) Albedos

The diffusion theory approximation within the infinite beryllium oxide reflector is valid. The neutrons produced by the $\text{Be}^9 (n, 2n) 2 \text{He}^4$ interaction in the reflector are accounted for in ${}_1\beta_1$, ${}_1\beta_2$, ${}_1\beta_3$, and ${}_1\beta_4$ results (see Appendix G). In Table 6-4 are the reflector albedo values obtained using the diffusion approximation (see Appendix D).

TABLE 6-4. REFLECTOR ALBEDOS OF THE GAS CORE REACTOR PROBLEM USING THE FOUR-GROUP ALBEDO THEORY

$i \rightarrow i'$	${}_1\beta_{i'}$			
	$i' = 1$	$i' = 2$	$i' = 3$	$i' = 4$
1	.28040	.30778	.11624	.18519
2	-	.62470	.15887	.13476
3	-	-	.68449	.23593
4	-	-	-	.93266

- Radius of the spherical core = 70.0 cm
- Infinite beryllium oxide reflector

From this table, the probability that a neutron from group 1 is reflected by the reflector is found to be large (${}_1\beta_1 + {}_1\beta_2 + {}_1\beta_3 + {}_1\beta_4 \approx .89$). Thus, only 11% of the group 1 neutrons impinging on the reflector are absorbed by nuclei of the BeO reflector. However, these neutron absorptions are responsible for the $n, 2n$ gain in the system.

The Ping-Pong Decision Process

In the albedo method as applied to a gas core reactor, the description of average neutron histories is accomplished through the specification of reflector albedos and core transmittances at each neutron traversal of the core-reflector interface.

Strategy for the Gas Core Reactor Problem

Considering the fission neutron energy distribution in Table 6-2, essentially all fission neutrons are born in the first two energy groups, i.e., above 5.53 keV and below 14.92 MeV. Figure 6-1 represents the average neutron history of a group 1 neutron source in the spherical core surrounded by an infinite reflector. Because of the lengthy algebraic manipulations involved, only a sketch of the method of determining the nonleakage probabilities is presented here.

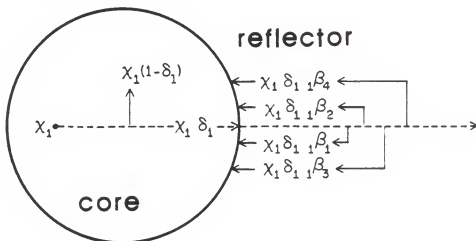


Figure 6-1. Schematic drawing showing the albedo method strategy for the gas core reactor problem.

For example, the first source of initial leakage neutrons for the evaluation of P_0 and P is χ_1 . The fraction of these neutrons transmitted through the core as group 1 is δ_1 . The quantities returned as group 1, 2, 3, and 4 on the first core boundary traverse are $\chi_1 \delta_{1-1}\beta_1$, $\chi_1 \delta_{1-1}\beta_2$, $\chi_1 \delta_{1-1}\beta_3$, and $\chi_1 \delta_{1-1}\beta_4$, respectively. The neutron fraction, $\chi_1 (1 - \delta_1)$, interacts in the core and can be dissected into:

- (1) The fraction absorbed in the core in group 1, i.e.,

$$\chi_1 (1 - \delta_1) \frac{\Sigma_{a1}}{\Sigma_{t1}},$$

- (2) The fraction scattered in the core into group 1, i.e.,

$$\chi_1 (1 - \delta_1) \frac{\Sigma_{s11}}{\Sigma_{t1}} = I_1$$

- (3) The fraction scattered in the core into group 2, i.e.,

$$\chi_1 (1 - \delta_1) \frac{\Sigma_{s12}}{\Sigma_{t1}} = I_2$$

- (4) The fraction scattered in the core into group 3, i.e.,

$$\chi_1 (1 - \delta_1) \frac{\Sigma_{s13}}{\Sigma_{t1}} = I_3$$

- (5) The fraction scattered in the core into group 4, i.e.,

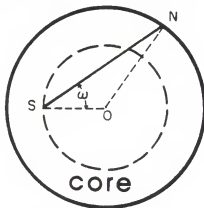
$$\chi_1 (1 - \delta_1) \frac{\Sigma_{s14}}{\Sigma_{t1}} = I_4.$$

Thus, in this first run, there are created four core secondary sources, whose strengths are I_1 , I_2 , I_3 , and I_4 , and four reflector secondary sources. $\chi_1 (1 - \delta_1) \frac{\Sigma_{a1}}{\Sigma_{t1}}$ is the quantity absorbed in the core as group 1.

For the next steps, it is required to consider the presence of all secondary sources in order to estimate the total quantities that will be absorbed as groups 1, 2, 3, and 4 in the core. To help understand the evaluation of P_0 and P , the contribution of one group 4 neutron source is detailed in Appendix F.

Transmission Coefficients for the Core and Reflector Secondary Sources

The core transmission coefficients for the reflector secondary sources for each energy group, γ_i , are evaluated as described in Chapter III. By calculating the core transmission coefficients for the core secondary sources for each energy group, δ_i , it is assumed that: (1) the neutron sources are isotropic and (2) the neutrons are homogeneously distributed at the surface of a sphere of radius R_0 given by $R_0 = \left[\frac{R_c^3}{2} \right]^{1/3}$, i.e., this imaginary sphere has half volume of the core. With the aid of Figure 6-2, the evaluation of δ_i values can be visualized.



$$R_c = ON,$$

$$R_0 = OS, \text{ and}$$

$$SN = R_0 \cos \omega + \sqrt{R_c^2 - R_0^2 \sin^2 \omega}$$

Figure 6-2. Schematic drawing for evaluation of the core transmission coefficients for the core secondary sources.

The δ_1 values are given by

$$\delta_1 = \frac{\int_{\omega=0}^{\pi} \exp \left[-\Sigma_{t_1} \times SN \right] \frac{\sin \omega}{2} d \omega}{\int_{\omega=0}^{\pi} \frac{\sin \omega}{2} d \omega} \quad (6-1)$$

where

$$SN = R_o \cos \omega + \sqrt{R_c^2 - R_o^2 \sin^2 \omega} \quad \text{and}$$

$$R_o = \frac{R_c}{\sqrt[3]{2}} .$$

In Table 6-5 are the transmission coefficients for all secondary sources.

TABLE 6-5. TRANSMISSION COEFFICIENTS FOR THE CORE (δ_1) AND REFLECTOR (γ_1) SECONDARY SOURCES OF THE GAS CORE REACTOR PROBLEM

	Four-Group Albedo Theory Method			
	i = 1	i = 2	i = 3	i = 4
γ_1	.98425	.97036	.94900	.78963
δ_1	.98811	.97759	.96134	.83643

- Radius of the spherical core = $R_c = 70$ cm.

- $R_o = R_c / \sqrt[3]{2}$

Dissection of P_o and P

As was done for the solid core reactor problem, it is also interesting to evaluate: (1) the fraction of neutrons which, having been produced by fission in the core, never traverse the core boundary, but are absorbed as group i , here denoted by A_{o_i} , and (2) the fraction of neutrons which are produced by fission in the core, spend time in the reflector but are eventually absorbed by nuclei of the core as group i , here denoted by A_i . The program ALB2 used for these numerical calculations is in Appendix G. In Table 6-6 are the probabilities, A_{o_i} and A_i , where

$$P_o = \sum_{i=1}^4 A_{o_i} \text{ and } P = \sum_{i=1}^4 A_i .$$

TABLE 6-6. NONLEAKAGE PROBABILITIES, A_{o_i} AND A_i , OF THE GAS CORE REACTOR PROBLEM USING THE FOUR-GROUP ALBEDO THEORY METHOD

	Four-Group Albedo Theory Method			
	$i = 1$	$i = 2$	$i = 3$	$i = 4$
A_{o_i}	.59772 E-03	.32561 E-03	.69758 E-06	.59687 E-09
A_i	.30723 E-03	.26766 E-02	.29157 E-01	.50246

- Radius of the spherical core = $R_c = 70$ cm.

- $R_o = R_c / \sqrt[3]{2}$

From this table, the total principal nonleakage probability, P_o , is found to be very small, .92404 E-03, while the total secondary nonleakage probability, P , is now very high, i.e., .53553. Thus, it is very clear that P almost represents 100% of the total nonleakage probability value for this gas core reactor problem.

Effective Neutron Multiplication Factor, k_{eff}

The neutrons produced by the $(n, 2n)$ interaction in the gas core are accounted for in k_{eff} evaluation. Thus, the k_{eff} value is given by

$$k_{eff} = \sum_{i=1}^4 (A_{o_i} + A_i) \frac{v_i \Sigma_{f_i} + 2 \Sigma_{n, 2n_i}}{\Sigma_{a_i}} \quad \text{and}$$

$$k_{eff} = 1.0752.$$

Comparison: XSDRNPM Versus Albedo Method Results

One-dimensional static neutronic analysis has also been performed on the gas core reactor problem with XSDRNPM, four-group, S_4 P_3 , computer code. It is desirable to compare the results obtained from this code because such comparisons aid in determining the accuracy and reliability of the obtained results. In Table 6-7 are the total nonleakage probability, P_{TNL} , and effective neutron multiplication factor, k_{eff} , for the system as obtained from XSDRNPM and the four-group albedo method.

TABLE 6-7. TOTAL NONLEAKAGE PROBABILITY, P_{TNL} , AND EFFECTIVE NEUTRON MULTIPLICATION FACTOR, k_{eff} , VALUES OF THE GAS CORE REACTOR PROBLEM USING THE XSDRNPM CODE AND THE ALBEDO METHOD

	XSDRNPM	Albedo
P_o	-	.92404 E-03
P	-	.53460
$P_{TNL} = P_o + P$.56869	.53553
k_{eff}	1.1374	1.0752
CPU Time	1.0 sec	< 1 sec

- Radius of the spherical core = $R_o = 70$ cm
- $R_o = R_c / \sqrt[3]{2}$
- Infinite reflector for the albedo results
- Reflector 100 cm thick for XSDRNPM values
- XSDRNPM parameters: four group, S_4 P_3 , number of spacial intervals = 60, overall convergence = $1.0 \text{ E-}04$
- XSDRNPM calculation on an IBM 3090-600J
- Four-group albedo method values on a 386-SX personal computer

Table 6-7 shows that, for the examined configuration, the albedo theory approximation is found to underpredict k_{eff} by about 5.5% relative to the S_4 P_3 result.

This rather large disagreement suggests that some of the approximations used in the albedo method deserve further attention. Sensitivity of results to some of these approximations are considered in the next section.

Analyses of Sensitivity

Comparative neutronic calculations are performed on the gas core reactor problem in one dimension to determine the impact of the variation of some parameters which are either

chosen intuitively in the method, or are of unconventional interest. The cases chosen are: (1) R_o variation, (2) assume $P_o = 0.0$, (3) $(n, 2n)$ reaction neglect in the reflector, (4) assume anisotropic secondary source for core transmittances, (5) assume anisotropic source for reflector albedos, and (6) combination of (4) and (5).

R_o Variation

It has been assumed that the core neutrons are uniformly produced on the surface of a sphere of radius R_o . R_o can range from zero to the core radius. In Table 6-8 are the nonleakage probability and effective neutron multiplication factor values for different R_o .

TABLE 6-8. NONLEAKAGE PROBABILITY AND EFFECTIVE NEUTRON MULTIPLICATION FACTOR VALUES OF THE GAS CORE REACTOR PROBLEM BY VARYING R_o

R_o	Four-Group Albedo Theory Method			
	P_o	P	$P_{TNL} = P_o + P$	k_{eff}
0.0	.12353 E-02	.53539	.53663	1.0774
R_o^*	.92404 E-03	.53460	.53553	1.0752
R_c	.60982 E-03	.53376	.53437	1.0730

- Radius of the spherical core = $R_c = 70$ cm.

- $R_o^* = R_c / \sqrt[3]{2} = .55559$ E+02 cm

Table 6-8 shows that less than 0.4% difference exists among k_{eff} values due to varying R_0 and the highest value of k_{eff} is for $R_0 = 0.0$. This is expected since the variation of R_0 yields effects on core transmittance which completely explain these results.

P_0 Neglect

From previous results it is appropriate to use the approximation that all fission neutrons escape from the gas core without core interaction. In Table 6-9 are the total nonleakage probability, P_{TNL} , and effective neutron multiplication factor, k_{eff} , using $P_0 = 0$.

TABLE 6-9. TOTAL NONLEAKAGE PROBABILITY, P_{TNL} , AND EFFECTIVE NEUTRON MULTIPLICATION FACTOR, k_{eff} , VALUES OF THE GAS CORE REACTOR PROBLEM NEGLECTING P_0 AND VARYING R_0

R_0	Four-Group Albedo Theory Method	
	$P_{TNL} = P_0 + P, P_0 = 0.0$	k_{eff}
0.0	.53653	1.0766
R_0^*	.53536	1.0747
R_c	.53433	1.0727

- Radius of the spherical core = $R_c = 70$ cm.
- $R_0^* = R_c / \sqrt[3]{2} = .55559 \text{ E}+02$ cm

Tables 6-8 and 6-9 show that the P_0 neglect gives smaller k_{eff} values but less than 0.07% difference exists among k_{eff} values considering the same imaginary R_0 . This is expected

since $P_0 = 0.0$ yields an increase in the fraction of neutrons that leaks from the core for the first time.

Be⁹(n, 2n) Be⁸ Reaction Neglect

The neutrons produced by the Be⁹ (n, 2n) 2 He⁴ interaction in the reflector are accounted for in ${}_1\beta_1$, ${}_1\beta_2$, ${}_1\beta_3$, and ${}_1\beta_4$ albedos, and result in k_{eff} variation. Because the ratio $\Sigma_{n, 2n}/\Sigma_{t1}$ (~1.6%) is small, it might seem appropriate to neglect the effect of the reaction. In Table 6-10 is compared the group 1 albedos and k_{eff} including the (n, 2n) reaction and neglecting it.

TABLE 6-10. GROUP 1 REFLECTOR ALBEDOS AND EFFECTIVE NEUTRON MULTIPLICATION FACTOR OF THE GAS CORE REACTOR PROBLEM FOR EITHER $\Sigma_{n, 2n} = 0$ OR $\Sigma_{n, 2n} \neq 0$

$\Sigma_{n, 2n}$	Four-Group Albedo Theory Method				
	${}_1\beta_{1'}$				k_{eff}
	$i' = 1$	$i' = 2$	$i' = 3$	$i' = 4$	
$= 0.0$.25992	.29529	.10932	.17031	1.0076
$\neq 0.0$.28040	.30778	.11624	.18519	1.0752

- Radius of the spherical core = $R_c = 70.0$ cm.

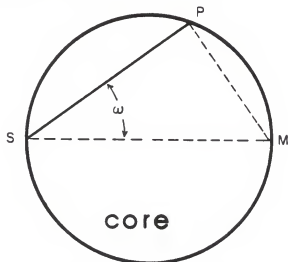
- $R_0 = R_c / \sqrt[3]{2}$

From Table 6-10, it is clear that neglect of the (n, 2n) interaction in the reflector yields for the probability that a neutron from group 1 be reflected by the reflector a value of .83, i.e., ${}_1\beta_1 + {}_1\beta_2 + {}_1\beta_3 + {}_1\beta_4 = .83$. Including the reaction yields a value of nearly .89. Finally, the Be⁹ (n, 2n) Be⁸ interaction neglect leads an error of about 6.3%

in k_{eff} value for the examined configuration. The $(n, 2n)$ reaction is certainly significant.

Anisotropic Secondary Source on the Reflector-Core Interface for Core Transmittance Calculation

Previously, the core transmission coefficients, γ_i , were evaluated by considering an isotropic neutron source on a spherical core of diameter $2R$. The angular neutron flux distribution deep inside a homogeneous medium is nearly isotropic because leakage processes, which are anisotropic by nature, are relatively unimportant and neutron production processes, such as fission, tend to be isotropic. However, near strong discontinuities in material properties, such as the core-reflector interface, the angular neutron flux distribution is usually relatively anisotropic. Consider a linear anisotropic neutron source emitting S_R neutrons per second (see Figure 6-3).



$$S_R = S_0 (1 + A \mu) \text{ where}$$

$$\mu = \cos \omega, \quad 0 \leq \mu \leq 1.$$

and the parameter A represents an angular coefficient.

$$SP = 2R \cos \omega$$

$$SM = 2R$$

Figure 6-3. Schematic drawing for evaluation of the core transmission coefficients, γ_i , due to a linear anisotropic neutron source on a spherical core.

The core transmission coefficients, γ_i , for this linear anisotropic neutron source are given by

$$\gamma_i = \frac{\int_{\omega=0}^{\pi/2} \exp \left[-\Sigma_{t1} \times 2R \times \cos \omega \right] (1 + A \cos \omega) 2\pi \sin \omega d\omega}{\int_{\omega=0}^{\pi/2} (1 + A \cos \omega) 2\pi \sin \omega d\omega} \quad (6-2)$$

In Table 6-11 are compared the core transmission coefficients, γ_i , and effective neutron multiplication factor for different values of the angular anisotropy coefficient, A .

TABLE 6-11. CORE TRANSMISSION COEFFICIENTS, γ_i , AND EFFECTIVE NEUTRON MULTIPLICATION FACTOR OF THE GAS CORE REACTOR PROBLEM BY VARYING THE ANGULAR ANISOTROPY COEFFICIENT, A .

A	Four-Group Albedo Theory Method				
	γ_i				k_{eff}
	$i = 1$	$i = 2$	$i = 3$	$i = 4$	
0.0	.98424	.97037	.94900	.78963	1.0752
0.5	.98321	.96841	.94566	.77671	1.0922
1.0	.98251	.96710	.94344	.76811	1.1028

Table 6-11 shows that when the angular anisotropy coefficient, A , increases, the core transmission coefficients, γ_i , decrease, which clearly yields an increase in the effective neutron multiplication factor.

Anisotropic Source on the Reflector-Core Interface
for Reflector Albedo Calculation

Directional dependence of the neutron flux is greatest close to a boundary. To this point, the reflector albedo values were evaluated by using the diffusion approximation (P_1) to the Boltzmann transport equation. Consider the reflector albedo evaluation employing the next relevant higher-order approximation (P_3) to the Boltzmann transport equation. To simplify the calculation, the following relationship among the P_1 and P_3 reflector albedo values is assumed:

$$\beta_{P_3}^{\text{spherical}} = \beta_{P_1}^{\text{spherical}} \times \frac{\beta_{P_3}^{\text{plane symmetry}} \text{ (for } B \neq 0 \text{)}}{\beta_{P_3}^{\text{plane symmetry}} \text{ (for } B = 0 \text{)}} . \quad (6-3)$$

Consider the problem of a linear anisotropic neutron source emitting S_0 neutrons per second that illuminates a plane surface of a homogeneous half-space as shown in Figure 6-4.

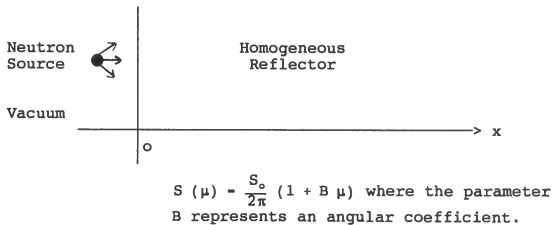


Figure 6-4. Schematic drawing for evaluation of the reflector albedo values due to a linear anisotropic neutron source in P_3 plane-symmetry geometry.

Plane-symmetry, P_3 , four-group transport
theory reflector albedo evaluation

Because of the lengthy algebraic manipulations involved, only an outline of the method of determining β_i value is presented here. Results for all the P_3 albedos are presented in Appendix G (see program ALB3). The angular neutron flux expression for group i is

$$\phi_i(x, \mu) = \sum_{n=0}^3 \frac{2n+1}{4\pi} \phi_{i,n}(x) P_n(\mu) \quad (6-4)$$

which yields a coupled set of sixteen linear, first-order ordinary differential equations:

(1) For $i = 1$

$$\begin{aligned} \frac{d}{dx} \phi_{11} + \Sigma_{R1} \phi_{10} &= 0 \\ \frac{d}{dx} \phi_{10} + 2 \frac{d}{dx} \phi_{12} + 3 \Sigma_{t1} \phi_{11} &= 0 \\ 2 \frac{d}{dx} \phi_{11} + 3 \frac{d}{dx} \phi_{13} + 5 \Sigma_{t1} \phi_{12} &= 0 \\ 3 \frac{d}{dx} \phi_{12} + 7 \Sigma_{t1} \phi_{13} &= 0 \end{aligned} \quad (6-5)$$

(2) For $i = 2$

$$\begin{aligned} \frac{d}{dx} \phi_{21} + \Sigma_{R2} \phi_{20} - \Sigma_{s12} \phi_{10} &= 0 \\ \frac{d}{dx} \phi_{20} + 2 \frac{d}{dx} \phi_{22} + 3 \Sigma_{t2} \phi_{21} &= 0 \\ 2 \frac{d}{dx} \phi_{21} + 3 \frac{d}{dx} \phi_{23} + 5 \Sigma_{t2} \phi_{22} &= 0 \\ 3 \frac{d}{dx} \phi_{22} + 7 \Sigma_{t2} \phi_{23} &= 0 \end{aligned} \quad (6-6)$$

(3) For $i = 3$

$$\begin{aligned}
 \frac{d}{dx} \phi_{31} + \Sigma_{R_3} \phi_{30} - \Sigma_{B_{13}} \phi_{10} + \Sigma_{B_{23}} \phi_{20} \\
 \frac{d}{dx} \phi_{30} + 2 \frac{d}{dx} \phi_{32} + 3 \Sigma_{t_3} \phi_{31} - 0 \\
 2 \frac{d}{dx} \phi_{31} + 3 \frac{d}{dx} \phi_{33} + 5 \Sigma_{t_3} \phi_{32} - 0 \\
 3 \frac{d}{dx} \phi_{32} + 7 \Sigma_{t_3} \phi_{33} - 0
 \end{aligned} \tag{6-7}$$

(4) For $i = 4$

$$\begin{aligned}
 \frac{d}{dx} \phi_{41} + \Sigma_{R_4} \phi_{40} - \Sigma_{B_{14}} \phi_{10} + \Sigma_{B_{24}} \phi_{20} + \Sigma_{B_{34}} \phi_{30} \\
 \frac{d}{dx} \phi_{40} + 2 \frac{d}{dx} \phi_{42} + 3 \Sigma_{t_4} \phi_{41} - 0 \\
 2 \frac{d}{dx} \phi_{41} + 3 \frac{d}{dx} \phi_{43} + 5 \Sigma_{t_4} \phi_{42} - 0 \\
 3 \frac{d}{dx} \phi_{42} + 7 \Sigma_{t_4} \phi_{43} - 0
 \end{aligned} \tag{6-8}$$

The expression for ${}_i\beta_4$ is

$${}_i\beta_4 = \frac{2\pi \int_{\mu=0}^{-1} \mu \phi_4(x=0, \mu) d\mu}{2\pi \int_{\mu=0}^{+1} \frac{S_0}{2\pi} (1 + B\mu) d\mu} . \tag{6-9}$$

All boundary conditions and albedo expressions are presented in program ALB3 (see Appendix G).

In Table 6-12 are compared the total nonleakage probability and effective neutron multiplication factor for different values of the angular coefficient, B.

TABLE 6-12. TOTAL NONLEAKAGE PROBABILITY, P_{TNL} , AND EFFECTIVE NEUTRON MULTIPLICATION FACTOR OF THE GAS CORE REACTOR PROBLEM BY VARYING THE ANGULAR COEFFICIENT, B

B	Four-Group Albedo Theory Method			
	P_0	P	P_{TNL}	k_{eff}
0.0	.92404 E-03	.53460	.53553	1.0752
0.5	.92404 E-03	.53548	.53640	1.0773
1.0	.92404 E-03	.53601	.53694	1.0786

- Radius of the spherical core = $R_c = 70.0$ cm

- $R_0 = R_c / \sqrt[3]{2}$

Table 6-12 shows that when the angular coefficient, B, increases, the effective neutron multiplication factor of the examined configuration increases. The P_3 reflector albedo values (in plane symmetry) are presented in Table 6-13. This table shows that when the angular coefficient, B, increases, ${}_1\beta_1$, ${}_2\beta_2$, ${}_3\beta_3$, ${}_4\beta_4$, and ${}_1\beta_2$ decrease, but ${}_1\beta_3$, ${}_1\beta_4$, ${}_2\beta_3$, ${}_2\beta_4$, and ${}_3\beta_4$ increase (outweighing these contributions for the k_{eff} value).

TABLE 6-13. P_3 REFLECTOR ALBEDO VALUES IN PLANE-SYMMETRY GEOMETRY BY VARYING THE ANGULAR ANISOTROPY COEFFICIENT, B.

${}_1\beta_1$	B		
	0.0	0.5	1.0
${}_2\beta_2$.40604	.40243	.40002
${}_1\beta_2$.30134	.29953	.29832
${}_1\beta_3$.10219	.10323	.10393
${}_1\beta_4$.14748	.15038	.15230
${}_2\beta_2$.68207	.67794	.67519
${}_1\beta_3$.15038	.15030	.15189
${}_2\beta_4$.12007	.12232	.12381
${}_3\beta_3$.73584	.73205	.72952
${}_3\beta_4$.21694	.21973	.22160
${}_4\beta_4$.95934	.95852	.95797

Anisotropic Neutron Source on the Reflector-Core Interface for Reflector Albedo and Core Transmittance Evaluations

The anisotropy of the neutron flux is greatest close to the reflector-core interface of the gas core reactor problem examined. In Table 6-14 is summarized the results of calculations of k_{eff} by XSDRNPM and the four-group albedo approach including core and reflector anisotropy factors $A = B = 1.5$.

TABLE 6-14. COMPARISON OF k_{eff} RESULTS OF THE GAS CORE REACTOR PROBLEM USING VARIOUS CORE DENSITY FACTOR (D.F.) FROM THE XSDRNPM CODE AND THE FOUR-GROUP ALBEDO METHOD WITH ANISOTROPY FACTORS $A = B = 1.5$

D.F.	P_{TNL}		k_{eff}		k_{eff}
	XSDRNPM	Albedo	XSDRNPM	Albedo	Relative Error
0.5	.47460	.47088	.95712	.95154	.583 E-02
1.0	.56869	.56469	1.1374	1.1329	.396 E-02
2.0	.63460	.63470	1.2511	1.2576	.520 E-02

- Radius of the spherical core = $R_c = 70.0$ cm
- $R_o = R_c / \sqrt[3]{2}$
- Angular coefficients: $A = 1.5$ and $B = 1.5$
- XSDRNPM parameters: Four-group, S_4 P_3 , number of spatial intervals = 60, overall convergence = $1.0 \text{ E-}04$.
- k_{eff} relative error = $\frac{k_{eff}(\text{XSDRNPM}) - k_{eff}(\text{albedo})}{k_{eff}(\text{XSDRNPM})}$

From Table 6-14, it is clear that, if $A = B = 1.5$ is assumed, the k_{eff} values for density factors in the range 0.5 to 2.0 as calculated by the albedo method are in close agreement with those calculated by XSDRNPM. In contrast, the results presented in Table 6-7 for the case D.F. = 1.0 and $A = B = 0.0$ (i.e., for isotropic-incident albedos) indicate significantly less correspondence. It is concluded that the anisotropic-incident effects are crucial for k_{eff} calculations. However, this effect should be far less important in reactivity (i.e., change in k_{eff} due to parameter variation) determination.

It should be emphasized, again, that the ability to recognize the physical phenomena origin of k_{eff} variation through albedo parameter dependence is the most interesting aspect of the multigroup albedo approach. In the next chapter, another aspect of the value of these methods is addressed, viz., the resulting ability to approximate effects on k_{eff} of symmetry perturbations which could not be included in conventional approaches (e.g., XSDRNPM application) without incurring considerable computational time and expense.

CHAPTER VII
STATIC, ONE-DIMENSIONAL NEUTRONIC ANALYSIS OF A SMALL
OPTICAL-PATH-LENGTH CORE REACTOR
(THREE-REGION GAS CORE REACTOR)

Introduction

As stated in Chapter VI the ideas employed in the multigroup neutron albedo approach are more useful and applicable to gaseous core reactors than to the "solid core" applications. Figure 7-1 illustrates a gaseous core (fuel radius R_c), intermediate gaseous region (thickness $R_p - R_c$), and solid moderator (thickness $R_m - R_p$) of a reactor system similar to a proposed concept of a gas core propulsion reactor [12]. For such a system, crucial considerations of nuclear reactor system dynamics are not approachable (at reasonable cost) with conventional calculational procedures.

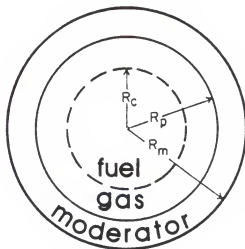


Figure 7-1. Schematic drawing of a three-region gas core reactor.

Three-Region Gas Core Reactor Atomic Densities

Consider a spherical cavity region consisting of a ball of gaseous fuel suspended in an intermediate gas (hydrogen). The cavity is surrounded by an infinite beryllium oxide reflector-moderator. In Table 7-1 are the plausible atomic densities for such a three-region gas core reactor problem.

TABLE 7-1. ATOMIC DENSITIES OF THE THREE-REGION GAS CORE REACTOR PROBLEM

Region	Material	Atomic Density (atoms/(barn x cm))
Fuel	Uranium - 238	.10400 E-04
	Uranium - 238	.18350 E-05
	Fluorine - 19	.49830 E-04
Intermediate Gas	Hydrogen - 1	.49830 E-04
Moderator	Be in BeO	.7223 E-01
Reflector	Oxygen - 16	.7223 E-01

Three-Region Gas Core Reactor Group Constants

The fuel and reflector group constants are obtained from XSDRNP calculations (see Tables 6-2 and 6-3, respectively). The fast group constants for hydrogen are obtained from PHROG [13] calculations while the thermal group constants are obtained from BRT [14] calculations by using the free gas kernel option. The molecular binding of the hydrogen molecule is ignored (i.e., atomic hydrogen is assumed).

In Table 7-2 are the hydrogen group constants.

TABLE 7-2. HYDROGEN GROUP CONSTANTS OF THE THREE-REGION GAS CORE REACTOR PROBLEM

	Four Energy-Group Structure			
	i = 1	i = 2	i = 1	i = 4
Σ_a	.00000	.00000	.36286 E-06	.69789 E-05
Σ_{tr}	.76168 E-04	.12559 E-03	.30068 E-03	.47675 E-03
Σ_t	.13829 E-03	.54005 E-03	.98166 E-03	.12778 E-02
i-i'	$\Sigma_{s_{ii'}}$			
	i' = 1	i' = 2	i' = 3	i' = 4
4	.67670 E-04	.70140 E-04	.47548 E-06	.15990 E-09
2	-	.43390 E-03	.10612 E-03	.35687 E-07
3	-	-	.85897 E-03	.12234 E-03
4	-	-	-	.12709 E-02

- Σ in cm^{-1}
- Fuel radius = $R_c = 100.0$ cm.
- Hydrogen thickness = $R_p - R_c = 50.0$ cm.

From Table 7-2, the ratio of hydrogen thickness to the total mean free paths in the hydrogen region is found to be very small (~ 0.06 for thermal neutrons). Thus, this hydrogen region is also of small optical dimensions.

Reflector (β) Albedos

The diffusion approximation within the infinite beryllium oxide reflector is employed. The neutrons produced by the $\text{Be}^9 (n, 2n) \text{Be}^8$ interaction are accounted for in ${}_1\beta_1$, ${}_1\beta_2$, ${}_1\beta_3$, and ${}_1\beta_4$ values. In Table 7-3 are the reflector albedo values obtained.

TABLE 7-3. REFLECTOR ALBEDOS OF THE THREE-REGION GAS CORE REACTOR PROBLEM USING THE FOUR-GROUP ALBEDO THEORY

$i-i'$	${}_1\beta_{i'}$			
	$i' = 1$	$i' = 2$	$i' = 3$	$i' = 4$
1	.29516	.31344	.11856	.18865
2	-	.63572	.16100	.13658
3	-	-	.69559	.23881
4	-	-	-	.94340

- Internal radius of the spherical reflector = $R_m = 150.0$ cm.
- Infinite beryllium oxide reflector.

From Table 7-3, the probability that a neutron from group 1 be reflected by the reflector is found to be very large (${}_1\beta_1 + {}_1\beta_2 + {}_1\beta_3 + {}_1\beta_4 \approx .92$). Thus, 8% of the incident neutrons from group 1 are absorbed by the reflector.

The Ping-Pong Decision Process

Assumptions established to facilitate the neutronic analyses are as follows: (1) the fuel is uniformly distributed within the central portion of the cavity volume and (2) no mixing is assumed to occur between fuel and hydrogen in the cavity. In the albedo method applied to the three-region gas core reactor problem, the description of average neutron histories is accomplished through the specification of reflector reflection and fuel and hydrogen transmission probabilities at each neutron traversal of the fuel-hydrogen-reflector interfaces.

Strategy for the Three-Region Gas Core Reactor Problem

As previously presumed, χ_3 and χ_4 are considered negligible, i.e., all fission neutrons are emitted in groups 1 and 2. Figure 7-2 represents the average neutron history of a group 1 neutron source in the three-region gas core reactor problem. Because of the lengthy algebraic manipulations involved, only a sketch of the method of determining the nonleakage probabilities for the fuel and hydrogen regions is presented here.

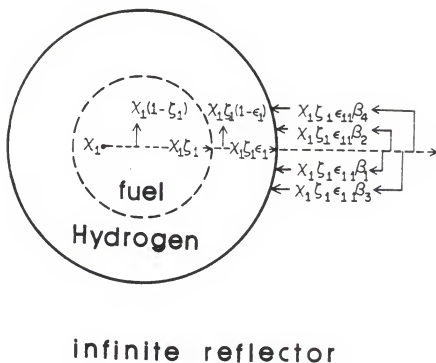


Figure 7-2. Schematic drawing showing the albedo method strategy for the three-region gas core reactor problem.

With reference to Figure 7-2, the first source of initial leakage neutrons for the evaluation of P_0 and P is χ_1 . The fraction of these neutrons transmitted through the fuel region as group 1 is ζ_1 . After that the fraction of these neutrons transmitted through the hydrogen region as group 1 is e_1 . The quantities returned as group 1, 2, 3, and 4 on the first cavity boundary traverse are $\chi_1 \zeta_1 e_1 {}_1\beta_1$, $\chi_1 \zeta_1 e_1 {}_1\beta_2$, $\chi_1 \zeta_1 e_1 {}_1\beta_3$, and $\chi_1 \zeta_1 e_1 {}_1\beta_4$, respectively. The neutron fraction, $\chi_1(1 - \zeta_1)$ interacts in the fuel region. It can be dissected into the contributions:

- (1) The fraction absorbed in the fuel as group 1, i.e.,

$$\chi_1(1 - \zeta_1)(\Sigma_{a_1}^F / \Sigma_{t_1}^F) ,$$

- (2) The fraction scattered in the fuel as group 1, i.e.,

$$\chi_1(1 - \zeta_1)(\Sigma_{s_{11}}^F / \Sigma_{t_1}^F) = F_1 ,$$

- (3) The fraction scattered in the fuel as group 2, i.e.,

$$\chi_1(1 - \zeta_1)(\Sigma_{s_{12}}^F / \Sigma_{t_1}^F) = F_2 ,$$

- (4) The fraction scattered in the fuel as group 3, i.e.,

$$\chi_1(1 - \zeta_1)(\Sigma_{s_{13}}^F / \Sigma_{t_1}^F) = F_3 , \text{ and}$$

- (5) The fraction scattered in the fuel as group 4, i.e.,

$$\chi_1(1 - \zeta_1)(\Sigma_{s_{14}}^F / \Sigma_{t_1}^F) = F_4 .$$

The neutron fraction $\chi_1 \zeta_1 (1 - e_1)$ interacts in the hydrogen region. It can be dissected into the contributions:

- (1) The fraction absorbed by hydrogen as group 1, i.e.,

$$\chi_1 \zeta_1 (1 - e_1) (\Sigma_{a_1}^H / \Sigma_{t_1}^H) ,$$

- (2) The fraction scattered in the hydrogen region as group 1, i.e.,

$$\chi_1 \zeta_1 (1 - e_1) (\Sigma_{s_{11}}^H / \Sigma_{t_1}^H) = E_1 ,$$

- (3) The fraction scattered in the hydrogen region as group 2, i.e.,

$$\chi_1 \zeta_1 (1 - e_1) (\Sigma_{s_{12}}^H / \Sigma_{t_1}^H) = E_2 ,$$

- (4) The fraction scattered in the hydrogen region as group 3, i.e.,

$$\chi_1 \zeta_1 (1 - e_1) (\Sigma_{s_{13}}^H / \Sigma_{t_1}^H) = E_3 , \text{ and}$$

- (5) The fraction scattered in the hydrogen region as group 4, i.e.,

$$\chi_1 \zeta_1 (1 - e_1) (\Sigma_{s_{14}}^H / \Sigma_{t_1}^H) = E_4 .$$

Thus, in the first cavity traversal, there are created four fuel secondary neutron sources whose strengths are F_1, F_2, F_3 , and F_4 , four hydrogen secondary neutron sources whose strengths are E_1, E_2, E_3 , and E_4 , and four reflector secondary neutron sources. Moreover, $\chi_1(1 - \zeta_1)(\Sigma_{a_1}^F / \Sigma_{t_1}^F)$ is the fraction absorbed in the fuel region as group 1 while $\chi_1 \zeta_1(1 - e_1)(\Sigma_{a_1}^H / \Sigma_{t_1}^H)$ is the fraction absorbed by hydrogen nuclei as group 1.

For the next steps, it is required to consider the presence of all secondary sources (i.e., 12 sources) in

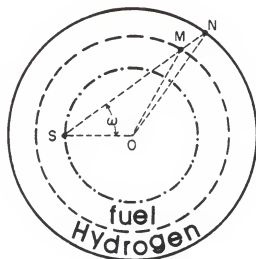
order to estimate the total fractions that are absorbed as group 1, 2, 3, and 4 in the fuel and hydrogen regions.

Transmission Coefficients for the Fuel, Hydrogen, and Reflector Secondary Neutron Sources

The fuel-hydrogen-reflector neutron transfer is described by means of three point-of-view neutron source transmission coefficients and are listed below.

First point of view: Fuel secondary neutron sources

The fuel secondary neutron sources can "see" the fuel region and the hydrogen region. Two transmission coefficients are of interest and have been identified as ζ_1 and e_1 . With the aid of Figure 7-3 these transmission coefficient evaluations can be better understood.



$$OS = R_o ; OM = R_c$$

$$ON = R_p$$

Thus,

$$SM = R_o \cos \omega + \sqrt{R_c^2 - R_o^2 \sin^2 \omega}$$

$$SN = R_o \cos \omega + \sqrt{R_p^2 - R_o^2 \sin^2 \omega}$$

$$MN = \sqrt{R_p^2 - R_o^2 \sin^2 \omega} - \sqrt{R_c^2 - R_o^2 \sin^2 \omega}$$

Figure 7-3. Schematic drawing for the transmission coefficient evaluation under fuel secondary neutron source point of view for the three-region gas core reactor problem.

In calculating the transmission coefficients for each energy group it is assumed that: (1) the neutron sources are isotropic and (2) the neutrons are homogeneously distributed at the surface of a sphere of radius R_0 given by

$R_0 = [\frac{R_c^3}{2}]^{1/3}$, i.e., this imaginary sphere has half volume of the fuel region volume. The ζ_1 are given by

$$\zeta_1 = \frac{\int_{\omega=0}^{\pi} \exp[-\Sigma_{t_1}^g \times SM] \frac{\sin \omega}{2} d\omega}{\int_{\omega=0}^{\pi} \frac{\sin \omega}{2} d\omega} \quad (7-1)$$

where

$$SM = R_0 \cos \omega + \sqrt{R_c^2 - R_0^2 \sin^2 \omega} \quad \text{and} \quad R_0 = R_c / \sqrt[3]{2}.$$

The ϵ_1 are given by

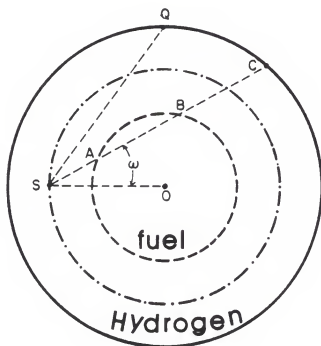
$$\epsilon_1 = \frac{\int_{\omega=0}^{\pi} \exp[-\Sigma_{t_1}^H \times MN] \frac{\sin \omega}{2} d\omega}{\int_{\omega=0}^{\pi} \frac{\sin \omega}{2} d\omega} \quad (7-2)$$

where

$$MN = \sqrt{R_p^2 - R_0^2 \sin^2 \omega} - \sqrt{R_c^2 - R_0^2 \sin^2 \omega}$$

Second point of view: Hydrogen
secondary neutron sources

For the hydrogen neutron sources there are four types of relevant transmission coefficients. They involve: (1) when a neutron crosses the hydrogen-fuel-hydrogen region, and (2) when a neutron crosses only the hydrogen region. With the aid of Figure 7-4, these transmission coefficient evaluations can be visualized and derived.



$$OB = R_c ; OC = R_p$$

$$OS = R_H$$

Thus,

$$SQ = R_H \cos \omega + \sqrt{R_p^2 - R_H^2 \sin^2 \omega}$$

$$SA = R_H \cos \omega - \sqrt{R_c^2 - R_H^2 \sin^2 \omega}$$

$$AB = 2 \sqrt{R_c^2 - R_H^2 \sin^2 \omega} \quad \text{and}$$

$$BC = \sqrt{R_p^2 - R_H^2 \sin^2 \omega} - \sqrt{R_c^2 - R_H^2 \sin^2 \omega}$$

$$\text{For } 0 \leq \omega \leq \pi$$

Figure 7-4. Schematic drawing for the transmission coefficient evaluation under hydrogen secondary neutron source point of view for the three-region gas core reactor problem.

In calculating the transmission coefficients for each energy group it is assumed that: (1) the neutron sources are isotropic and (2) the neutrons are homogeneously distributed at the surface of a sphere of radius R_H given by

$R_H = \left[\frac{R_p^3 + R_c^3}{2} \right]^{1/3}$, i.e., this imaginary sphere has the average volume of the fuel region and cavity region volumes. The relevant transmission coefficient values are given by

(1) SA hydrogen path

$$\xi_{SA,1} = \frac{\int_{\omega=0}^{\omega_0} \exp \left[-\Sigma_{t1}^H \times SA \right] \frac{\sin \omega}{2} d\omega}{\int_{\omega=0}^{\omega_0} \frac{\sin \omega}{2} d\omega} \quad (7-3)$$

where

$$SA = R_H \cos \omega - \sqrt{R_c^2 - R_H^2 \sin^2 \omega}, \quad R_H = \left[\frac{R_p^3 + R_c^3}{2} \right]^{1/3},$$

$$\text{and } \omega_0 = \sin^{-1} \left[\frac{R_c}{R_H} \right]$$

(2) AB fuel path

$$\xi_{AB,1} = \frac{\int_{\omega=0}^{\omega_0} \exp \left[-\Sigma_{t1}^F \times AB \right] \frac{\sin \omega}{2} d\omega}{\int_{\omega=0}^{\omega_0} \frac{\sin \omega}{2} d\omega} \quad (7-4)$$

where

$$AB = 2 \sqrt{R_c^2 - R_H^2 \sin^2 \omega}$$

(3) BC hydrogen path

$$\xi_{BC,1} = \frac{\int_{\omega=0}^{\omega_0} \exp \left[-\Sigma_{t,1}^H \times BC \right] \frac{\sin \omega}{2} d\omega}{\int_{\omega=0}^{\omega_0} \frac{\sin \omega}{2} d\omega} \quad (7-5)$$

where

$$BC = \sqrt{R_p^2 - R_H^2 \sin^2 \omega} - \sqrt{R_c^2 - R_H^2 \sin^2 \omega}$$

(4) SQ hydrogen path

$$\xi_{SQ,1} = \frac{\int_{\omega=\omega_0}^{\pi} \exp \left[-\Sigma_{t,1}^H \times SQ \right] \frac{\sin \omega}{2} d\omega}{\int_{\omega=\omega_0}^{\pi} \frac{\sin \omega}{2} d\omega} \quad (7-6)$$

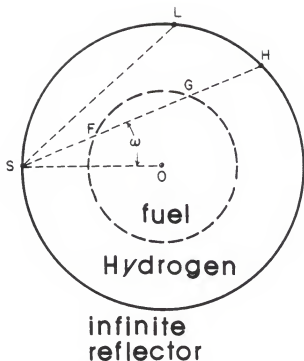
where

$$SQ = R_H \cos \omega + \sqrt{R_p^2 - R_H^2 \sin^2 \omega} .$$

Third point of view: Reflector secondary neutron sources

For the reflector neutron sources there are four types of relevant transmission coefficients. They involve:

(1) when a neutron crosses the hydrogen-fuel-hydrogen region, and (2) when a neutron crosses only the hydrogen region. With the aid of Figure 7-5, these transmission coefficient evaluations can be comprehended.



$$OS = R_p, OG = R_c$$

Thus,

$$SL = 2 R_p \cos \omega$$

$$SF = R_p \cos \omega - \sqrt{R_c^2 - R_p^2 \sin^2 \omega}$$

$$GH = SF$$

$$FG = 2 \sqrt{R_c^2 - R_p^2 \sin^2 \omega}$$

$$\text{for } 0 \leq \omega \leq \pi/2$$

Figure 7-5. Schematic drawing for the transmission coefficient evaluation under reflector secondary neutron source point of view for the three-region gas core reactor problem.

The relevant transmission coefficients are given by

(1) SF hydrogen path

$$\eta_{SF,1} = \frac{\int_{\omega=0}^{\omega_c} \exp \left[- \Sigma_{t,1}^H \times SF \right] \frac{\sin \omega}{2} d \omega}{\int_{\omega=0}^{\omega_c} \frac{\sin \omega}{2} d \omega} \quad (7-7)$$

where

$$SF = R_p \cos \omega - \sqrt{R_c^2 - R_p^2 \sin^2 \omega} \text{ and } \omega_c = \sin^{-1} \left[\frac{R_c}{R_p} \right]$$

(2) FG fuel path

$$\eta_{FG,1} = \frac{\int_{\omega=0}^{\omega_L} \exp \left[-\Sigma_{t,1}^F \times FG \right] \frac{\sin \omega}{2} d\omega}{\int_{\omega=0}^{\omega_L} \frac{\sin \omega}{2} d\omega} \quad (7-8)$$

where

$$FG = 2 \sqrt{R_c^2 - R_p^2 \sin^2 \omega}$$

(3) GH hydrogen path

$$\eta_{GH,1} = \frac{\int_{\omega=0}^{\omega_L} \exp \left[-\Sigma_{t,1}^H \times GH \right] \frac{\sin \omega}{2} d\omega}{\int_{\omega=0}^{\omega_L} \frac{\sin \omega}{2} d\omega} \quad (7-9)$$

where

$$GH = R_p \cos \omega - \sqrt{R_c^2 - R_p^2 \sin^2 \omega}$$

(4) SL hydrogen path

$$\eta_{SL,1} = \frac{\int_{\omega=\omega_L}^{\pi/2} \exp \left[-\Sigma_{t,1}^S \times SL \right] \frac{\sin \omega}{2} d\omega}{\int_{\omega=\omega_L}^{\pi/2} \frac{\sin \omega}{2} d\omega} \quad (7-10)$$

where

$$SL = 2 R_p \cos \omega$$

In Table 7-4 are the transmission coefficient values for the fuel, hydrogen, and reflector secondary neutron sources.

TABLE 7-4. TRANSMISSION COEFFICIENTS FOR ALL SECONDARY NEUTRON SOURCES OF THE THREE-REGION GAS CORE REACTOR PROBLEM

	Four-Group Albedo Theory Method			
	i = 2	i = 2	i = 3	i = 4
ζ_i	.98308	.96821	.94542	.77782
ϵ_i	.99913	.96815	.94288	.92631
$\xi_{EX,i}$.99411	.97722	.99907	.94699
$\xi_{AB,i}$.97156	.94675	.90913	.64883
$\xi_{HF,i}$.99087	.96457	.93658	.91828
$\xi_{HD,i}$.96486	.96109	.93092	.91147
$\eta_{EF,i}$.99087	.96109	.93092	.91892
$\eta_{FD,i}$.97110	.94590	.90771	.64408
$\eta_{DH,i}$.99087	.96486	.93092	.91147
$\eta_{EL,i}$.98470	.94198	.89786	.86982

- Fuel radius = $R_c = 100.0$ cm
- Hydrogen thickness = $R_p - R_c = 50.0$ cm
- Infinite beryllium oxide reflector

$$- R_o = R_c / \sqrt[3]{2}$$

$$- R_H = \left[\frac{R_p^3 + R_c^3}{2} \right]^{1/3}$$

Dissection of P_o and P

As was done for the two-region gas core reactor problem, it is of interest to evaluate: (1) the fraction of neutrons which, having been produced by fission in the fuel region, never traverse the fuel-hydrogen interface, and are absorbed as group i ; here denoted by $A_{o_i}^F$, (2) the fraction of neutrons which are produced by fission in the fuel region, spend time out of the fuel region, but are eventually absorbed by nuclei of the fuel as group i ; here denoted by A_i^F , and (3) the fraction of neutrons which are produced by fission in the fuel region and are absorbed by hydrogen nuclei as group i ; here denoted by A_i^H . In Table 7-5 are the nonleakage probabilities $A_{o_i}^F$, A_i^F , and A_i^H where

$$P_o^F = \sum_{i=1}^4 A_{o_i}^F, \quad P^F = \sum_{i=1}^4 A_i^F \text{ and } P_{TNL}^F = P_o^F + P^F.$$

TABLE 7-5. NONLEAKAGE PROBABILITY RESULTS OF THE THREE-REGION GAS CORE REACTOR PROBLEM USING THE FOUR-GROUP ALBEDO THEORY METHOD

	Four-Group Albedo Theory Method			
	$i = 1$	$i = 2$	$i = 3$	$i = 4$
$A_{o_1}^F$.84746 E-03	.46700 E-03	.14235 E-05	.16362 E-08
A_1^F	.15303 E-03	.10442 E-02	.15243 E-01	.45393
A_1^H	.00000	.00000	.51114 E-04	.38556 E-02

- Fuel region radius = $R_c = 100.0$ cm
- Hydrogen region thickness = $R_p - R_c = 50.0$ cm
- Infinite beryllium oxide reflector

$$- R_o = \frac{R_c}{\sqrt[3]{2}} \text{ and } R_H = \left[\frac{R_p^3 + R_c^3}{2} \right]^{1/3}$$

From Table 7-5, the fuel total principal nonleakage probability, P_o^p , is found to be very small, .17815 E-02, while the fuel total secondary nonleakage probability, P^p , is very high, .46991. It is clear that P^p represents about 100% of the fuel total nonleakage probability value for this three-region gas core reactor problem. The hydrogen total nonleakage probability, P_{TNL}^H , value is found to be very small, .39067 E-02. The hydrogen total nonleakage probability is given by

$$P_{TNL}^H = \sum_{i=1}^4 A_i^H .$$

Effective Neutron Multiplication Factor, k_{eff}

The neutrons produced by the (n, 2n) interaction in the fuel region are accounted for in k_{eff} value, i.e.,

$$k_{eff} = \sum_{i=1}^4 (A_{o_i}^p + A_i^p) \frac{v_i \Sigma_{f_i} + 2 \Sigma_{n,2n_i}}{\Sigma_{a_i}} .$$

Using the numerical values for this problem, $k_{eff} = .95437$. The program ALB4 used for these numerical calculations is in Appendix G.

Analyses of Sensitivity

Because of its location between the fuel region and the reflector, the hydrogen in the reactor cavity region notably affects neutron multiplication due to both neutron absorption and scattering collisions. Comparative neutronic calculations are performed on the three-region gas core reactor problem in one dimension to determine the impact of the variation of the hydrogen density factor.

Hydrogen Limiting Density Factor

At constant fuel mass and radius, in order to characterize the reference reactor configuration the effective neutron multiplication factor is calculated by employing hydrogen at different density factors (D.F.) in the cavity of the reactor. The hydrogen limiting density factor (D.F.*) corresponds to a maximum in k_{eff} . The variation of k_{eff} is shown in Table 7-6.

TABLE 7-6. EFFECTS OF THE HYDROGEN DENSITY FACTOR VARIATION ON THE THREE-REGION GAS CORE REACTOR PROBLEM PROPERTIES USING THE FOUR-GROUP ALBEDO METHOD

Hydrogen D.F.	Four-Group Albedo Theory Method		
	k_{eff}	P_{TNL}^H	P_{TNL}^F
1.0	.95437	.39067 E-02	.48108
2.0	.95951	.78195 E-02	.47417
8.5	.96381	.11753 E-01	.48020
1.0	.96713	.15719 E-01	.47791
8.5	.97007	.19730 E-01	.47921
8.5	.97211	.23795 E-01	.48016
8.5	.97348	.27926 E-01	.48078
8.5	.97421	.32130 E-01	.48109
2.0	.97433	.33834 E-01	.48111
8.5	.97435	.34263 E-01	.48111
8.5	.97435	.34692 E-01	.48110
8.7*	.97435	.35984 E-01	.48110
8.5	.97435	.35552 E-01	.48109
8.9	.97434	.35984 E-01	.48108
1.0	.97433	.36416 E-01	.48107
11.0	.97386	.19730 E-01	.48078
11.0	.97283	.45265 E-01	.48020
12.0	.97126	.49842 E-01	.47936
13.0	.96917	.54529 E-01	.47826
14.0	.96659	.59332 E-01	.47693
15.0	.96353	.64257 E-01	.47535

- Fuel region radius = $R_o = 100.0$ cm
- Hydrogen region thickness = $R_p - R_o = 50.0$ cm
- Infinite beryllium oxide reflector

$$- R_o = R_o / \sqrt[3]{2} \text{ and } R_H = \left[(R_p^3 + R_o^3) / 2 \right]^{1/3}$$

Table 7-6 shows that the highest k_{eff} value is obtained for a hydrogen density factor = 8.7 and neutron absorption in hydrogen varies almost proportionately with hydrogen density factor. Illustrated in Figure 7-6 is the hydrogen limiting density factor for the reference reactor configuration ($D.F.^* = 8.7$).

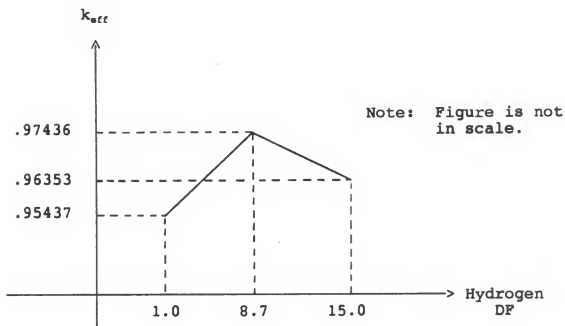


Figure 7-6. Hydrogen limiting density factor for the three-region gas core reactor problem.

Oscillations of the Fuel-Hydrogen Boundary

As a particular example of what can be accomplished using the multigroup albedo theory ideas, consider the problem of predicting and understanding the reactivity (and, thus neutronic feedback) variation due to perturbation of the fuel-hydrogen boundary (e.g., the boundary shape shown as a dashed line in Figure 7-7).

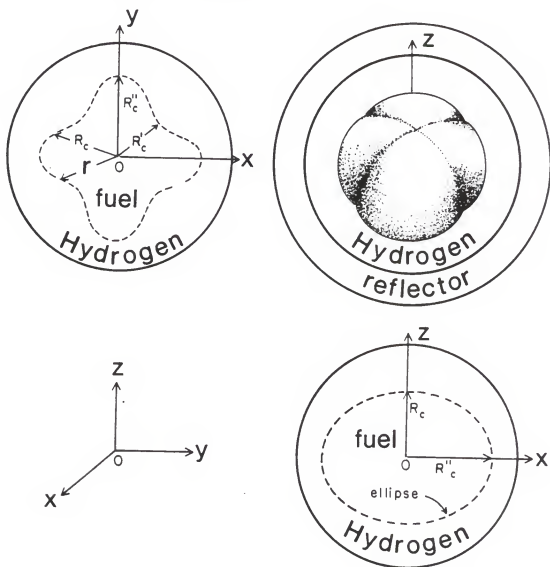


Figure 7-7. Schematic drawing showing a particular oscillation of the fuel-hydrogen boundary of the three-region gas core reactor problem.

The various possible modes of the three-dimensional fuel-hydrogen boundary of the three-region gas core reactor problem cannot be effectively handled by conventional neutronic calculation methods such as XSDRNPM.

Transmission coefficient values for the fuel, hydrogen, and reflector secondary neutron sources

The variation of all forty transmission coefficient values generated by such standing waves of the fuel-hydrogen boundary is an algebraic-geometric exercise that can be solved by using: (1) a plastic mock-up to obtain the length and solid angle values required to estimate the transmission coefficients, and (2) an analytical approach that for complex standing wave geometry can be tougher than the mock-up solution. After the evaluation of the transmission coefficients, the perturbed fuel-hydrogen boundary reactor problem is performed exactly as the one-dimensional three-region gas core reactor.

k_{eff} variation due to a perturbed fuel-hydrogen boundary

For convenience, the two-dimensional standing wave perturbation of the fuel-hydrogen boundary in Figure 7-7 can be described by the magnitude position vector function

$$r = R_c + (R_c'' - R_c) \cos \left(\frac{2\pi l}{\lambda} \right) \quad (7-11)$$

where λ is the wavelength of the surface wave, R_c'' is its maximum amplitude and ℓ is the arc length measured on the circle of radius R_c . It is assumed that

$$(1) \quad \frac{R_c'' - R_c}{R_c} = \frac{R_c - R_c'}{R_c} \text{ variation is small where } R_c' \text{ is}$$

the wave minimum amplitude, and

$$(2) \quad \text{the wavelength of the wave is given by } \lambda = \frac{\pi}{2} R_c.$$

The resulting changes in k_{eff} values by varying $(R_c'' - R_c)$ is shown in Table 7-7.

TABLE 7-7. EFFECTS DUE TO OSCILLATING FUEL-HYDROGEN BOUNDARY ON THE THREE-REGION GAS CORE REACTOR PROBLEM PROPERTIES USING THE FOUR-GROUP ALBEDO METHOD

$\frac{R_c'' - R_c}{R_c} \times 100\%$	Four-Group Albedo Theory Method		
	k_{eff}	P_{TNL}^H	P_{TNL}^F
2.0	.95437	.39067 E-02	.47169
1.0	.95083	.39427 E-02	.46994
2.0	.95101	.39411 E-02	.47003
3.0	.95121	.39395 E-02	.47013

- $R_c = 100.0$ cm
- $R_p = 150.0$ cm
- Infinite beryllium oxide reflector
- $R_c - R_c / \sqrt[3]{2}$ and $R_H = \left[(R_p^3 + R_c^3) / 2 \right]^{1/3}$

Table 7-7 shows that for $1\% \leq \frac{R_c'' - R_c}{R_c} \leq 3\%$ the k_{eff} values are smaller than the k_{eff} result obtained for the unperturbed spherical fuel-hydrogen boundary case. This result is

due to the increase of the fuel transmittance and consequent reduction in the absorption rate of neutrons by the nuclei of the fuel region.

Nonuniform Atomic Density Distribution in the Fuel Region

For the previous neutronic calculations performed on the gas core reactor problems the gaseous fuel was assumed to be uniformly distributed within the fuel region. Assuming that the total fuel mass is constant, a nonuniform atomic density distribution in the fuel region is used to determine the impact of such behavior on the value of k_{eff} .

Linear atomic density distribution

A linear atomic density distribution, $\rho(r)$, is assumed and given by the expression

$$\rho(r) = \rho_0' (1 + A r)$$

where

A = the angular coefficient and

r = the position vector magnitude

The fuel mass conservation condition is

$$\int_{r=0}^{R_0} \rho_0' (1 + A r) 4\pi r^2 dr = \int_{r=0}^{R_0} \rho_0 4\pi r^2 dr \quad (7-12)$$

where

ρ_0 = the initial uniform atomic density and

R_0 = the radius of the fuel region.

Equation 7-12 reduces to

$$\rho'_0 = \frac{\rho_0}{1 + A \frac{3R_c}{4}}, \quad (7-13)$$

thus,

$$\rho(r) = \frac{\rho_0}{1 + A \frac{3R_c}{4}} (1 + A r) \quad . \quad (7-14)$$

The corresponding fuel macroscopic cross section, $\Sigma(r)$, is given by

$$\Sigma(r) = \frac{1 + A r}{1 + A \frac{3R_c}{4}} \Sigma_0 \quad (7-15)$$

where

Σ_0 = the fuel macroscopic cross section for the
initial uniform atomic density.

Fuel secondary source position

As done previously, it is assumed that the neutron sources are isotropic and are homogeneously distributed at the surface of a sphere of radius R_0 given by the equation

$$\int_{r=0}^{R_0} \rho'_0 (1 + A r) 4\pi r^2 dr = \int_{r=R_0}^{R_c} \rho'_0 (1 + A r) 4\pi r^2 dr, \quad (7-16)$$

i.e., this imaginary sphere contains half of the total fuel mass. Thus, R_0 must satisfy the quartic equation given by

$$R_o^4 + \frac{4}{3A} R_o^3 - \frac{2R_c^3}{A} \left(\frac{1}{3} + A \frac{R_c}{4} \right) = 0 \quad (7-17)$$

and

$$0 < R_o < R_c \quad .$$

Cases investigated: $A = \pm 0.01$

Figure 7-8 illustrates the variation of $\rho(r)$ and $\Sigma(r)$ in the fuel region for the cases $A = \pm 0.01$.

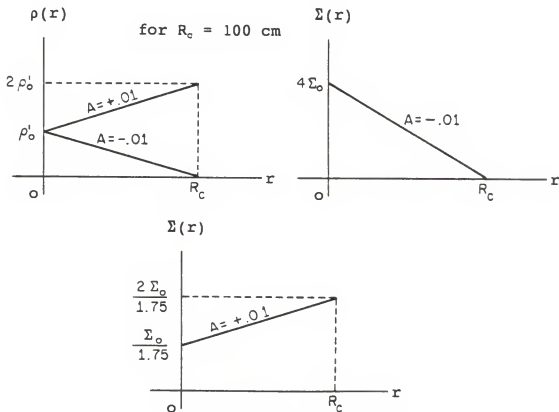


Figure 7-8. Qualitative graph of $\rho(r)$ and $\Sigma(r)$ in the fuel region of a three-region gas core reactor problem.

The resulting k_{eff} variation is shown in Table 7-8.

TABLE 7-8. EFFECTS DUE TO NONUNIFORM ATOMIC DENSITY DISTRIBUTION IN THE FUEL REGION ON THE THREE-REGION GAS CORE REACTOR PROBLEM USING THE FOUR-GROUP ALBEDO METHOD

A	Four-Group Albedo Theory Method			
	k_{eff}	P_{TNL}^H	P_{TNL}^{FUEL}	R_o
- 0.01	1.1231	.27303 E-02	.55867	61.43
0.00	.95437	.39067 E-02	.46956	79.37
+ 0.01	.90812	.42273 E-02	.44840	81.58

- $R_o = 100$ cm
- $R_p = 150$ cm
- $R_B = \left[(R_p^3 + R_o^3) / 2 \right]^{1/3}$
- R_o in cm

The results in Table 7-8 indicate that relative to a uniform density core $A = -0.01$ yields an increase in k_{eff} , while $A = +0.01$ yields a decrease in k_{eff} . It appears that the required variation of R_o yields effects on core transmittance which completely explain these results.

The program ALB4 employed for these numerical calculations is in Appendix G.

CHAPTER VIII CONCLUSIONS AND RECOMMENDATIONS

Introduction

In this dissertation is described a multigroup neutron albedo theory method is described which can be used to perform neutronic calculations for thermal nuclear reactors. In the neutron albedo method, description of average neutron histories is accomplished through the specification of reflection and transmission probabilities at each neutron traversal of the core-reflector interface. For each energy group and for energy group transfers, these neutron albedo parameters are expressed in terms of the cross sections and geometry of the core and the reflector. The effective neutron multiplication factor and reactivity can be expressed as analytical functions of the neutron albedos, and thereby, as direct functions of cross sections and geometry. Thus, the neutron albedo theory method is an analytical method that does not try to solve an explicit describing equation for the neutron flux but uses such relations to obtain partial information that is combined with the average behavior of the neutrons in the physical system to obtain nonleakage probability values and effective and infinite neutron multiplication factors. Therefore, the

neutron albedo theory methodology is very different from the methodology used in conventional discrete ordinate computer codes such as XSDRNPM.

Final Conclusion

The primary goal of developing a multigroup neutron albedo method has been achieved herein. Moreover, application to neutronic calculations such as nonleakage probabilities and effective neutron multiplication factors, and investigation of the main parameters of the four-group albedo method have been accomplished with the aid of three calculational reactor problems. In Chapters V, VI, and VII of this work, the four-group albedo theory is used to calculate the total nonleakage probability and effective neutron multiplication factor for reactors with large and small optical-dimension cores displaying the matrix of parameter dissection of the total principal and secondary nonleakage probabilities. In Table 8-1 is shown the importance of the total principal and secondary nonleakage probabilities for the analyzed large and small optical-path-length (O.P.L.) core (two-region) reactor problems.

TABLE 8-1. LARGE VERSUS SMALL OPTICAL-PATH-LENGTH CORE (TWO-REGION) REACTOR PROBLEM VALUES USING THE FOUR-GROUP ALBEDO METHOD

	Four-Group Albedo Method	
	Large O.P.L.	Small O.P.L.
P_o	.89433	.92404 E-03
P	.46129 E-01	.53460
$P_{TNL} = P_o + P$.94046	.53553
P_o/P_{TNL}	~ 95.1%	0.2%
P/P_{TNL}	~ 4.9%	99.8%

From Table 8-1, the extremely important role of the reflector-moderator in small O.P.L. gas core concepts is demonstrated. The total secondary nonleakage probability, P , represents approximately 100% of the total nonleakage probability, P_{TNL} , for such cores.

Unfortunately, the multigroup neutron albedo method usually depends on assumptions that make the final solution less accurate than the more rigorous solution methods (such as XSDRNPM). Nevertheless, for many cases (such as gas core reactor problems), if the approximate solution is reasonably acceptable, the loss in accuracy and generality is more than outweighed by the increased simplicity in the calculations, the resultant reduction in computer time, and especially the clarity of identification of various physical phenomena influences. What is sought here is the development of a simplified approach to neutronic problems under the neutron albedo point of view with a reasonably quick and inexpensive

method that is clearly more understandable than conventional neutron transport methods. In order to quickly see relations among the albedos, transmittances, nonleakage probabilities, and the effective neutron multiplication factor, the multigroup neutron albedo method represents a powerful and helpful tool.

Recommendations for Future Work

The significant amount of research performed on the small optical-dimension core reactors (two- and three-region), as presented in this dissertation, has furnished some insights and intuitive procedures. Future research should focus on developing analytical and intuitive methods of determining the core transmission coefficient values which actually constitute the key item for accurate effective neutron multiplication factor results.

Application of the albedo approach to the determination of neutronic-fluid dynamic coupling in gas core reactor concepts [15] could be an important practical application that is uniquely suited to the method. Such an effort, if successful, would certainly justify, on practical grounds, the considerations and ideas developed in this dissertation.

APPENDIX A ALGORITHM FOR mS_i

The following algorithm for calculating mS_i is developed employing the known sources of initial leakage of $m-1$ group albedo theory, ${}^{m-1}S_i$.

Initial step: Evaluation of mS_i ($i < m$)

Using the assumed known expression of ${}^{m-1}S_i$, replace the superscript ($m-1$) by the superscript m ; thus,

$${}^mS_i = {}^{m-1}S_i \quad (m - 1 \rightarrow m). \quad (A-1)$$

Final step: Evaluation of mS_m

$${}^mS_m = (1 - P_{o_m}) \left(\chi_m + \sum_{i=1}^{m-1} \frac{{}^mS_i}{1 - P_{o_i}} P_{o_i} \frac{\Sigma_{S_{im}}}{\Sigma_{R_i}} \right). \quad (A-2)$$

APPENDIX B
TWO-GROUP PARTIAL SECONDARY
NONLEAKAGE PROBABILITY DERIVATION

Here Equation 2-11 is derived under the point of view of two-group albedo theory. The total secondary nonleakage probability, P , is a sum of two separate processes.

Specifically,

- (1) The probability that a neutron produced by fission in the core leaks as thermal ($i=2$) and is absorbed in the core, here denoted by 2P_2 , and
- (2) The probability that a neutron produced by fission in the core leaks as fast ($i=1$) neutron and is absorbed in the core, here denoted by 2P_1 .

Hence, one can write $P = {}^2P_1 + {}^2P_2$ or $P = {}^2S_1 {}^2Q_1 + {}^2S_2 {}^2Q_2$

where 2Q_1 is the probability that if one neutron leaks as a group 1 neutron from the core, it will be absorbed in the core and 2Q_2 is the probability that if one neutron leaks as a group 2 neutron from the core, it will be absorbed in the core. 2S_1 and 2S_2 values are shown in Figure 2-1.

First, with the aid of Figure B-1, the 2Q_2 expression is determined.

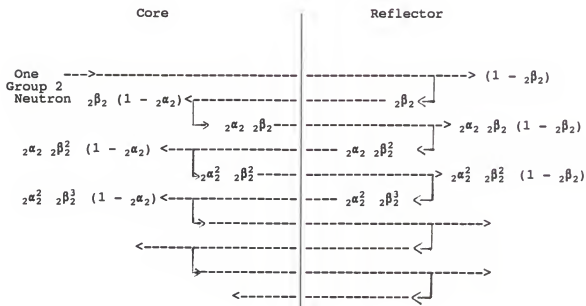


Figure B-1. Schematic drawing showing the two-group albedo approach strategy for determining 2Q_2 expression.

In the passage of group 2 neutrons from the core to the reflector, ${}_2\beta_2$ is the probability that group 2 neutrons incident on the reflector from the core will be reflected with group 2 energy. Similarly, ${}_2\alpha_2$ is the probability that group 2 neutrons incident on the core from the reflector will be reflected with group 2 energy. From Figure B-1, the 2Q_2 is given by

$$\begin{aligned}
 {}^2Q_2 &= {}_2\beta_2 (1 - {}_2\alpha_2) + {}_2\alpha_2 {}_2\beta_2^2 (1 - {}_2\alpha_2) + {}_2\alpha_2^2 {}_2\beta_2^3 (1 - {}_2\alpha_2) + \dots = \\
 &= {}_2\beta_2 (1 - {}_2\alpha_2) \left[1 + {}_2\alpha_2 {}_2\beta_2 + ({}_2\alpha_2 {}_2\beta_2)^2 + \dots \right] = \\
 &= \frac{(1 - {}_2\alpha_2) {}_2\beta_2}{1 - {}_2\alpha_2 {}_2\beta_2}
 \end{aligned}
 \tag{B-1}$$

From Figure B-1 and Equation B-1, one determines the probability that if one neutron leaks as a group 2 neutron from the reflector, it will be absorbed in the reflector, here denoted by ${}^2\bar{Q}_2$. Thus,

$${}^2\bar{Q}_2 = \frac{(1 - {}_2\beta_2) {}_2\alpha_2}{1 - {}_2\alpha_2 {}_2\beta_2} \quad (\text{B-2})$$

and

$$1 - {}^2\bar{Q}_2 = \frac{1 - {}_2\alpha_2}{1 - {}_2\alpha_2 {}_2\beta_2} \quad (\text{B-3})$$

represents the probability that if one neutron leaks as group 2 neutron from the reflector, it will be absorbed in the core.

Finally, with the aid of Figure B-2, one determines the 2Q_1 expression.

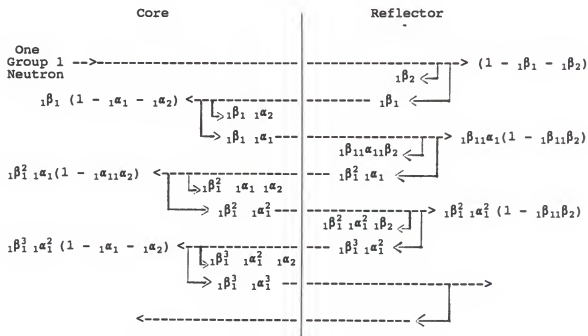


Figure B-2. Schematic drawing showing the two-group albedo approach strategy for determining 2Q_1 expression.

In the passage of group 1 neutrons from the core to reflector, ${}_1\beta_1$ is the probability that group 1 neutrons incident on the reflector from the core will be reflected with group 1 energy and ${}_1\beta_2$ is the probability that group 1 neutrons incident on the reflector from the core will be reflected with group 2 energy. There are similar definitions for ${}_1\alpha_1$ and ${}_1\alpha_2$. From Figure B-2 and Equations B-1 and B-3, the 2Q_1 equation is obtained given by

$$\begin{aligned}
 {}^2Q_1 &= [{}_1\beta_1 (1 - {}_1\alpha_1 - {}_1\alpha_2) + {}_1\beta_1^2 {}_1\alpha_1 (1 - {}_1\alpha_1 - {}_1\alpha_2) + {}_1\beta_1^3 {}_1\alpha_1^2 (1 - {}_1\alpha_1 - {}_1\alpha_2) + \dots] + \\
 &\quad [{}_1\beta_2 + {}_1\beta_1 {}_1\alpha_1 {}_1\beta_2 + {}_1\beta_1^2 {}_1\alpha_1^2 {}_1\beta_2 + \dots] (1 - {}^2Q_2) + \\
 &\quad [{}_1\beta_1 {}_1\alpha_2 + {}_1\beta_1^2 {}_1\alpha_1 {}_1\alpha_2 + {}_1\beta_1^3 {}_1\alpha_1^2 {}_1\alpha_2 + \dots] {}^2Q_2 = \\
 &= \frac{{}_1\beta_1(1 - {}_1\alpha_1 - {}_1\alpha_2)}{1 - {}_1\beta_1 {}_1\alpha_1} + \frac{{}_1\beta_2}{1 - {}_1\beta_1 {}_1\alpha_1} \cdot \frac{1 - {}_2\alpha_2}{1 - {}_2\alpha_2 {}_2\beta_2} \\
 &\quad + \frac{{}_1\beta_1 {}_1\alpha_2}{1 - {}_1\alpha_1 {}_1\beta_1} \cdot \frac{(1 - {}_2\alpha_2) {}_2\beta_2}{(1 - {}_2\alpha_2 {}_2\beta_2)} = \\
 &= \frac{(1 - {}_1\alpha_1) {}_1\beta_1}{1 - {}_1\alpha_1 {}_1\beta_1} - \frac{{}_1\beta_1 {}_1\alpha_2}{1 - {}_1\alpha_1 {}_1\beta_1} \cdot \frac{1 - {}_2\beta_2}{1 - {}_2\alpha_2 {}_2\beta_2} \\
 &\quad + \frac{{}_1\beta_2}{1 - {}_1\alpha_1 {}_1\beta_1} \cdot \frac{1 - {}_2\alpha_2}{1 - {}_2\alpha_2 {}_2\beta_2} .
 \end{aligned}$$

(B-4)

Equations B-1 and B-4 represent the desired result of the probabilities 2Q_2 and 2Q_1 , respectively.

APPENDIX C FOUR-GROUP CORE ALBEDO DERIVATION

Here the four-group core albedos are derived for a spherical core of radius R . All group constants refer to the homogeneous core.

Determination of ${}_1\alpha_1$, ${}_1\alpha_2$, ${}_1\alpha_3$, and ${}_1\alpha_4$

The governing equations and the boundary conditions for the determination of ${}_1\alpha_1$, ${}_1\alpha_2$, ${}_1\alpha_3$, and ${}_1\alpha_4$ are

$$-D_1 \nabla^2 \phi_1 + \Sigma_{R_1} \phi_1 = 0$$

$$-\Sigma_{S_{12}} \phi_1 - D_2 \nabla^2 \phi_2 + \Sigma_{R_2} \phi_2 = 0$$

$$-\Sigma_{S_{13}} \phi_1 - \Sigma_{S_{23}} \phi_2 - D_3 \nabla^2 \phi_3 + \Sigma_{R_3} \phi_3 = 0$$

$$-\Sigma_{S_{14}} \phi_1 - \Sigma_{S_{24}} \phi_2 - \Sigma_{S_{34}} \phi_3 - D_4 \nabla^2 \phi_4 + \Sigma_{R_4} \phi_4 = 0$$

at $r = 0$, ϕ_1 , ϕ_2 , ϕ_3 , and ϕ_4 are finite and

at $r = R$, $J_-]_{i-1} = 1$, $J_-]_{i-2,3,4} = 0$.

Finally, at $r = R$, ${}_1\alpha_1 = \frac{J_+]_{i-1}}{J_-]_{i-1}}$,

$${}_1\alpha_2 = \frac{J_+]_{i-2}}{J_-]_{i-1}}, \quad {}_1\alpha_3 = \frac{J_+]_{i-3}}{J_-]_{i-1}}, \quad \text{and} \quad {}_1\alpha_4 = \frac{J_+]_{i-4}}{J_-]_{i-1}}.$$

Determination of ${}_2\alpha_2$, ${}_2\alpha_3$, and ${}_2\alpha_4$

The governing equations and the boundary conditions for the determination of ${}_2\alpha_2$, ${}_2\alpha_3$, and ${}_2\alpha_4$ are

$$- D_2 \nabla^2 \phi_2 + \Sigma_{R_2} \phi_2 = 0$$

$$- \Sigma_{S_{23}} \phi_2 - D_3 \nabla^2 \phi_3 + \Sigma_{R_3} \phi_3 = 0$$

$$- \Sigma_{S_{24}} \phi_2 - \Sigma_{S_{34}} \phi_3 - D_4 \nabla^2 \phi_4 + \Sigma_{R_4} \phi_4 = 0$$

at $r = 0$, ϕ_2 , ϕ_3 , and ϕ_4 are finite and

at $r = R$, $J_-]_{i=2} = 1$, $J_-]_{i=3,4} = 0$.

Finally, at $r = R$, ${}_2\alpha_2 = \frac{J_+]_{i=2}}{J_-]_{i=2}}$,

${}_2\alpha_3 = \frac{J_+]_{i=3}}{J_-]_{i=2}}$, and ${}_2\alpha_4 = \frac{J_+]_{i=4}}{J_-]_{i=2}}$.

Determination of ${}_3\alpha_3$ and ${}_3\alpha_4$

The governing equations and the boundary conditions for the determination of ${}_3\alpha_3$ and ${}_3\alpha_4$ are

$$-D_3 \nabla^2 \phi_3 + \Sigma_{R_3} \phi_3 = 0$$

$$- \Sigma_{S_{34}} \phi_3 - D_4 \nabla^2 \phi_4 + \Sigma_{R_4} \phi_4 = 0$$

at $r = 0$, ϕ_3 and ϕ_4 are finite and

at $r = R$, $J_-]_{i=3} = 1$, $J_-]_{i=4} = 0$.

Finally, at $r = R$, ${}_3\alpha_3 = \frac{J_+]_{i=3}}{J_-]_{i=3}}$ and ${}_3\alpha_4 = \frac{J_+]_{i=4}}{J_-]_{i=3}}$.

Determination of ${}_4\alpha_4$

The governing equation and the boundary conditions for the determination of ${}_4\alpha_4$ are

$$-D_4 \nabla^2 \phi_4 + \Sigma_{R_4} \phi_4 = 0$$

at $r = 0$, ϕ_4 is finite and at $r = R$, $J_-]_{i=4} = 1$.

Finally, at $r = R$, ${}_4\alpha_4 = \frac{J_+]_{i=4}}{J_-]_{i=4}}$.

The analytical expressions of the core albedos (${}_i\alpha_i$) are given below.

$$\underline{{}_i\alpha_i \quad (i = 1, 2, 3, 4)}$$

$${}_i\alpha_i = \frac{1 - 2 D_i \left[L_i \coth (L_i R) - \frac{1}{R} \right]}{1 + 2 D_i \left[L_i \coth (L_i R) - \frac{1}{R} \right]} \quad (C-1)$$

where

$$L_i = \sqrt{\frac{\Sigma_{R_i}}{D_i}}.$$

$$\underline{{}_i\alpha_{i+1} \quad (i = 1, 2, 3)}$$

$${}_i\alpha_{i+1} = \frac{4 \Sigma_{S_{i+1}}}{L_i^2 - L_{i+1}^2}.$$

(C-2)

$$\frac{L_i \coth (L_i R) - L_{i+1} \coth (L_{i+1} R)}{\left[1 + 2 D_i \left[L_i \coth (L_i R) - \frac{1}{R} \right] \right] \left[1 + 2 D_{i+1} \left[L_{i+1} \coth (L_{i+1} R) - \frac{1}{R} \right] \right]}$$

$$\underline{{}_i\alpha_{i+2} \quad (i = 1, 2)}$$

$${}_i\alpha_{i+2} = \frac{1}{2R} [E_{i3} \sinh (L_{i+2} R) + F_{i2} \sinh (L_i R) + F_{i3} \sinh (L_{i+1} R)]$$

where

$$E_{i1} = \frac{4R}{\left[1 + 2D_i \left[L_i \coth (L_i R) - \frac{1}{R} \right] \right]} \sinh (L_i R)$$

$$F_{i1} = \frac{\Sigma_{s_{i1,1}} E_{i1}}{D_{i+1} (L_{i+1}^2 - L_i^2)}$$

$$E_{i2} = - \frac{F_{i1} \sinh (L_i R) \left[1 + 2D_{i+1} \left[L_i \coth (L_i R) - \frac{1}{R} \right] \right]}{\sinh (L_{i+1} R) \left[1 + 2D_{i+1} \left[L_{i+1} \coth (L_{i+1} R) - \frac{1}{R} \right] \right]}$$

$$F_{i2} = \frac{\Sigma_{s_{i1,2}} E_{i1} + \Sigma_{s_{i+1,1,2}} F_{i1}}{D_{i+2} (L_{i+2}^2 - L_i^2)} \quad (C-3)$$

$$F_{i3} = \frac{\Sigma_{s_{i+1,1,2}} E_{i2}}{D_{i+2} (L_{i+2}^2 - L_{i+1}^2)} \text{ and}$$

$$\begin{aligned} E_{i3} = & - \left[F_{i2} \sinh (L_i R) \left[1 + 2D_{i+2} \left[L_i \coth (L_i R) - \frac{1}{R} \right] \right] + \right. \\ & F_{i3} \sinh (L_{i+1} R) \left[1 + 2D_{i+2} \left[L_{i+1} \coth (L_{i+1} R) - \frac{1}{R} \right] \right] \Big] / \\ & \left[\sinh (L_{i+2} R) \left[1 + 2D_{i+2} \left[L_{i+2} \coth (L_{i+2} R) - \frac{1}{R} \right] \right] \right] . \end{aligned}$$

$$\underline{{}_1\alpha_4 \text{ (i = 1)}}$$

$${}_1\alpha_4 = \frac{1}{2R} [E_{i4} \sinh (L_4 R) + F_{i4} \sinh (L_1 R) + F_{i5} \sinh (L_2 R) + \\ F_{i6} \sinh (L_3 R)]$$

where

$$F_{i4} = \frac{\Sigma_{s_{14}} E_{i1} + \Sigma_{s_{24}} F_{i1} + \Sigma_{s_{34}} F_{i2}}{D_4 (L_4^2 - L_1^2)}$$

$$F_{i5} = \frac{\Sigma_{s_{24}} E_{i2} + \Sigma_{s_{34}} F_{i3}}{D_4 (L_4^2 - L_1^2)} \quad (C-4)$$

$$F_{i6} = \frac{\Sigma_{s_{34}} E_{i3}}{D_4 (L_4^2 - L_1^2)} , \text{ and}$$

$$E_{i4} = - \left[F_{i4} \sinh (L_1 R) \left[1 + 2 D_4 \left[L_1 \coth (L_1 R) - \frac{1}{R} \right] \right] + \right.$$

$$F_{i5} \sinh (L_2 R) \left[1 + 2 D_4 \left[L_2 \coth (L_2 R) - \frac{1}{R} \right] \right] +$$

$$F_{i6} \sinh (L_3 R) \left[1 + 2 D_4 \left[L_3 \coth (L_3 R) - \frac{1}{R} \right] \right] \Big] /$$

$$\left[\sinh (L_4 R) \left[1 + 2 D_4 \left[L_4 \coth (L_4 R) - \frac{1}{R} \right] \right] \right] .$$

APPENDIX D
FOUR-GROUP REFLECTOR ALBEDO
DERIVATION FOR A FINITE REFLECTOR

In this appendix, the four-group reflector albedos are derived for a reflector of inner radius R and outer radius H. The reflector thickness is given by the expression $H - R$.

Determination of ${}_1\beta_1$, ${}_1\beta_2$, ${}_1\beta_3$, and ${}_1\beta_4$

The neutrons produced by the (n, 2n) interaction in the reflector are only accounted for in ${}_1\beta_1$, ${}_1\beta_2$, ${}_1\beta_3$, and ${}_1\beta_4$ values because the (n, 2n) reaction neutrons are born in the first energy group.

The governing equations and the boundary conditions for the determination of ${}_1\beta_1$, ${}_1\beta_2$, ${}_1\beta_3$, and ${}_1\beta_4$ are

$$-D_1 \nabla^2 \phi_1 + \Sigma_{R_1} \phi_1 - 2 \Sigma_{n,2n_1} \phi_1$$

$$- \Sigma_{S_{12}} \phi_1 - D_2 \nabla^2 \phi_2 + \Sigma_{R_2} \phi_2 = 0$$

$$- \Sigma_{S_{13}} \phi_1 - \Sigma_{S_{23}} \phi_2 - D_3 \nabla^2 \phi_3 + \Sigma_{R_3} \phi_3 = 0$$

$$- \Sigma_{S_{14}} \phi_1 - \Sigma_{S_{24}} \phi_2 - \Sigma_{S_{34}} \phi_3 - D_4 \nabla^2 \phi_4 + \Sigma_{R_4} \phi_4 = 0$$

at $r = R$, $J_{+} |_{i=1} = 1$, $J_{+} |_{i=2,3,4} = 0$ and

at $r = H$, $J_{-} |_{i=1,2,3,4} = 0$.

Finally, at $r = R$, ${}_1\beta_1 = \frac{J_- J_{i-1}}{J_+ J_{i-1}}$,

$${}_1\beta_2 = \frac{J_- J_{i-2}}{J_+ J_{i-1}}, \quad {}_1\beta_3 = \frac{J_- J_{i-3}}{J_+ J_{i-1}}, \quad \text{and} \quad {}_1\beta_4 = \frac{J_- J_{i-4}}{J_+ J_{i-1}}.$$

Determination of ${}_2\beta_2$, ${}_2\beta_3$, and ${}_2\beta_4$

The governing equations and the boundary conditions for the determination of ${}_2\beta_2$, ${}_2\beta_3$, and ${}_2\beta_4$ are

$$-D_2 \nabla^2 \phi_2 + \Sigma_{R_2} \phi_2 = 0$$

$$-\Sigma_{S_{23}} \phi_2 - D_3 \nabla^2 \phi_3 + \Sigma_{R_3} \phi_3 = 0$$

$$-\Sigma_{S_{24}} \phi_2 - \Sigma_{S_{34}} \phi_3 - D_4 \nabla^2 \phi_4 + \Sigma_{R_4} \phi_4 = 0$$

at $r = R$, $J_+ J_{i-2} = 1$, $J_+ J_{i-3,4} = 0$ and

at $r = H$, $J_- J_{i-2,3,4} = 0$.

Finally, at $r = R$, ${}_2\beta_2 = \frac{J_- J_{i-2}}{J_+ J_{i-2}}$,

$${}_2\beta_3 = \frac{J_- J_{i-3}}{J_+ J_{i-2}}, \quad \text{and} \quad {}_2\beta_4 = \frac{J_- J_{i-4}}{J_+ J_{i-2}}.$$

Determination of ${}_3\beta_3$ and ${}_3\beta_4$

The governing equations and the boundary conditions for the determination of ${}_3\beta_3$ and ${}_3\beta_4$ are

$$-D_3 \nabla^2 \phi_3 + \Sigma_{R_3} \phi_3 = 0$$

$$- \Sigma_{S_{34}} \phi_3 - D_4 \nabla^2 \phi_4 + \Sigma_{R_4} \phi_4 = 0$$

at $r = R$, $J_+]_{i=3} = 1$, $J_+]_{i=4} = 0$ and

at $r = H$, $J_-]_{i=3,4} = 0$.

Finally, at $r = R$, ${}_3\beta_3 = \frac{J_-]_{i=3}}{J_+]_{i=3}}$ and ${}_3\beta_4 = \frac{J_-]_{i=4}}{J_+]_{i=3}}$.

Determination of ${}_4\beta_4$

The governing equations and the boundary conditions for the determination of ${}_4\beta_4$ are

$$-D_4 \nabla^2 \phi_4 + \Sigma_{R_4} \phi_4 = 0$$

at $r = R$, $J_+]_{i=4} = 1$ and

at $r = H$, $J_-]_{i=4} = 0$.

Finally, at $r = R$, ${}_4\beta_4 = \frac{J_-]_{i=4}}{J_+]_{i=4}}$.

The analytical expressions of the reflector albedos (${}_i\beta_i$) are given below.

$$\underline{{}_i\beta_i \quad (i = 1, 2, 3, 4)}$$

$${}_i\beta_i = C_{i1} \frac{e^{-k_i R}}{R} \left[\frac{1}{4} - \frac{D_i}{2} \left(k_i + \frac{1}{R} \right) \right] + \quad (D-1)$$

$$C_{i2} \frac{e^{+k_i R}}{R} \left[\frac{1}{4} + \frac{D_i}{2} \left(k_i - \frac{1}{R} \right) \right]$$

where

$$k_i = \sqrt{\frac{\Sigma_{R_i} - 2 \Sigma_{n, 2n_i}}{D_i}},$$

and C_{i1} and C_{i2} are the solution of the following system

$$0 - C_{i1} e^{-k_i H} \left[1 - 2 D_i \left(k_i + \frac{1}{H} \right) \right] + C_{i2} e^{+k_i H} \left[1 + 2 D_i \left(k_i - \frac{1}{H} \right) \right]$$

$$4R - C_{i1} e^{-k_i R} \left[1 + 2 D_i \left(k_i + \frac{1}{R} \right) \right] + C_{i2} e^{+k_i R} \left[1 - 2 D_i \left(k_i - \frac{1}{R} \right) \right] .$$

$$\underline{{}_i\beta_{i+1} \quad (i = 1, 2, 3)}$$

$${}_i\beta_{i+1} = \frac{1}{2R} [C_{i3} e^{-k_{i+1} R} + C_{i4} e^{+k_{i+1} R} + G_{i1} e^{-k_i R} + G_{i2} e^{+k_i R}] \quad (D-2)$$

where

$$G_{i1} = - \frac{\Sigma g_{i1+i}}{D_{i+1} (k_i^2 - k_{i+1}^2)} C_{i1} ,$$

$$G_{i2} = - \frac{\Sigma g_{i2+i}}{D_{i+1} (k_i^2 - k_{i+1}^2)} C_{i2} ,$$

and C_{i3} and C_{i4} are the solutions of the system

$$- G_{i1} e^{-k_i R} \left[1 + 2 D_{i+1} (k_i + R) \right] - G_{i2} e^{+k_i R} \left[1 - 2 D_{i+1} \left(k_i - \frac{1}{R} \right) \right] -$$

$$C_{i3} e^{-k_{i+1} R} \left[1 + 2 D_{i+1} (k_{i+1} + R) \right] + C_{i4} e^{+k_{i+1} R} \left[1 - 2 D_{i+1} (k_{i+1} - R) \right]$$

and

$$- G_{i1} e^{-k_i R} \left[1 - 2 D_{i+1} \left(k_i + \frac{1}{H} \right) \right] - G_{i2} e^{+k_i R} \left[1 + 2 D_{i+1} \left(k_i - \frac{1}{H} \right) \right] -$$

$$C_{i3} e^{-k_{i+1} R} \left[1 - 2 D_{i+1} \left(k_{i+1} + \frac{1}{H} \right) \right] + C_{i4} e^{+k_{i+1} R} \left[1 + 2 D_{i+1} \left(k_{i+1} - \frac{1}{H} \right) \right] .$$

$$\underline{{}_i\beta_{i+2} \quad (i = 1, 2)}$$

$$\begin{aligned} {}_i\beta_{i+2} = & \frac{1}{2R} \left[C_{i5} e^{-k_{i+2}R} + C_{i6} e^{+k_{i+2}R} + G_{i3} e^{-k_iR} + \right. \\ & \left. G_{i4} e^{+k_iR} + G_{i5} e^{-k_{i+1}R} + G_{i6} e^{+k_{i+1}R} \right] \end{aligned} \quad (D-3)$$

where

$$G_{i3} = - \frac{\Sigma_{s_{i+2}} C_{i1} + \Sigma_{s_{i+1, i+2}} G_{i1}}{D_{i+2} (k_i^2 - k_{i+2}^2)}$$

$$G_{i4} = - \frac{\Sigma_{s_{i+2}} C_{i2} + \Sigma_{s_{i+1, i+2}} G_{i2}}{D_{i+2} (k_i^2 - k_{i+2}^2)}$$

$$G_{i5} = - \frac{\Sigma_{s_{i+1, i+2}}}{D_{i+2} (k_{i+1}^2 - k_{i+2}^2)} C_{i3}$$

$$G_{i6} = - \frac{\Sigma_{s_{i+1, i+2}}}{D_{i+2} (k_{i+1}^2 - k_{i+2}^2)} C_{i4} ,$$

and C_{i5} and C_{i6} are the solutions of the following system

$$\begin{aligned}
& - G_{13} e^{-k_1 R} \left[1 + 2 D_{1+2} \left(k_1 + \frac{1}{R} \right) \right] - G_{14} e^{+k_1 R} \left[1 - 2 D_{1+2} \left(k_1 - \frac{1}{R} \right) \right] \\
& - G_{15} e^{-k_{1+1} R} \left[1 + 2 D_{1+2} \left(k_{1+1} + \frac{1}{R} \right) \right] - G_{16} e^{+k_{1+1} R} \left[1 - 2 D_{1+2} \left(k_{1+1} - \frac{1}{R} \right) \right] - \\
& C_{15} e^{-k_{1+2} R} \left[1 + 2 D_{1+2} \left(k_{1+2} + \frac{1}{R} \right) \right] + C_{16} e^{+k_{1+2} R} \left[1 - 2 D_{1+2} \left(k_{1+2} - \frac{1}{R} \right) \right]
\end{aligned}$$

and

$$\begin{aligned}
& - G_{13} e^{-k_1 H} \left[1 - 2 D_{1+2} \left(k_1 + \frac{1}{H} \right) \right] - G_{14} e^{+k_1 H} \left[1 + 2 D_{1+2} \left(k_1 - \frac{1}{H} \right) \right] \\
& - G_{15} e^{-k_{1+1} H} \left[1 - 2 D_{1+2} \left(k_{1+1} + \frac{1}{H} \right) \right] - G_{16} e^{+k_{1+1} H} \left[1 + 2 D_{1+2} \left(k_{1+1} - \frac{1}{H} \right) \right] - \\
& C_{15} e^{-k_{1+2} H} \left[1 - 2 D_{1+2} \left(k_{1+2} + \frac{1}{H} \right) \right] + C_{16} e^{+k_{1+2} H} \left[1 + 2 D_{1+2} \left(k_{1+2} - \frac{1}{H} \right) \right] .
\end{aligned}$$

$$i\beta_4 \quad (i = 1)$$

$$i\beta_4 = \frac{1}{2R} \left[C_{17} e^{-k_4 R} + C_{18} e^{+k_4 R} + G_{17} e^{-k_1 R} + G_{18} e^{+k_1 R} + \right. \\ \left. G_{19} e^{-k_2 R} + G_{110} e^{+k_2 R} + G_{111} e^{-k_3 R} + G_{112} e^{+k_3 R} \right] \quad (D-4)$$

where

$$G_{17} = - \frac{\Sigma_{s_{14}} C_{11} + \Sigma_{s_{24}} G_{11} + \Sigma_{s_{34}} G_{13}}{D_4 (k_1^2 - k_4^2)}$$

$$G_{18} = - \frac{\Sigma_{s_{14}} C_{12} + \Sigma_{s_{24}} G_{12} + \Sigma_{s_{34}} G_{14}}{D_4 (k_1^2 - k_4^2)}$$

$$G_{19} = - \frac{\Sigma_{s_{24}} C_{13} + \Sigma_{s_{34}} G_{15}}{D_4 (k_2^2 - k_4^2)}$$

$$G_{110} = - \frac{\Sigma_{s_{24}} C_{14} + \Sigma_{s_{34}} G_{16}}{D_4 (k_2^2 - k_4^2)}$$

$$G_{111} = - \frac{\Sigma_{s_{34}} C_{15}}{D_4 (k_3^2 - k_4^2)}$$

$$G_{112} = - \frac{\Sigma_{s_{34}} C_{16}}{D_4 (k_3^2 - k_4^2)},$$

and C_{17} and C_{18} are the solutions of the following system

$$\begin{aligned}
& - G_{17} e^{-k_1 R} \left[1 + 2 D_4 \left(k_1 + \frac{1}{R} \right) \right] - G_{18} e^{+k_1 R} \left[1 - 2 D_4 \left(k_1 - \frac{1}{R} \right) \right] \\
& - G_{19} e^{-k_2 R} \left[1 + 2 D_4 \left(k_2 + \frac{1}{R} \right) \right] - G_{110} e^{+k_2 R} \left[1 - 2 D_4 \left(k_2 - \frac{1}{R} \right) \right] \\
& - G_{111} e^{-k_3 R} \left[1 + 2 D_4 \left(k_3 + \frac{1}{R} \right) \right] - G_{112} e^{+k_3 R} \left[1 - 2 D_4 \left(k_3 - \frac{1}{R} \right) \right] - \\
& C_{17} e^{-k_4 R} \left[1 + 2 D_4 \left(k_4 + \frac{1}{R} \right) \right] + C_{18} e^{+k_4 R} \left[1 - 2 D_4 \left(k_4 - \frac{1}{R} \right) \right]
\end{aligned}$$

and

$$\begin{aligned}
& - G_{17} e^{-k_1 H} \left[1 - 2 D_4 \left(k_1 + \frac{1}{H} \right) \right] - G_{18} e^{+k_1 H} \left[1 + 2 D_4 \left(k_1 - \frac{1}{H} \right) \right] \\
& - G_{19} e^{-k_2 H} \left[1 - 2 D_4 \left(k_2 + \frac{1}{H} \right) \right] - G_{110} e^{+k_2 H} \left[1 + 2 D_4 \left(k_2 - \frac{1}{H} \right) \right] \\
& - G_{111} e^{-k_3 H} \left[1 - 2 D_4 \left(k_3 + \frac{1}{H} \right) \right] - G_{112} e^{+k_3 H} \left[1 + 2 D_4 \left(k_3 - \frac{1}{H} \right) \right] - \\
& C_{17} e^{-k_4 H} \left[1 - 2 D_4 \left(k_4 + \frac{1}{H} \right) \right] + C_{18} e^{+k_4 H} \left[1 + 2 D_4 \left(k_4 - \frac{1}{H} \right) \right] .
\end{aligned}$$

Additionally, the four-group reflector albedo expressions for an infinite reflector are given by

$$\underline{{}_i\beta_i \quad (i = 1, 2, 3, 4)}$$

$${}_i\beta_i = \frac{1 - 2 D_i (k_i + \frac{1}{R})}{1 + 2 D_i (k_i + \frac{1}{R})} \quad . \quad (D-5)$$

$$\underline{{}_i\beta_{i+1} \quad (i = 1, 2, 3)}$$

$${}_i\beta_{i+1} = \frac{1}{2 R} [C_{i10} e^{-k_{i+1} R} + G_{i13} e^{-k_i R}] \quad (D-6)$$

where

$$G_{i13} = - \frac{\Sigma_{g_{i+1}}}{D_{i+1} (k_i^2 - k_{i+1}^2)} C_{i9}$$

$$C_{i9} = \frac{4 R e^{+k_i R}}{1 + 2 D_i (k_i + \frac{1}{R})} \quad \text{and}$$

$$C_{i10} = \frac{e^{-k_i R} [1 + 2 D_{i+1} (k_i + R)]}{e^{-k_{i+1} R} [1 + 2 D_{i+1} (k_{i+1} + R)]} G_{i13} \quad .$$

${}_i\beta_{i+2}$ ($i = 1, 2$)

$${}_i\beta_{i+2} = \frac{1}{2R} [C_{i11} e^{-k_{i+2}R} + G_{i14} e^{-k_i R} + G_{i15} e^{-k_{i+1}R}] \quad (D-7)$$

where

$$G_{i14} = - \frac{\Sigma_{S_{i+2}} C_{i9} + \Sigma_{S_{i+1, i+2}} G_{i13}}{D_{i+2} (k_1^2 - k_{i+2}^2)} ,$$

$$G_{i15} = - \frac{\Sigma_{S_{i+1, i+2}}}{D_{i+2} (k_{i+1}^2 - k_{i+2}^2)} C_{i10} , \text{ and}$$

$$C_{i11} = - \left[e^{-k_i R} \left[1 + 2 D_{i+2} (k_1 + \frac{1}{R}) \right] G_{i14} + e^{-k_{i+1} R} \left[1 + 2 D_{i+2} (k_{i+2} + \frac{1}{R}) \right] G_{i15} \right] /$$

$$\left[e^{-k_{i+2} R} \left[1 + 2 D_{i+2} (k_{i+2} + \frac{1}{R}) \right] \right] .$$

${}_i\beta_4$ ($i = 1$)

$${}_i\beta_4 = \frac{1}{2R} [C_{i12} e^{-k_4 R} + G_{i16} e^{-k_1 R} + G_{i17} e^{-k_2 R} + G_{i18} e^{-k_3 R}] \quad (D-8)$$

where

$$G_{i16} = \frac{- \Sigma_{S_{i4}} C_{i9} + \Sigma_{S_{24}} G_{i13} + \Sigma_{S_{34}} G_{i14}}{D_4 (k_1^2 - k_4^2)}$$

$$G_{117} = - \frac{\Sigma_{S_{24}} C_{110} + \Sigma_{S_{24}} G_{115}}{D_4 (k_2^2 - k_4^2)}$$

$$G_{118} = - \frac{\Sigma_{S_{24}} C_{111}}{D_4 (k_3^2 - k_4^2)} , \text{ and}$$

$$C_{112} = - \left[e^{-k_1 R} \left[1 + 2 D_4 \left(k_1 + \frac{1}{R} \right) \right] G_{116} + e^{-k_2 R} \left[1 + 2 D_4 \left(k_2 + \frac{1}{R} \right) \right] G_{117} + \right. \\ \left. e^{-k_3 R} \left[1 + 2 D_4 \left(k_3 + \frac{1}{R} \right) \right] G_{118} \right] / \left[e^{-k_4 R} \left[1 + 2 D_4 \left(k_4 + \frac{1}{R} \right) \right] \right] .$$

Finally, to illustrate the concepts developed in this appendix, three different BeO reflector thicknesses are employed in the albedo calculations. The impact of the reflector thickness on the performance of the reflector albedos can be seen in Table D-1.

TABLE D-1. EFFECTS OF THE BeO REFLECTOR THICKNESS VARIATION ON THE REFLECTOR ALBEDO VALUES USING THE FOUR-GROUP ALBEDO METHOD

${}_1\beta_1$	Reflector Thickness in cm		
	100	60	30
${}_1\beta_1$.28041	.28041	.28041
${}_2\beta_2$.30778	.30778	.30777
${}_1\beta_3$.18006	.18006	.11610
${}_3\beta_3$.18481	.18006	.14602
${}_2\beta_2$.62471	.62471	.62470
${}_2\beta_4$.15887	.15887	.15883
${}_2\beta_4$.13455	.13202	.11357
${}_3\beta_3$.68448	.68448	.68447
${}_3\beta_4$.23575	.23344	.21657
${}_4\beta_4$.93254	.93111	.92088

- Inner radius of the reflector = $R = 70.0$ cm.
- Reflector thickness = $H - R$.
- BeO four-group constants from Table 6-3.

The BeO reflector having a thickness of 100 cm behaves virtually as an infinite reflector. From Table D-1, it is seen that, as the BeO reflector thickness decreases, β_i ($i = 1, 2, 3, 4$) values decrease quicker than $\beta_{i'}$ ($i' \neq 4$) results. This expected behavior is a direct consequence of the fact that (1) fast neutrons require larger numbers of collisions with BeO nuclei to decrease the neutron energy to the thermal energy range and (2) thermal neutron transmission probability increases as BeO reflector becomes thinner.

The program ALB5 used for these numerical calculations is in Appendix G.

APPENDIX E FOUR-GROUP $P_{0,r}$ DERIVATION

Here the derivation of the partial principal nonleakage probabilities, $P_{0,r}$, is derived employing the consistent four-group diffusion theory method. Table 5-2 shows that all fission neutrons are essentially born in the first two energy groups. Thus, it is assumed $\chi_3 = 0$ and $\chi_4 = 0$. The subscript, r , is used to denote the reflector.

Four-Group Diffusion Equations for the Core

The characteristic equations for the core fluxes can be seen as follows

$$-D_1 \nabla^2 \phi_1 + \Sigma_{R1} \phi_1 = \frac{\chi_1}{k_{eff}} \left[\left(v_1 \Sigma_{f1} + \frac{2 \Sigma_{a,2n1}}{\chi_1} \right) \phi_1 + v_2 \Sigma_{f2} \phi_2 + v_3 \Sigma_{f3} \phi_3 + v_4 \Sigma_{f4} \phi_4 \right]$$

$$-D_2 \nabla^2 \phi_2 + \Sigma_{R2} \phi_2 - \Sigma_{S12} \phi_1 + \frac{\chi_2}{k_{eff}} (v_1 \Sigma_{f1} \phi_1 + v_2 \Sigma_{f2} \phi_2 + v_3 \Sigma_{f3} \phi_3 + v_4 \Sigma_{f4} \phi_4)$$

(E-1)

$$-D_3 \nabla^2 \phi_3 + \Sigma_{R3} \phi_3 - \Sigma_{S13} \phi_1 + \Sigma_{S23} \phi_2$$

$$-D_4 \nabla^2 \phi_4 + \Sigma_{R4} \phi_4 - \Sigma_{S14} \phi_1 + \Sigma_{S24} \phi_2 + \Sigma_{S34} \phi_3$$

Four-Group Diffusion Equations for the Reflector

The characteristic equations for the reflector fluxes can be written as follows

$$\begin{aligned}
 & -D_{1r} \nabla^2 \phi_{1r} + \Sigma_{R1r} \phi_{1r} = 0 \\
 & -D_{2r} \nabla^2 \phi_{2r} + \Sigma_{R2r} \phi_{2r} - \Sigma_{S12r} \phi_{1r} \\
 & -D_{3r} \nabla^2 \phi_{3r} + \Sigma_{R3r} \phi_{3r} - \Sigma_{S13r} \phi_{1r} + \Sigma_{S23r} \phi_{2r} \\
 & -D_{4r} \nabla^2 \phi_{4r} + \Sigma_{R4r} \phi_{4r} - \Sigma_{S14r} \phi_{1r} + \Sigma_{S24r} \phi_{2r} + \Sigma_{S34r} \phi_{3r}
 \end{aligned} \tag{E-2}$$

Boundary Conditions for the Core and Reflector Fluxes

The boundary conditions for the core and reflector fluxes are described below

- 1) at $r = 0$, ϕ_i are finite,
- 2) at $r = R$, $\phi_i = \phi_{ir}$ and $-D_i \phi_i' = -D_{ir} \phi_{ir}'$, i.e., neutron fluxes and currents must be continuous at the boundary between core and reflector. The prime (') denotes the gradient of the flux, and
- 3) at $r = \infty$, ϕ_{ir} are finite.

Core Flux Distributions, ϕ_i ($i = 1, 2, 3, 4$)

Solving the simultaneous linear differential equations with constant coefficients for the core and reflector fluxes and applying the boundary conditions, one can write the core flux solutions

$$\phi_1(r) = M_{11} \frac{\sin(\lambda_1 r)}{r} + M_{12} \frac{\sinh(\lambda_2 r)}{r} + \quad (E-3)$$

$$M_{13} \frac{\sinh(\lambda_3 r) \cos(\lambda_4 r)}{r} + M_{14} \frac{\cosh(\lambda_3 r) \sin(\lambda_4 r)}{r} .$$

The roots $\lambda_1, \lambda_2, \lambda_3, \lambda_4$ are obtained from the characteristic equations of the core fluxes. The coefficients M_{11}, M_{12}, M_{13} , and M_{14} are known from Equation E-1 combined with the boundary condition equations.

Evaluation of P_{o_i}

For a spherical core with an infinite reflector, the ϕ_{o_i} are given by

$$\phi_{o_i}(r) = \phi_i(r) - \frac{H_i}{r} \sinh(L_i r)$$

where

$$L_i = \sqrt{\frac{\Sigma_{R_i}}{D_i}} \quad \text{and subject to the boundary condition}$$

$$\text{at } r = R, \quad J_{o_i} = -\frac{\phi_{o_i}}{4} + \frac{D_i}{2} \frac{d}{dr} \phi_{o_i} = 0 .$$

After applying the boundary condition above, one can write the expression for H_1 as follows

$$\begin{aligned}
 H_1 = & \left[R \left[M_{11} \sin (\lambda_1 R) + M_{12} \sinh (\lambda_2 R) + M_{13} \sinh (\lambda_3 R) \cos (\lambda_4 R) + \right. \right. \\
 & M_{14} \cosh (\lambda_3 R) \sin (\lambda_4 R) \left. \right] + 2 D_1 \left[M_{11} \left[R \lambda_1 \cos (\lambda_1 R) - \sin (\lambda_1 R) \right] + \right. \\
 & M_{12} \left[R \lambda_2 \cosh (\lambda_2 R) - \sinh (\lambda_2 R) \right] + M_{13} R \left[\lambda_3 \cosh (\lambda_3 R) \cos (\lambda_4 R) - \right. \\
 & \left. \lambda_4 \sinh (\lambda_3 R) \sin (\lambda_4 R) - (\sinh (\lambda_3 R) \cos (\lambda_4 R)) \right] / R \left. \right] + \\
 & M_{14} R \left[\lambda_3 \sinh (\lambda_3 R) \sin (\lambda_4 R) + \lambda_4 \cosh (\lambda_3 R) \cos (\lambda_4 R) - \right. \\
 & \left. (\cosh (\lambda_3 R) \sin (\lambda_4 R)) / R \right] \left. \right] / \left[R \sinh (L_1 R) + 2 D_1 \left[R L_1 \cosh (L_1 R) - \right. \right. \\
 & \left. \left. \sinh (L_1 R) \right] \right] \quad . \quad (E-4)
 \end{aligned}$$

Finally, the initial nonleakage probabilities, P_{o_1} , are given by

$$P_{o_1} = \left[1 + \frac{B_1^2}{L_1^2} \right]^{-1}$$

where

$$\begin{aligned}
 B_1^2 = & - \frac{\int_S \nabla \phi_{o_1} \cdot d\vec{S}}{\int_V \phi_{o_1} dv} = - \frac{\int_{core} \nabla^2 \phi_{o_1} dv}{\int_{core} \phi_{o_1} dv} = \\
 = & - \frac{\int_{r=0}^R \frac{d^2}{dr^2} (r \phi_{o_1}) r dr}{\int_{r=0}^R \phi_{o_1} r^2 dr} \quad . \quad (E-5)
 \end{aligned}$$

APPENDIX F
THE PING-PONG DECISION PROCESS FOR A SOURCE
OF GROUP 4 NEUTRONS IN A TWO-REGION
GAS CORE REACTOR PROBLEM

Here is considered the contribution of one group 4 neutron source for evaluation of P_0 and P . The neutron fraction transmitted through the core region is δ_4 and the fraction returned as group 4 on this first core boundary traverse is $\delta_4 \beta_4$ as can be seen in Figure F-1.

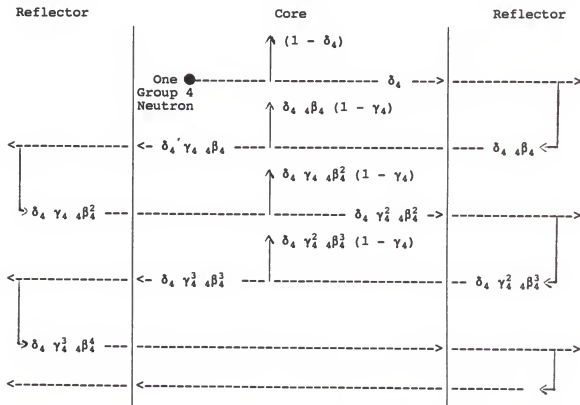


Figure F-1. Schematic drawing showing the ping-pong decision process for a core source of thermal neutrons in a two-region gas core reactor problem.

From Figure F-1, one can infer:

- (1) The fraction that interacts into the core is given by

$$H_4 = (1 - \delta_4) + \delta_4 \beta_4 (1 - \delta_4) \left[1 + \gamma_4 \beta_4 + (\gamma_4 \beta_4)^2 + \dots \right] \quad \text{or}$$

$$H_4 = (1 - \delta_4) + \frac{\delta_4 \beta_4 (1 - \gamma_4)}{1 - \gamma_4 \beta_4} , \quad (\text{F-1})$$

- (2) The initial fraction absorbed in the core as group 4 is given by

$$H_4 \times \frac{\Sigma_{a4}^{\text{core}}}{\Sigma_{t4}^{\text{core}}} , \quad (\text{F-2})$$

- (3) The fraction scattered in the core as group 4 is given by

$$H_4 \times \frac{\Sigma_{s44}^{\text{core}}}{\Sigma_{t4}^{\text{core}}} . \quad (\text{F-3})$$

Thus, in this first run, one initial group 4 neutron creates another secondary source of group 4 neutron whose strength is given by Equation F-3, and the fraction absorbed in the core as group 4 neutron is given by Equation F-2. Repeating this ping-pong decision process infinite times, i.e., until all secondary sources disappear, the total quantity absorbed as group 4 neutron is given by

$$H_4 \times \frac{\Sigma_{a_4}^{\text{core}}}{\Sigma_{c_4}^{\text{core}}} \left[1 + H_4 \frac{\Sigma_{a_{44}}^{\text{core}}}{\Sigma_{c_4}^{\text{core}}} + \left(H_4 \frac{\Sigma_{a_{44}}^{\text{core}}}{\Sigma_{c_4}^{\text{core}}} \right)^2 + \dots \right] \text{ or}$$

$$\frac{H_4 \times \frac{\Sigma_{a_4}^{\text{core}}}{\Sigma_{c_4}^{\text{core}}}}{1 - H_4 \times \frac{\Sigma_{a_{44}}^{\text{core}}}{\Sigma_{c_4}^{\text{core}}}} \quad . \quad (\text{F-4})$$

Finally, for each group 4 neutron created in the core the quantity given by Equation F-4 is absorbed in the core as group 4 neutron.

APPENDIX G FOUR-GROUP ALBEDO METHOD COMPUTER PROGRAMS

The four-group albedo approach computer programs used in this work are written in FORTRAN-77 and run on a personal computer. These programs are described below.

Source Listing of ALB1

```

C
C *****
C FOUR-GROUP ALBEDOS , NONLEAKAGE PROBABILITIES AND EFFECTIVE
C NEUTRON MULTIPLICATION FACTOR FOR A LARGE OPTICAL-PATH-
C LENGTH CORE REACTOR (TWO-REGION SOLID CORE REACTOR)
C *****
C
C PROGRAM ALB1
C NTIN=01
C NTOUT=02
C OPEN(NTIN, FILE='ALB1.IN', STATUS='OLD')
C OPEN(NTOUT, FILE='ALB1.OUT', STATUS='UNKNOWN')
C
C *****
C INPUT DATA
C
C R = RADIUS OF THE SPHERICAL CORE
C DF = CORE DENSITY FACTOR
C EK = STARTING GUESS FOR THE EFFECTIVE NEUTRON
C MULTIPLICATION FACTOR
C
C CORE GROUP PARAMETERS ( SEE TABLE 5-2 )
C
C D1 = DIFFUSION COEFFICIENT FOR GROUP 1
C D2 = DIFFUSION COEFFICIENT FOR GROUP 2 , AND SO ON.
C EN1 = MACROSCOPIC (n,2n) CROSS SECTION FOR GROUP 1
C EA1 = MACROSCOPIC ABSORPTION CROSS SECTION FOR GROUP 1
C EA2 = MACROSCOPIC ABSORPTION CROSS SECTION FOR GROUP 2 ,
C AND SO ON.
C FV1 = ( AVERAGE NUMBER OF FISSION NEUTRONS RELEASED IN A
C FISSION REACTION ) * ( MACROSCOPIC FISSION CROSS
C SECTION ) FOR GROUP 1
C FV2 = ( AVERAGE NUMBER OF FISSION NEUTRONS RELEASED IN A
C FISSION REACTION ) * ( MACROSCOPIC FISSION CROSS
C SECTION ) FOR GROUP 2 , AND SO ON.
C CHI1 = PROBABILITY THAT A FISSION NEUTRON WILL BE BORN
C WITH AN ENERGY IN GROUP 1
C CHI2 = PROBABILITY THAT A FISSION NEUTRON WILL BE BORN
C WITH AN ENERGY IN GROUP 2
C ES12 = MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C GROUP 1 TO GROUP 2
C ES13 = MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C GROUP 1 TO GROUP 3 , AND SO ON.
C
C REFLECTOR GROUP PARAMETERS ( SEE TABLE 5-3 )
C
C DR1 = DIFFUSION COEFFICIENT FOR GROUP 1
C DR2 = DIFFUSION COEFFICIENT FOR GROUP 2 , AND SO ON.
C ERA1 = MACROSCOPIC ABSORPTION CROSS SECTION FOR GROUP 1
C ERA2 = MACROSCOPIC ABSORPTION CROSS SECTION FOR GROUP 2 ,
C AND SO ON.

```

```

C   ERS12= MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C   GROUP 1 TO GROUP 2
C   ERS13= MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C   GROUP 1 TO GROUP 3 , AND SO ON.
C
C   *****
C
C   READ(01,1)R,DF,EK
C   READ(01,2)D1,D2,D3,D4,EN1,EA1,EA2,EA3,EA4,FV1,FV2,FV3,FV4,
C   #ES12,ES13,ES14,ES23,ES24,ES34,CHI1,CHI2,DR1,DR2,DR3,DR4,
C   #ER1,ER2,ER3,ER4,ERS12,ERS13,ERS14,ERS23,ERS24,ERS34
1   FORMAT(3E17.7)
2   FORMAT(4E17.7)
C   D1=D1/DF
C   D2=D2/DF
C   D3=D3/DF
C   D4=D4/DF
C   EA1=EA1*DF
C   EA2=EA2*DF
C   EA3=EA3*DF
C   EA4=EA4*DF
C   ES12=ES12*DF
C   ES13=ES13*DF
C   ES14=ES14*DF
C   ES23=ES23*DF
C   ES24=ES24*DF
C   ES34=ES34*DF
C   FV1=FV1*DF
C   FV2=FV2*DF
C   FV3=FV3*DF
C   FV4=FV4*DF
C   EN1=EN1*DF
C   EA1=EA1+EN1
C   ER1=EA1+ES12+ES13+ES14
C   ER2=EA2+ES23+ES24
C   ER3=EA3+ES34
C   ER4=EA4
C   ERR1=ER1+ERS12+ERS13+ERS14
C   ERR2=ER2+ERS23+ERS24
C   ERR3=ER3+ERS34
C   ERR4=ER4
C
C   *****
C   PARTIAL PRINCIPAL NOWLEAKAGE PROBABILITY CALCULATION
C   *****
C
C   CK1=SQRT(ER1/D1)
C   CK2=SQRT(ER2/D2)
C   CK3=SQRT(ER3/D3)
C   CK4=SQRT(ER4/D4)
C   RK1=SQRT(ERR1/DR1)
C   RK2=SQRT(ERR2/DR2)
C   RK3=SQRT(ERR3/DR3)
C   RK4=SQRT(ERR4/DR4)
C   Z1=EXP(-RK1*R)/R
C   Z2=EXP(-RK2*R)/R
C   Z3=EXP(-RK3*R)/R
C   Z4=EXP(-RK4*R)/R
C   ZP1=-(R*RK1+1.)*EXP(-RK1*R)/(R*R)
C   ZP2=-(R*RK2+1.)*EXP(-RK2*R)/(R*R)
C   ZP3=-(R*RK3+1.)*EXP(-RK3*R)/(R*R)
C   ZP4=-(R*RK4+1.)*EXP(-RK4*R)/(R*R)
C   F1=ERS12/(DR2*(RK2*RK2-RK1*RK1))
C   F2=(ERS13+ERS23*F1)/(DR3*(RK3*RK3-RK1*RK1))
C   F3=ERS23/(DR3*(RK3*RK3-RK2*RK2))
C   F4=(ERS14+ERS24*F1+ERS34*F2)/(DR4*(RK4*RK4-RK1*RK1))
C   F5=(ERS24+ERS34*F3)/(DR4*(RK4*RK4-RK2*RK2))
C   F6=ERS34/(DR4*(RK4*RK4-RK3*RK3))
C   Q11=((CHI1*FV1+2.*EN1)/(EK-ER1))/D1
C   Q12=CHI1*FV2/(EK*D1)
C   Q13=CHI1*FV3/(EK*D1)
C   Q14=CHI1*FV4/(EK*D1)
C   Q21=(ES12+CHI2*FV1/EK)/D2
C   Q22=(CHI2*FV2/EK-ER2)/D2
C   Q23=CHI2*FV3/(EK*D2)
C   Q24=CHI2*FV4/(EK*D2)
C   Q31=ES13/D3
C   Q32=ES23/D3
C   Q33=-ER3/D3

```

```

Q41=ES14/D4
Q42=ES24/D4
Q43=ES34/D4
Q44=FR4/D4
A1=Q11+Q22+Q33+Q44
A2=Q11*(Q22+Q33+Q44)+Q22*(Q33+Q44)+Q33*Q44-(Q14*Q41+Q24*
#Q42+Q13*Q31+Q23*Q32+Q12*Q21)
A3=Q14*Q41*(Q22+Q33)+Q14*Q21*Q42+Q14*Q31*Q43+Q24*Q32*Q43+
#Q24*Q12*Q41-Q24*Q42*(Q11+Q33)+Q44*(Q11*Q22+Q11*Q33+Q22*
#Q33-Q44*Q13*Q31-Q44*Q23*Q32-Q44*Q12*Q21+Q11*Q22*Q33+Q12*
#Q23*Q13*Q21+Q32*Q13-Q13*Q31*Q22-Q23*Q32*Q11-Q12*Q21*Q33
A4=Q14*(Q21*Q32+Q43+Q41*Q22+Q33+Q23*Q31*Q42-Q23*Q32*Q41-
#Q21*Q42*Q33-Q31*Q43*Q22)+Q24*(Q32*Q43*Q11+Q12*Q41*Q33+Q13*
#Q31*Q42-Q13*Q32*Q41-Q42*Q33*Q11-Q12*Q13*Q43)+Q44*(Q11*Q22*
#Q33+Q12*Q23*Q31+Q21*Q32*Q13-Q13*Q31*Q22-Q23*Q32*Q11-Q12*
#Q21*Q33)
PI=2.*ASIN(1.)
CA1=-A2
CA2=-A1*A3-4.*A4
CA3=-4.*A2*A4-A3*A3-A1*A1*A4
QP1=(3.*CA2-CA1*CA1)/9.
QP2=(9.*CA1*CA2-27.*CA3-2.*CA1**3)/54.
QP3=QP1**3+QP2*QP2
QP4=QP2+(ABS(QP3))**0.5
QP5=QP2-(ABS(QP3))**0.5
QP6=(ABS(QP4))** (1./3.)
QP7=(ABS(QP5))** (1./3.)
IF (QP3) 10,52,52
10  TETA=ACOS(QP2/((-QP1**3)**0.5))
R1=2.*(-QP1**0.5*COS(TETA/3.))-CA1/3.
R2=2.*(-QP1)**0.5*COS(TETA/3.+2.*PI/3.))-CA1/3.
R3=2.*(-QP1)**0.5*COS(TETA/3.+4.*PI/3.))-CA1/3.
DSR1=A1*A1-4.*A2+4.*R1
ESR1=R1*R1-4.*A4
IF (DSR1) 25,20,20
20  CONTINUE
IF (ESR1) 25,22,22
22  XR1=RL
GO TO 64
25  DSR2=A1*A1-4.*A2+4.*R2
ESR2=R2*R2-4.*A4
IF (DSR2) 35,30,30
30  CONTINUE
IF (ESR2) 35,32,32
32  XR1=R2
GO TO 64
35  DSR3=A1*A1-4.*A2+4.*R3
ESR3=R3*R3-4.*A4
IF (DSR3) 999,40,40
40  CONTINUE
IF (ESR3) 999,42,42
42  XR1=R3
GO TO 64
52  CONTINUE
IF (QP4) 54,56,56
54  QP6=QP6
GO TO 57
56  QP6=QP6
57  CONTINUE
IF (QP5) 58,60,60
58  QP7=QP7
GO TO 61
60  QP7=QP7
61  CONTINUE
XR1=QP6+QP7-CA1/3.
XRR=(QP6+QP7)/2.-CA1/3.
XIR=(QP6-QP7)*COS(PI/6.)
64  DIS=A1*A1-4.*A2+4.*XR1
EIS=XR1*XR1-4.*A4
IF (DIS) 999,66,66
66  CONTINUE
IF (EIS) 999,67,67
67  CONTINUE
QP8=(A1-SQRT(DIS))/2.
QP9=(XR1+SQRT(EIS))/2.
QP10=(A1+SQRT(DIS))/2.
QP11=(XR1-SQRT(EIS))/2.
DIS1=QP8*QP8-4.*QP9
QP12=(ABS(DIS1))**0.5

```

```

IF(DIS1)74,70,70
70 ZZ1=(-QP8+QP12)/2.
   ZZ2=(-QP8-QP12)/2.
   GO TO 81
74 RZZ1=-QP8/2.
   RIZ1=QP12/2.
   TETA1=ATAN(ABS(RIZ1/RZZ1))
   IF(RZZ1)77,78,78
77 AA=(RZZ1*RZZ1+RIZ1*RIZ1)**0.25*COS((PI-TETA1)/2.)
   BB=(RZZ1*RZZ1+RIZ1*RIZ1)**0.25*SIN((PI-TETA1)/2.)
   A=AA
   B=BB
   GO TO 81
78 AA=(RZZ1*RZZ1+RIZ1*RIZ1)**0.25*COS(TETA1/2.)
   BB=(RZZ1*RZZ1+RIZ1*RIZ1)**0.25*SIN(TETA1/2.)
   A=AA
   B=BB
81 DIS2=QP10*QP10-4.*QP11
   QP13=(ABS(DIS2))**0.5
   IF(DIS2)86,82,82
82 ZZ3=(-QP10+QP13)/2.
   ZZ4=(-QP10-QP13)/2.
   ZZ1=ZZ3
   ZZ2=ZZ4
   GO TO 99
86 RZZ3=-QP10/2.
   RIZ3=QP13/2.
   TETA2=ATAN(ABS(RIZ3/RZZ3))
   IF(RZZ3)90,91,91
90 A=(RZZ3*RZZ3+RIZ3*RIZ3)**0.25*COS((PI-TETA2)/2.)
   B=(RZZ3*RZZ3+RIZ3*RIZ3)**0.25*SIN((PI-TETA2)/2.)
   GO TO 99
91 A=(RZZ3*RZZ3+RIZ3*RIZ3)**0.25*COS(TETA2/2.)
   B=(RZZ3*RZZ3+RIZ3*RIZ3)**0.25*SIN(TETA2/2.)
99 CONTINUE
   SLAMB=ZZ1
   SMU=ZZ2
   SU=SQRT(SMU)
   SL=SQRT(SLAMB)
   X1=SIN(SU*R)/R
   X2= SINH(SL*R)/R
   X3= SINH(A*R)*COS(B*R)/R
   X4= COSH(A*R)*SIN(B*R)/R
   XP1=(R*SU*COS(SU*R)-SIN(SU*R))/(R*R)
   XP2=(R*SL*COSH(SL*R)-SINH(SL*R))/(R*R)
   XP3=(R*A*COSH(A*R)*COS(B*R)-R*B*SINH(A*R)*SIN(B*R)-
#SINH(A*R)*COS(B*R))/(R*R)
   XP4=(R*A*SINH(A*R)*SIN(B*R)+R*B*COSH(A*R)*COS(B*R)-
#COSH(A*R)*SIN(B*R))/(R*R)
   Q34=Q24*(Q43*Q32-Q42*(Q33-SMU))-(Q44-SMU)*(Q32*Q23-(Q22-
#SMU)*(Q33-SMU))
   Q50=(Q24*(Q31*Q42-Q32*Q41)-(Q44-SMU)*(Q31*(Q22-SMU)-
#Q32*Q21))/Q34
   Q51=(-Q31-(Q33-SMU)*Q50)/Q32
   Q52=(-Q41-Q43*Q50-Q42*Q51)/(Q44-SMU)
   Q35=Q24*(Q32*Q43-Q42*(Q33+SL*SL))+(Q44+SL*SL)*{(Q22+SL*SL)
#(Q33+SL*SL)-Q23*Q32}
   Q53=(Q24*(Q31*Q42-Q32*Q41)+(Q44+SL*SL)*(Q21*Q32-Q31*(Q22+
#SL*SL)))/Q35
   Q54=(-Q31-(Q33+SL*SL)*Q53)/Q32
   Q55=(-Q41-Q43*Q53-Q42*Q54)/(Q44+SL*SL)
   Q60=2.*A*B
   Q61=Q11+A*A-B*B
   Q62=Q33+A*A-B*B
   Q63=Q44+A*A-B*B
   Q64=Q42-Q63*Q12/Q14
   Q65=-Q60*Q12/Q14
   Q66=Q43-Q63*Q13/Q14
   Q67=-Q60*Q13/Q14
   Q68=-Q63*Q61/Q14+Q60*Q60/Q14+Q41
   Q69=-Q63*Q60/Q14-Q60*Q61/Q14
   Q70=Q60*Q12/Q14
   Q71=Q42-Q63*Q12/Q14
   Q72=Q60*Q13/Q14
   Q73=Q43-Q63*Q13/Q14
   Q75=Q60*(Q61+Q63)/Q14
   Q76=Q60*Q60/Q14-Q63*Q61/Q14+Q41
   Q77=-Q64*Q62/Q32+Q65*Q60/Q32+Q66
   Q78=-Q64*Q60/Q32-Q65*Q62/Q32+Q67

```

Q79=Q64*Q31/Q32+Q68
 Q80=Q65*Q31/Q32+Q69
 Q81=Q70*Q62/Q32+Q71*Q60/Q32+Q72
 Q82=Q70*Q60/Q32-Q71*Q62/Q32+Q73
 Q83=Q70*Q31/Q32+Q75
 Q84=Q71*Q31/Q32+Q76
 Q36=Q77*Q82-Q78*Q81
 Q85=(Q78*Q83-Q79*Q82)/Q36
 Q86=(Q78*Q84-Q80*Q82)/Q36
 Q87=(Q79*Q81-Q77*Q83)/Q36
 Q88=(Q80*Q81-Q77*Q84)/Q36
 Q89=(Q31+Q62*Q85+Q60*Q87)/Q32
 Q90=(Q62*Q86+Q60*Q88)/Q32
 Q91=(Q60*Q85-Q62*Q87)/Q32
 Q92=(-Q31+Q60*Q86-Q62*Q88)/Q32
 Q93=(Q61+Q12*Q89+Q13*Q85)/Q14
 Q94=(Q60+Q12*Q90+Q13*Q86)/Q14
 Q95=(Q60-Q12*Q91-Q13*Q87)/Q14
 Q96=(Q61+Q12*Q92+Q13*Q88)/Q14
 Q26=A*A-B*B+Q22
 Q100=Q26*Q43/Q42+Q23
 Q101=-Q60*Q43/Q42
 Q102=Q26*Q63/Q42+Q24+Q60*Q60/Q42
 Q103=(Q26+Q63)*Q60/Q42
 Q104=Q26*Q41/Q42-Q21
 Q105=Q60*Q41/Q42
 Q106=Q60*Q43/Q42
 Q107=Q26*Q43/Q42+Q23
 Q108=(Q63+Q26)*Q60/Q42
 Q109=Q60*Q60/Q42-Q26*Q63/Q42+Q24
 Q110=Q60*Q41/Q42
 Q111=Q26*Q41/Q42-Q21
 Q112=Q62/Q32+Q43/Q42
 Q113=Q60/Q32
 Q114=Q63/Q42
 Q115=Q60/Q42
 Q116=Q41/Q42+Q31/Q32
 Q117=Q60/Q32
 Q118=Q62/Q32+Q43/Q42
 Q119=Q60/Q42
 Q120=Q63/Q42
 Q121=Q41/Q42+Q31/Q32
 Q122=Q112-Q115*Q117/Q120
 Q123=Q113-Q115*Q118/Q120
 Q124=Q114-Q115*Q119/Q120
 Q125=Q115*Q121/Q120
 Q126=Q106-Q108*Q122/Q124-Q109*Q117/Q120
 Q127=Q107-Q108*Q123/Q124-Q109*Q118/Q120
 Q128=Q109*Q119/Q120
 Q129=Q110-Q108*Q116/Q124
 Q130=Q111-Q108*Q125/Q124-Q109*Q121/Q120
 Q131=Q126-Q128*Q122/Q124
 Q132=Q127-Q128*Q123/Q124
 Q133=Q129-Q128*Q116/Q124
 Q134=Q130-Q128*Q125/Q124
 Q135=Q100-Q102*Q122/Q124-Q103*Q117/Q120
 Q136=Q101-Q102*Q123/Q124-Q103*Q118/Q120
 Q137=Q103*Q119/Q120
 Q138=Q104-Q102*Q116/Q124
 Q139=Q105-Q103*Q121/Q120-Q102*Q125/Q124
 Q140=Q135-Q137*Q122/Q124
 Q141=Q136-Q137*Q123/Q124
 Q142=Q138-Q137*Q116/Q124
 Q143=Q139-Q137*Q125/Q124
 Q144=Q131*Q141-Q132*Q140
 Q145=(Q141*Q133-Q132*Q142)/Q144
 Q146=(Q141*Q134-Q132*Q143)/Q144
 Q147=(Q131*Q142-Q133*Q140)/Q144
 Q148=(Q131*Q143-Q134*Q140)/Q144
 Q149=(Q31+Q62*Q145+Q60*Q147)/Q32
 Q150=(Q62*Q146+Q60*Q148)/Q32
 Q151=(Q60*Q145-Q62*Q147)/Q32
 Q152=(Q60*Q146-Q31-Q62*Q148)/Q32
 Q153=(Q116-Q122*Q145-Q123*Q147)/Q124
 Q154=(Q125-Q122*Q146-Q123*Q148)/Q124
 Q155=(Q117*Q145+Q118*Q147+Q119*Q153)/Q120
 Q156=(Q121-Q117*Q146-Q118*Q148-Q119*Q154)/Q120
 S11=Q51
 S22=Q54

S33-Q149
 S34-Q150
 S43-Q151
 S44-Q152
 AM11-Q50
 AM22-Q53
 AM33-Q145
 AM34-Q146
 AM43-Q147
 AM44-Q148
 AN11-Q52
 AN22-Q55
 AN33-Q153
 AN34-Q154
 AN43-Q155
 AN44-Q156
 Y11=X1
 Y12=X2
 Y13=X3
 Y14=X4
 Y15=Z1
 Y21=S11*X1
 Y22=S22*X2
 Y23=S33*X3+S43*X4
 Y24=S34*X3+S44*X4
 Y25=F1*Z1
 Y26=Z2
 Y31=AM11*X1
 Y32=AM22*X2
 Y33=AM33*X3+AM43*X4
 Y34=AM34*X3+AM44*X4
 Y35=F2*Z1
 Y36=F3*Z2
 Y37=Z3
 Y41=AN11*X1
 Y42=AN22*X2
 Y43=AN33*X3+AN43*X4
 Y44=AN34*X3+AN44*X4
 Y45=F4*Z1
 Y46=F5*Z2
 Y47=F6*Z3
 Y48=Z4
 Y51=XP1
 Y52=XP2
 Y53=XP3
 Y54=XP4
 Y55=DR1*ZP1/D1
 Y61=S11*XP1
 Y62=S22*XP2
 Y63=S33*XP3+S43*XP4
 Y64=S34*XP3+S44*XP4
 Y65=DR2*F1*ZP1/D2
 Y66=DR2*ZP2/D2
 Y71=AM11*XP1
 Y72=AM22*XP2
 Y73=AM33*XP3+AM43*XP4
 Y74=AM34*XP3+AM44*XP4
 Y75=DR3*F2*ZP1/D3
 Y76=DR3*F3*ZP2/D3
 Y77=DR3*ZP3/D3
 X60=Y31-Y35*Y11/Y15-Y36*(Y21-Y25*Y11/Y15)/Y26
 X61=Y32-Y35*Y12/Y15-Y36*(Y22-Y25*Y12/Y15)/Y26
 X62=Y33-Y35*Y13/Y15-Y36*(Y23-Y25*Y13/Y15)/Y26
 X63=Y34-Y35*Y14/Y15-Y36*(Y24-Y25*Y14/Y15)/Y26
 X65=Y71-Y75*Y11/Y15-Y76*(Y21-Y25*Y11/Y15)/Y26
 X66=Y72-Y75*Y12/Y15-Y76*(Y22-Y25*Y12/Y15)/Y26
 X67=Y73-Y75*Y13/Y15-Y76*(Y23-Y25*Y13/Y15)/Y26
 X68=Y74-Y75*Y14/Y15-Y76*(Y24-Y25*Y14/Y15)/Y26
 X50=Y11*Y55-Y15*Y51
 X51=Y12*Y55-Y15*Y52
 X52=Y13*Y55-Y15*Y53
 X53=Y14*Y55-Y15*Y54
 X55=Y66*(Y21*Y15-Y25*Y11)-Y26*(Y61*Y15-Y65*Y11)
 X56=Y66*(Y22*Y15-Y25*Y12)-Y26*(Y62*Y15-Y65*Y12)
 X57=Y66*(Y23*Y15-Y25*Y13)-Y26*(Y63*Y15-Y65*Y13)
 X58=Y66*(Y24*Y15-Y25*Y14)-Y26*(Y64*Y15-Y65*Y14)
 X70=Y77*X60-Y37*X65
 X71=Y77*X61-Y37*X66
 X72=Y77*X62-Y37*X67

```

X73=Y77*X63-Y37*X68
DELT=X51*(X57*X73-X58*X72)-X52*(X56*X73-X58*X71)+X53*(X56*
X72-X57*X71)
DELT2=X50*(X57*X73-X58*X72)-X52*(X58*X70-X55*X73)+X53*
(X57*X70-X55*X72)
DELT3=X51*(X58*X70-X55*X73)+X50*(X56*X73-X58*X71)+X53*(X55*
X71-X56*X70)
DELT4=X51*(X55*X72-X57*X70)-X52*(X55*X71-X56*X70)-X50*(X56*
X72-X57*X71)
Q2R=DELT2/DELT
Q3R=DELT3/DELT
Q4R=DELT4/DELT
V=A*A*B*B
V1=A*A-B*B
V2=2.*A*B
V3=SINH(A*R)
V4=COSH(A*R)
V5=SIN(B*R)
V6=COS(B*R)
V7=SINH(CK1*R)
V8=COSH(CK1*R)
V9=SINH(CK2*R)
V10=COSH(CK2*R)
V11=SINH(CK3*R)
V12=COSH(CK3*R)
V13=SINH(CK4*R)
V14=COSH(CK4*R)
AA1=V7/R+2.*D1*(R*CK1*V8-V7)/(R*R)
A1R=(X1+2.*D1*XP1+Q2R*(X2+2.*D1*XP2)+Q3R*(X3+2.*D1*XP3)+
Q4R*(X4+2.*D1*XP4))/AA1
W1=SIN(SU*R)/(SU*SU)-R*COS(SU*R)/SU
W2=R*COSH(SL*R)/SL-SINH(SL*R)/(SL*SL)
W3=(R*A*V4+V6+R*B*V3+V5)/V-(A*A*V3+V6+A*B*V4+V5)/V**2-
(B*A*V4+V5-B*B*V3+V6)/V**2
W4=(R*A*V3+V5-R*B*V4+V6)/V-(A*A*V4+V5-R*B*V3+V6)/V**2+
(B*A*V3+V6+B*B*V4+V5)/V**2
WN1=R*V8/CK1-V7/(CK1*CK1)
BUC1=(-(SU*SU*W1+Q2R*SL*SL*W2+Q3R*(V1*W3-V2*W4)+Q4R*(
V1*W4+V2*W3)-A1R*CK1*CK1*WN1)/(W1+Q2R*W2+Q3R*W3+Q4R*W4-
A1R*WN1)
FO1=-1./(1.+BUC1/(CK1*CK1))
AA2=V9/R+2.*D2*(R*CK2*V10-V9)/(R*R)
A2R=(S11*(X1+2.*D2*XP1)+S22*Q2R*(X2+2.*D2*XP2)+(S33*Q3R+
S34*Q4R)*(X3+2.*D2*XP3)+(S43*Q3R+S44*Q4R)*(X4+2.*D2*XP4))
/AA2
WN2=R*V10/CK2-V9/(CK2*CK2)
BUC2=(-(S11*SU*SU*W1+S22*Q2R*SL*SL*W2+(S33*Q3R+S34*Q4R)*(
V1*W3-V2*W4)+(S43*Q3R+S44*Q4R)*(V1*W4+V2*W3)-A2R*
CK2*CK2*WN2)/(S11*W1+S22*Q2R*W2+(S33*Q3R+S34*Q4R)*W3+(S43*
Q3R+S44*Q4R)*W4-A2R*WN2)
FO2=-1./(1.+BUC2/(CK2*CK2))
AA3=V11/R+2.*D3*(R*CK3*V12-V11)/(R*R)
A3R=(AM11*(X1+2.*D3*XP1)+AM22*Q2R*(X2+2.*D3*XP2)+(AM33*
Q3R+AM34*Q4R)*(X3+2.*D3*XP3)+(AM43*Q3R+AM44*Q4R)*(X4+2.*
D3*XP4))/AA3
WN3=R*V12/CK3-V11/(CK3*CK3)
BUC3=(-(AM11*SU*SU*W1+AM22*Q2R*SL*SL*W2+(AM33*Q3R+AM34*
Q4R)*(V1*W3-V2*W4)+(AM43*Q3R+AM44*Q4R)*(V1*W4+V2*W3)-
A3R*CK3*CK3*WN3)/(AM11*W1+AM22*Q2R*W2+(AM33*Q3R+AM34*Q4R)*
W3+(AM43*Q3R+AM44*Q4R)*W4-A3R*WN3)
FO3=-1./(1.+BUC3/(CK3*CK3))
AA4=V13/R+2.*D4*(R*CK4*V14-V13)/(R*R)
A4R=(AN11*(X1+2.*D4*XP1)+AN22*Q2R*(X2+2.*D4*XP2)+(AN33*
Q3R+AN34*Q4R)*(X3+2.*D4*XP3)+(AN43*Q3R+AN44*Q4R)*(X4+2.*
D4*XP4))/AA4
WN4=R*V14/CK4-V13/(CK4*CK4)
BUC4=(-(AN11*SU*SU*W1+AN22*Q2R*SL*SL*W2+(AN33*Q3R+AN34*
Q4R)*(V1*W3-V2*W4)+(AN43*Q3R+AN44*Q4R)*(V1*W4+V2*W3)-
A4R*CK4*CK4*WN4)/(AN11*W1+AN22*Q2R*W2+(AN33*Q3R+AN34*Q4R)*
W3+(AN43*Q3R+AN44*Q4R)*W4-A4R*WN4)
FO4=-1./(1.+BUC4/(CK4*CK4))

*****
CORE ALBEDO CALCULATION
*****

G11=CK1/TANH(CK1*R)-1./R
G12=CK2/TANH(CK2*R)-1./R
G13=CK3/TANH(CK3*R)-1./R

```

C
C
C
C
C

```

G14=CK4/TANH(CK4*R)-1./R
CA11=(1.-2.*D1*G11)/(1.+2.*D1*G11)
CA22=(1.-2.*D2*G12)/(1.+2.*D2*G12)
CA33=(1.-2.*D3*G13)/(1.+2.*D3*G13)
CA44=(1.-2.*D4*G14)/(1.+2.*D4*G14)
G15=4.*R/((1.+2.*D1*G11)*SINH(CK1*R))
G16=G15*RS12/(D2*(CK2*CK2-CK1*CK1))
G17=G15*SINH(CK1*R)*(1.+2.*D2*G11)/(SINH(CK2*R)*(1.+2.*
#D2*G12))
G18=(G15*RS13+G16*RS23)/(D3*(CK3*CK3-CK1*CK1))
G19=(G17*RS23)/(D3*(CK3*CK3-CK2*CK2))
G20=(G18*SINH(CK1*R)*(1.+2.*D3*G11)+G19*SINH(CK2*R)*(1.+
#2.*D3*G12))/(SINH(CK3*R)*(1.+2.*D3*G13))
G21=(G15*RS14+G16*RS24+G18*RS34)/(D4*(CK4*CK4-CK1*CK1))
G22=(G17*RS24+G19*RS34)/(D4*(CK4*CK4-CK1*CK1))
G23=G20*RS34/(D4*(CK4*CK4-CK1*CK1))
G24=(G21*V7*(1.+2.*D4*G11)+G22*V9*(1.+2.*D4*
#G12)+G23*V11*(1.+2.*D4*G13))/(V13*(1.+2.*D4*G14))
CA12=(G17*SINH(CK2*R)+G16*SINH(CK1*R))/(2.*R)
CA13=(G20*V11+G18*V7+G19*V9)/(2.*R)
CA14=(G24*V13+G21*V7+G22*V9+G23*V11)/(2.*R)
G25=4.*R/((1.+2.*D2*G12)*SINH(CK2*R))
G26=G25*RS23/(D3*(CK3*CK3-CK2*CK2))
G27=(G26*V9*(1.+2.*D3*G12))/(V11*(1.+2.*D3*G13))
G28=(G25*RS24+G26*RS34)/(D4*(CK4*CK4-CK2*CK2))
G29=G27*RS34/(D4*(CK4*CK4-CK3*CK3))
G30=(G28*V9*(1.+2.*D4*G12)+G29*V11*(1.+2.*D4*G13))/
#(V13*(1.+2.*D4*G14))
CA23=(G27*SINH(CK3*R)+G26*SINH(CK2*R))/(2.*R)
CA24=(G30*V13+G28*V9+G29*V11)/(2.*R)
G31=4.*R/((1.+2.*D3*G13)*SINH(CK3*R))
G32=G31*RS34/(D4*(CK4*CK4-CK3*CK3))
G33=G32*V11*(1.+2.*D4*G13)/(V13*(1.+2.*D4*G14))
CA34=(G33*V13+G32*V11)/(2.*R)

```

```

*****
REFLECTOR ALBEDO CALCULATION
*****

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```

G34=RK1+1./R
G35=RK2+1./R
G36=RK3+1./R
G37=RK4+1./R
RA11=(1.-2.*DR1*G34)/(1.+2.*DR1*G34)
RA22=(1.-2.*DR2*G35)/(1.+2.*DR2*G35)
RA33=(1.-2.*DR3*G36)/(1.+2.*DR3*G36)
RA44=(1.-2.*DR4*G37)/(1.+2.*DR4*G37)
G38=4./Z1*(1.+2.*DR1*G34)
G39=G38*RS12/(DR2*(RK2*RK2-RK1*RK1))
G40=G39*Z1*(1.+2.*DR2*G34)/(Z2*(1.+2.*DR2*G35))
G41=(G38*RS13+G39*RS23)/(DR3*(RK3*RK3-RK1*RK1))
G42=G40*RS23/(DR3*(RK3*RK3-RK2*RK2))
G43=(G41*Z1*(1.+2.*DR3*G34)+G42*Z2*(1.+2.*DR3*G35))/(Z3*(
#1.+2.*DR3*G36))
G44=(G38*RS14+G39*RS24+G41*RS34)/(DR4*(RK4*RK4-RK1*RK1))
G45=(G40*RS24+G42*RS34)/(DR4*(RK4*RK4-RK2*RK2))
G46=G43*RS34/(DR4*(RK4*RK4-RK3*RK3))
G47=(G44*Z1*(1.+2.*DR4*G34)+G45*Z2*(1.+2.*DR4*G35)+G46*Z3
#(1.+2.*DR4*G36))/(Z4*(1.+2.*DR4*G37))
RA12=(G40*Z2+G39*Z1)/2.
RA13=(G43*Z3+G41*Z1+G42*Z2)/2.
RA14=(G47*Z4+G44*Z1+G45*Z2+G46*Z3)/2.
G48=4./Z2*(1.+2.*DR2*G35)
G49=G48*RS23/(DR3*(RK3*RK3-RK2*RK2))
G50=G49*Z2*(1.+2.*DR3*G35)/(Z3*(1.+2.*DR3*G36))
G51=(G48*RS24+G49*RS34)/(DR4*(RK4*RK4-RK2*RK2))
G52=G50*RS34/(DR4*(RK4*RK4-RK3*RK3))
G53=(G51*Z2*(1.+2.*DR4*G35)+G52*Z3*(1.+2.*DR4*G36))/(Z4
#(1.+2.*DR4*G37))
RA23=(G50*Z3+G49*Z2)/2.
RA24=(G53*Z4+G51*Z2+G52*Z3)/2.
G54=4./Z3*(1.+2.*DR3*G36)
G55=G54*RS34/(DR4*(RK4*RK4-RK3*RK3))
G56=G55*Z3*(1.+2.*DR4*G36)/(Z4*(1.+2.*DR4*G37))
RA34=(G56*Z4+G55*Z3)/2.

```

```

*****
SOURCE OF INITIAL LEAKAGE CALCULATION
*****

```



```

C      S11=CHI1*(1.-PO1)
      S61=CHI2+CHI1*PO1*ES12/ER1
      S62=CHI1*PO1*ES13/ER1+S61*PO2*ES23/ER2
      S63=CHI1*PO1*ES14/ER1+S61*PO2*ES24/ER2+S62*PO3*ES34/ER3
      S12=S61*(1.-PO2)
      S13=S62*(1.-PO3)
      S14=S63*(1.-PO4)

C      *****
C      FRACTIONAL SECONDARY NONLEAKAGE PROBABILITY CALCULATION
C      *****

      G57=1.-CA11*RA11
      G58=1.-CA22*RA22
      G59=1.-CA33*RA33
      G60=1.-CA44*RA44
      PR44=RA44*(1.-CA44)/G60
      P4=SI4*PR44
      PR34=RA34*(1.-CA44)/(G59*G60)
      PR33=RA33*(1.-CA33)/G59-RA33*CA34*(1.-RA44)/(G59*G60)
      P3=SI3*(PR33+PR34)
      PR24=RA24*(1.-CA44)/(G58*G60)
      G61=(1.-CA33)/G59+CA33*RA34*(1.-CA44)/(G59*G60)
      G62=CA34*(1.-RA44)/(G59*G60)
      PR23=RA23*(G61-G62)/G58
      G63=RA22*(1.-CA22)/G58
      G64=RA22*CA23*(1.-P3/SI3)/G58
      G65=RA22*CA24*(1.-PR44)/G58
      PR22=G63-G64-G65
      P2=SI2*(PR24+PR23+PR22)
      PR14=RA14*(1.-CA44)/(G57*G60)
      PR13=RA13*(G61-G62)/G57
      G66=(1.-CA22)*G59*G60-CA24*(1.-RA44)*G59+CA22*RA24*(1.-
      #CA44)*G59
      G67=CA23*(1.-RA33)*G60+RA33*CA34*(1.-RA44)-RA34*(1.-CA44)
      G68=CA22*RA23*(1.-CA33)*G60+CA33*RA34*(1.-CA44)-CA34*(1.-
      #RA44)
      PR12=RA12*(G66-G67+G68)/(G57*G58*G59*G60)
      PR11=RA11*(1.-CA11)/G57-RA11*CA14*(1.-P4/SI4)/G57-RA11*
      #CA13*(1.-P3/SI3)/G57-RA11*CA12*(1.-P2/SI2)/G57
      P1=SI1*(PR14+PR13+PR12+PR11)
      PNP=1.-(S11+S12+S13+S14)
      SNP=P1+P2+P3+P4
      TNP=PNP+SNP

C      *****
C      CALCULATION OF THE NEUTRON ABSORPTION BY THE CORE
C      *****

      POB1=CHI1*PO1*EA1/ER1
      POB2=S61*PO2*EA2/ER2
      POB3=S62*PO3*EA3/ER3
      POB4=S63*PO4
      BPO=POB1+POB2+POB3+POB4
      PA34=(RA34*PR44/RA44+RA33*CA34*PR44)/G59
      PA33=RA33*(1.-CA34-CA33)/G59
      TA34=(CA34*PR44+CA33*RA34*PR44/RA44)/G59
      TA33=(1.-CA34-CA33)/G59
      PA24=(RA22*CA24*PR44+RA24*PR44/RA44+RA22*CA23*PA34+
      #RA23*TA34)/G58
      EA23=(RA22*CA23*PA33+RA23*TA33)/G58
      PA22=RA22*(1.-CA22-CA23-CA24)/G58
      TA24=(CA24*PR44+CA22*RA24*PR44/RA44+CA23*PA34+CA22*
      #RA23*TA34)/G58
      TA23=(CA23*PA33+CA22*RA23*TA33)/G58
      TA22=(1.-CA22-CA23-CA24)/G58
      PA14=(RA11*CA14*PR44+RA14*PR44/RA44+RA11*CA13*PA34+
      #RA13*TA34+RA11*CA12*PA24+RA12*TA24)/G57
      PA13=(RA11*CA13*PA33+RA13*TA33+RA11*CA12*PA23+
      #RA12*TA23)/G57
      PA12=(RA11*CA12*PA22+RA12*TA22)/G57
      PA11=RA11*(1.-CA11-CA12-CA13-CA14)/G57
      FF34=ES34/ER3
      FF33=EA3/ER3
      FF24=ES24/ER2+ES23*ES34/(ER2*ER3)
      FF23=ES23*EA3/(ER2*ER3)
      FF22=EA2/ER2
      FF14=ES14/ER1+ES12*(ES24+ES23*ES34/ER3)/(ER1*ER2)+

```

```

#ES13=ES34/(ER1*ER3)
FF13=ES13*EA3/(ER1*ER3)+ES12*ES23*EA3/(ER1*ER2*ER3)
FF12=ES12*EA2/(ER1*ER2)
FF11=EA1/ER1
E1A1=PA11*FF11
E1A2=PA11*FF12+PA12*FF22
E1A3=PA11*FF13+PA12*FF23+PA13*FF33
E1A4=PA11*FF14+PA12*FF24+PA13*FF34+PA14
E2A2=PA22*FF22
E2A3=PA22*FF23+PA23*FF33
E2A4=PA22*FF24+PA23*FF34+PA24
E3A3=PA33*FF33
E3A4=PA33*FF34+PA34
E4A4=PR44
PPB1=SI1*E1A1
PPB2=SI1*E1A2+SI2*E2A2
PPB3=SI1*E1A3+SI2*E2A3+SI3*E3A3
PPB4=SI1*E1A4+SI2*E2A4+SI3*E3A4+SI4*E4A4
BPP=PPB1+PPB2+PPB3+PPB4
TAB1=POB1+PPB1
TAB2=POB2+PPB2
TAB3=POB3+PPB3
TAB4=POB4+PPB4
BTNL=TAB1+TAB2+TAB3+TAB4

```

```

*****
EFFECTIVE NEUTRON MULTIPLICATION FACTOR CALCULATION
*****

```

```

#KEFF=TAB1*(FV1+2.*EN1)/EA1+TAB2*FV2/EA2+TAB3*FV3/EA3+TAB4*
#FV4/EA4

```

***** OUTPUT RESULTS

CORE ALBEDOS (SEE TABLE 5-4)

CA11 = PROBABILITY THAT A NEUTRON FROM GROUP 1 BE REFLECTED
WITH GROUP 1 ENERGY
CA12 = PROBABILITY THAT A NEUTRON FROM GROUP 1 BE REFLECTED
WITH GROUP 2 ENERGY , AND SO ON.

REFLECTOR ALBEDOS (SEE TABLE 5-5)

RA11 = PROBABILITY THAT A NEUTRON FROM GROUP 1 BE REFLECTED
WITH GROUP 1 ENERGY
RA12 = PROBABILITY THAT A NEUTRON FROM GROUP 1 BE REFLECTED
WITH GROUP 2 ENERGY , AND SO ON.

PARTIAL PRINCIPAL NONLEAKAGE PROBABILITIES (SEE
TABLE 5-6)

PO1 = PROBABILITY THAT GROUP 1 NEUTRONS NEVER SPEND TIME
IN THE REFLECTOR
PO2 = PROBABILITY THAT GROUP 2 NEUTRONS NEVER SPEND TIME
IN THE REFLECTOR , AND SO ON.

SOURCES OF INITIAL LEAKAGE (SEE TABLE 5-6)

SI1 = FRACTION OF GROUP 1 NEUTRONS THAT LEAKS FROM THE
CORE FOR THE FIRST TIME
SI2 = FRACTION OF GROUP 2 NEUTRONS THAT LEAKS FROM THE
CORE FOR THE FIRST TIME , AND SO ON.

FRACTIONAL SECONDARY NONLEAKAGE PROBABILITIES (SEE
TABLE 5-7)

PR11 = PROBABILITY THAT IF ONE NEUTRON LEAKS AS A GROUP
1 NEUTRON FROM THE CORE , AND RETURNS AS A GROUP
1 NEUTRON , IT WILL BE ABSORBED IN THE CORE
PR12 = PROBABILITY THAT IF ONE NEUTRON LEAKS AS A GROUP
1 NEUTRON FROM THE CORE , AND RETURNS AS A GROUP
2 NEUTRON , IT WILL BE ABSORBED IN THE CORE ,
AND SO ON.

NEUTRON ABSORPTION (SEE TABLE 5-8)

POB1 = FRACTION OF FISSION NEUTRONS (NEVER SPEND TIME
IN THE REFLECTOR) WHICH ARE ABSORBED BY THE

```

C      CORE AS GROUP 1
C      POB2 = FRACTION OF FISSION NEUTRONS ( NEVER SPEND TIME
C      IN THE REFLECTOR ) WHICH ARE ABSORBED BY THE
C      CORE AS GROUP 2 , AND SO ON.
C
C      PPB1 = FRACTION OF FISSION NEUTRONS ( SPEND TIME IN
C      THE REFLECTOR ) WHICH ARE ABSORBED BY THE
C      CORE AS GROUP 1
C      PPB2 = FRACTION OF FISSION NEUTRONS ( SPEND TIME IN
C      THE REFLECTOR ) WHICH ARE ABSORBED BY THE
C      CORE AS GROUP 2 , AND SO ON.
C
C      EFFECTIVE NEUTRON MULTIPLICATION FACTOR ( SEE
C      TABLE 5-10 )
C
C      EKEF = EFFECTIVE NEUTRON MULTIPLICATION FACTOR
C      *****
C
C      WRITE(02,100)R,DF
C      WRITE(02,102)CA11,CA12,CA13,CA14
C      WRITE(02,104)CA22,CA23,CA24
C      WRITE(02,106)CA33,CA34,CA44
C      WRITE(02,108)RA11,RA12,RA13,RA14
C      WRITE(02,110)RA22,RA23,RA24
C      WRITE(02,112)RA33,RA34,RA44
C      WRITE(02,114)PO1,PO2,PO3,PO4
C      WRITE(02,116)SI1,SI2,SI3,SI4
C      WRITE(02,118)PR11,PR12,PR13,PR14
C      WRITE(02,120)PR22,PR23,PR24
C      WRITE(02,122)PR33,PR34,PR44
C      WRITE(02,124)POB1,POB2,POB3,POB4
C      WRITE(02,126)PPB1,PPB2,PPB3,PPB4
C      WRITE(02,128)EKEF
100  FORMAT(3X,'R =',E12.5,1X,'DF =',E12.5)
102  FORMAT(3X,'CA11=',E12.5,1X,'CA12=',E12.5,1X,'CA13=',E12.5,
      #1X,'CA14=',E12.5)
104  FORMAT(3X,'CA22=',E12.5,1X,'CA23=',E12.5,1X,'CA24=',E12.5)
106  FORMAT(3X,'CA33=',E12.5,1X,'CA34=',E12.5,1X,'CA44=',E12.5)
108  FORMAT(3X,'RA11=',E12.5,1X,'RA12=',E12.5,1X,'RA13=',E12.5,
      #1X,'RA14=',E12.5)
110  FORMAT(3X,'RA22=',E12.5,1X,'RA23=',E12.5,1X,'RA24=',E12.5)
112  FORMAT(3X,'RA33=',E12.5,1X,'RA34=',E12.5,1X,'RA44=',E12.5)
114  FORMAT(3X,'PO1 =',E12.5,1X,'PO2 =',E12.5,1X,'PO3 =',E12.5,
      #1X,'PO4 =',E12.5)
116  FORMAT(3X,'SI1 =',E12.5,1X,'SI2 =',E12.5,1X,'SI3 =',E12.5,
      #1X,'SI4 =',E12.5)
118  FORMAT(3X,'PR11=',E12.5,1X,'PR12=',E12.5,1X,'PR13=',E12.5,
      #1X,'PR14=',E12.5)
120  FORMAT(3X,'PR22=',E12.5,1X,'PR23=',E12.5,1X,'PR24=',E12.5)
122  FORMAT(3X,'PR33=',E12.5,1X,'PR34=',E12.5,1X,'PR44=',E12.5)
124  FORMAT(3X,'POB1=',E12.5,1X,'POB2=',E12.5,1X,'POB3=',E12.5,
      #1X,'POB4=',E12.5)
126  FORMAT(3X,'PPB1=',E12.5,1X,'PPB2=',E12.5,1X,'PPB3=',E12.5,
      #1X,'PPB4=',E12.5)
128  FORMAT(3X,'EFFECTIVE MULTIPLICATION FACTOR =',E12.5)
999  STOP
      END

```

Source Listing of ALB2

```

C      *****
C      FOUR-GROUP ALBEDOS , NONLEAKAGE PROBABILITIES AND EFFECTIVE
C      NEUTRON MULTIPLICATION FACTOR FOR A SMALL OPTICAL-PATH-
C      LENGTH CORE REACTOR (TWO-REGION GAS CORE REACTOR)
C      *****
C
C      PROGRAM ALB2
C      NTIN=01
C      NTOUT=02
C      OPEN(NTIN, FILE='ALB2.IN', STATUS='OLD')
C      OPEN(NTOUT, FILE='ALB2.OUT', STATUS='UNKNOWN')
C
C      *****
C      INPUT DATA

```

```

C      R      = RADIUS OF THE SPHERICAL CORE
C      DF      = CORE DENSITY FACTOR
C      A      = ANGULAR COEFFICIENT OF ANISOTROPIC SECONDARY
C              SOURCE ON THE REFLECTOR-CORE INTERFACE ( SEE
C              TABLE 6-11 )
C
C      CORE GROUP PARAMETERS ( SEE TABLE 6-2 )
C
C      EN1      = MACROSCOPIC (n,2n) CROSS SECTION FOR GROUP 1
C      EA1      = MACROSCOPIC ABSORPTION CROSS SECTION FOR GROUP 1
C      EA2      = MACROSCOPIC ABSORPTION CROSS SECTION FOR GROUP 2 ,
C              AND SO ON.
C      FV1      = ( AVERAGE NUMBER OF FISSION NEUTRONS RELEASED IN A
C              FISSION REACTION ) * ( MACROSCOPIC FISSION CROSS
C              SECTION ) FOR GROUP 1
C      FV2      = ( AVERAGE NUMBER OF FISSION NEUTRONS RELEASED IN A
C              FISSION REACTION ) * ( MACROSCOPIC FISSION CROSS
C              SECTION ) FOR GROUP 2 , AND SO ON.
C      CHI1     = PROBABILITY THAT A FISSION NEUTRON WILL BE BORN
C              WITH AN ENERGY IN GROUP 1
C      CHI2     = PROBABILITY THAT A FISSION NEUTRON WILL BE BORN
C              WITH AN ENERGY IN GROUP 2
C      ES11     = MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C              GROUP 1 TO GROUP 1
C      ES12     = MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C              GROUP 1 TO GROUP 2 , AND SO ON.
C
C      REFLECTOR GROUP PARAMETERS ( SEE TABLE 6-3 )
C
C      DR1      = DIFFUSION COEFFICIENT FOR GROUP 1
C      DR2      = DIFFUSION COEFFICIENT FOR GROUP 2 , AND SO ON.
C      ERN1     = MACROSCOPIC (n,2n) CROSS SECTION FOR GROUP 1
C      ERA1     = MACROSCOPIC ABSORPTION CROSS SECTION FOR GROUP 1
C      ERA2     = MACROSCOPIC ABSORPTION CROSS SECTION FOR GROUP 2 ,
C              AND SO ON.
C      ERS12    = MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C              GROUP 1 TO GROUP 2
C      ERS13    = MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C              GROUP 1 TO GROUP 3 , AND SO ON.
C      *****
C
C      READ(01,10)R,DF,A
C      READ(01,11)EN1,EA1,EA2,EA3,EA4,FV1,FV2,FV3,FV4,ES11,ES12,
C      ES13,ES14,ES22,ES23,ES24,ES33,ES34,ES44,CHI1,CHI2,DR1,DR2,
C      DR3,DR4,ERN1,ERA1,ERA2,ERA3,ERA4,ERS12,ERS13,ERS14,ERS23,
C      ERS24,ERS34
10     FORMAT(4E17.7)
11     FORMAT(4E17.7)
C      D1=D1/DF
C      D2=D2/DF
C      D3=D3/DF
C      D4=D4/DF
C      ET1=(EA1+EN1+ES11+ES12+ES13+ES14)*DF
C      ET2=(EA2+ES22+ES23+ES24)*DF
C      ET3=(EA3+ES33+ES34)*DF
C      ET4=(EA4+ES44)*DF
C      EA1=(EA1+EN1)*DF
C      EA2=EA2*DF
C      EA3=EA3*DF
C      EA4=EA4*DF
C      ES11=ES11*DF
C      ES12=ES12*DF
C      ES13=ES13*DF
C      ES14=ES14*DF
C      ES22=ES22*DF
C      ES23=ES23*DF
C      ES24=ES24*DF
C      ES33=ES33*DF
C      ES34=ES34*DF
C      ES44=ES44*DF
C      FV1=FV1*DF
C      FV2=FV2*DF
C      FV3=FV3*DF
C      FV4=FV4*DF
C      ERR1=ERA1+ERS12+ERS13+ERS14-ERN1
C      ERR2=ERA2+ERS23+ERS24
C      ERR3=ERA3+ERS34
C      ERR4=ERA4

```

```

RK1=SQRT(ERR1/DR1)
RK2=SQRT(ERR2/DR2)
RK3=SQRT(ERR3/DR3)
RK4=SQRT(ERR4/DR4)

```

```

*****
REFLECTOR ALBEDO CALCULATION
*****

```

```

Z1=EXP(-RK1*R)/R
Z2=EXP(-RK2*R)/R
Z3=EXP(-RK3*R)/R
Z4=EXP(-RK4*R)/R
G34=RK1+1./R
G35=RK2+1./R
G36=RK3+1./R
G37=RK4+1./R
RA11=(1.-2.*DR1*G34)/(1.+2.*DR1*G34)
RA22=(1.-2.*DR2*G35)/(1.+2.*DR2*G35)
RA33=(1.-2.*DR3*G36)/(1.+2.*DR3*G36)
RA44=(1.-2.*DR4*G37)/(1.+2.*DR4*G37)
G38=4./(Z1*(1.+2.*DR1*G34))
G39=G38*ERS12/(DR2*(RK2*RK2-RK1*RK1))
G40=G39*Z1*(1.+2.*DR2*G34)/(Z2*(1.+2.*DR2*G35))
G41=(G38*ERS13+G39*ERS23)/(DR3*(RK3*RK3-RK1*RK1))
G42=G40*ERS23/(DR3*(RK3*RK3-RK2*RK2))
G43=(G41*Z1*(1.+2.*DR3*G34)+G42*Z2*(1.+2.*DR3*G35))/(Z3*
  (1.+2.*DR3*G36))
G44=(G38*ERS14+G39*ERS24+G41*ERS34)/(DR4*(RK4*RK4-RK1*RK1))
G45=(G40*ERS24+G42*ERS34)/(DR4*(RK4*RK4-RK2*RK2))
G46=G43*ERS34/(DR4*(RK4*RK4-RK3*RK3))
G47=(G44*Z1*(1.+2.*DR4*G34)+G45*Z2*(1.+2.*DR4*G35)+G46*Z3*
  (1.+2.*DR4*G36))/(Z4*(1.+2.*DR4*G37))
RA12=(G40*Z2+G39*Z1)/2.
RA13=(G43*Z3+G41*Z1+G42*Z2)/2.
RA14=(G47*Z4+G44*Z1+G45*Z2+G46*Z3)/2.
G48=4./(Z2*(1.+2.*DR2*G35))
G49=G48*ERS23/(DR3*(RK3*RK3-RK2*RK2))
G50=G49*Z2*(1.+2.*DR3*G35)/(Z3*(1.+2.*DR3*G36))
G51=(G48*ERS24+G49*ERS34)/(DR4*(RK4*RK4-RK2*RK2))
G52=G50*ERS34/(DR4*(RK4*RK4-RK3*RK3))
G53=(G51*Z2*(1.+2.*DR4*G35)+G52*Z3*(1.+2.*DR4*G36))/(Z4*
  (1.+2.*DR4*G37))
RA23=(G50*Z3+G49*Z2)/2.
RA24=(G53*Z4+G51*Z2+G52*Z3)/2.
G54=4./(Z3*(1.+2.*DR3*G36))
G55=G54*ERS34/(DR4*(RK4*RK4-RK3*RK3))
G56=G55*Z3*(1.+2.*DR4*G36)/(Z4*(1.+2.*DR4*G37))
RA34=(G56*Z4+G55*Z3)/2.
RO=R/(2.*(1./3))
S1=CHI1*RA11
S2=CHI1*RA12+CHI2*RA22
S3=CHI1*RA13+CHI2*RA23
S4=CHI1*RA14+CHI2*RA24

```

```

*****
TRANSMITTANCE CALCULATION FOR THE REFLECTOR SECONDARY
SOURCES
*****

```

```

C1=-2.*ET1*R
C2=-2.*ET2*R
C3=-2.*ET3*R
C4=-2.*ET4*R
TEF1=(EXP(C1)*(1.+A-A/C1)+A/C1-1.)/(C1*(1.+A/2.))
TEF2=(EXP(C2)*(1.+A-A/C2)+A/C2-1.)/(C2*(1.+A/2.))
TEF3=(EXP(C3)*(1.+A-A/C3)+A/C3-1.)/(C3*(1.+A/2.))
TEF4=(EXP(C4)*(1.+A-A/C4)+A/C4-1.)/(C4*(1.+A/2.))
X1=(1.-TEF1)/(1.-TEF1*RA11)
X2=(1.-TEF2)/(1.-TEF2*RA22)
X3=(1.-TEF3)/(1.-TEF3*RA33)
X4=(1.-TEF4)/(1.-TEF4*RA44)
X5=TEF1/(1.-TEF1*RA11)
X6=TEF2/(1.-TEF2*RA22)
X7=TEF3/(1.-TEF3*RA33)

```

```

*****
TRANSMITTANCE CALCULATION FOR THE CORE SECONDARY
SOURCES

```

```

C *****
C
N=180
H=3.14159/FLOAT(N)
NN=N-1
XX=0.0
DO 100 K=1,NN
Y1=FLOAT(K)*H
Y2=SIN(Y1)
Y3=RO*COS(Y1)+SQRT(R*R-RO*RO*Y2*Y2)
Y4=EXP(-ET1*Y3)*Y2
XX=XX+Y4
100 CONTINUE
FFT1=H*XX/2.
YY=0.0
DO 101 L=1,NN
Y5=FLOAT(L)*H
Y6=SIN(Y5)
Y7=RO*COS(Y5)+SQRT(R*R-RO*RO*Y6*Y6)
Y8=EXP(-ET2*Y7)*Y6
YY=YY+Y8
101 CONTINUE
FFT2=H*YY/2.
ZZ=0.0
DO 102 I=1,NN
Y9=FLOAT(I)*H
Y10=SIN(Y9)
Y11=RO*COS(Y9)+SQRT(R*R-RO*RO*Y10*Y10)
Y12=EXP(-ET3*Y11)*Y10
ZZ=ZZ+Y12
102 CONTINUE
FFT3=H*ZZ/2.
WW=0.0
DO 103 J=1,NN
Y13=FLOAT(J)*H
Y14=SIN(Y13)
Y15=RO*COS(Y13)+SQRT(R*R-RO*RO*Y14*Y14)
Y16=EXP(-ET4*Y15)*Y14
WW=WW+Y16
103 CONTINUE
FFT4=H*WW/2.
C *****
C CALCULATION OF THE NEUTRON ABSORPTION BY THE CORE
C *****
C
Y17=1.-FFT4+FFT4*RA44*X4
Y18=Y17*ES44/ET4
X30=(Y17*EA4/ET4)/(1.-Y18)
X44=X4*EA4/ET4+(X4*ES44/ET4)*X30
X31=(1.-FFT3)+FFT3*RA33*X3
X32=FFT3*RA34/(1.-TEF3*RA33)
X33=X31*ES34/ET3
X34=X31*ES33/ET3
X35=(X32*X44+X33*X30)/(1.-X34)
X36=(X31*EA3/ET3)/(1.-X34)
X37=X7*RA34
X38=X3*ES34/ET3
X39=X3*ES33/ET3
X40=X37*X44+X38*X30+X39*X35
X41=X3*EA3/ET3+X39*X36
X42=FFT2*RA24/(1.-TEF2*RA22)
X43=FFT2*RA23/(1.-TEF2*RA22)
X44=1.-FFT2+FFT2*RA22*X2
X45=X44*ES24/ET2
X46=X44*ES23/ET2
X47=X44*ES22/ET2
X48=(X42*X44+X45*X30+X43*X40+X46*X35)/(1.-X47)
X49=(X43*X41+X46*X36)/(1.-X47)
X50=(X44*RA2/ET2)/(1.-X47)
X51=X6*RA24
X52=X6*RA23
X53=X2*ES24/ET2
X54=X2*ES23/ET2
X55=X2*ES22/ET2
X56=X51*X44+X53*X30+X52*X40+X54*X35+X55*X48
X57=X52*X41+X54*X36+X55*X49
X58=X2*RA2/ET2+X56*X50
X59=FFT1*RA14/(1.-TEF1*RA11)

```

```

X60=FFT1*RA13/(1.-TEF1*RA11)
X61=FFT1*RA12/(1.-TEF1*RA11)
X62=1.-FFT1+FFT1*RA11*X1
X63=X62*ES14/ET1
X64=X62*ES13/ET1
X65=X62*ES12/ET1
X66=X62*ES11/ET1
X67=(X59*XX4+X63*X30+X60*X40+X64*X35+X61*X56+X65*X48)/(1.-
#X66)
X68=(X60*X41+X64*X36+X61*X57+X65*X49)/(1.-X66)
X69=(X61*X58+X65*X50)/(1.-X66)
X70=(X62*RA1/ET1)/(1.-X66)
X71=X5*RA14
X72=X5*RA13
X73=X5*RA12
X74=X1*ES14/ET1
X75=X1*ES13/ET1
X76=X1*ES12/ET1
X77=X1*ES11/ET1
X78=X71*XX4+X74*X30+X72*X40+X75*X35+X73*X56+X76*X48+X77*X67
X79=X72*X41+X75*X36+X73*X57+X76*X49+X77*X68
X80=X73*X58+X76*X50+X77*X69
X81=X1*EA1/ET1+X77*X70
X82=S1*X81
X83=S1*X80+S2*X58
X84=S1*X79+S2*X57+S3*X41
X85=S1*X78+S2*X56+S3*X40+S4*XX4
X86=X82+X83+X84+X85
POX=0.0
XKEF=X82*(FV1+2.*EH1)/RA1+X83*FV2/EA2+X84*FV3/EA3+X85*
#FV4/EA4
X87=CHI1*X70
X88=CHI1*X69+CHI2*X50
X89=CHI1*X68+CHI2*X49
X90=CHI1*X67+CHI2*X48
X91=X87+X88+X89+X90
Q500=(1.-FFT4)*EA4/ET4
Q501=(1.-FFT4)*ES44/ET4
Q502=Q500/(1.-Q501)
Q503=(1.-FFT3)*EA3/ET3
Q504=(1.-FFT3)*ES33/ET3
Q505=(1.-FFT3)*ES34/ET3
Q506=Q503/(1.-Q504)
Q507=Q505/(1.-Q504)
Q508=Q507*Q502
Q509=(1.-FFT2)*EA2/ET2
Q510=(1.-FFT2)*ES22/ET2
Q511=(1.-FFT2)*ES23/ET2
Q512=(1.-FFT2)*ES24/ET2
Q513=Q509/(1.-Q510)
Q514=Q511/(1.-Q510)
Q515=Q512/(1.-Q510)
Q516=Q515*Q502+Q514*Q508
Q517=Q514*Q506
Q518=(1.-FFT1)*EA1/ET1
Q519=(1.-FFT1)*ES11/ET1
Q520=(1.-FFT1)*ES12/ET1
Q521=(1.-FFT1)*ES13/ET1
Q522=(1.-FFT1)*ES14/ET1
Q523=Q518/(1.-Q519)
Q524=Q520/(1.-Q519)
Q525=Q521/(1.-Q519)
Q526=Q522/(1.-Q519)
Q527=Q524*Q513
Q528=Q524*Q517+Q525*Q506
Q529=Q524*Q516+Q525*Q508+Q526*Q502
Q530=CHI1*Q523
Q531=CHI1*Q527+CHI2*Q513
Q532=CHI1*Q528+CHI2*Q517
Q533=CHI1*Q529+CHI2*Q516
Q534=Q530+Q531+Q532+Q533
Q535=X91-Q534
Q540=X87-Q530
Q541=X88-Q531
Q542=X89-Q532
Q543=X90-Q533

```

C
C
C

EFFECTIVE NEUTRON MULTIPLICATION FACTOR CALCULATION

```

C *****
C
C EKEF=X87*(FV1+2.*EN1)/EA1+X88*FV2/EA2+X89*FV3/EA3+X90*
C #FV4/EA4
C
C *****
C          OUTPUT RESULTS
C
C REFLECTOR ALBEDOS ( SEE TABLE 6-4 )
C
C RA11 = PROBABILITY THAT A NEUTRON FROM GROUP 1 BE REFLECTED
C WITH GROUP 1 ENERGY
C RA12 = PROBABILITY THAT A NEUTRON FROM GROUP 1 BE REFLECTED
C WITH GROUP 2 ENERGY , AND SO ON.
C
C CORE TRANSMITTANCES ( SEE TABLE 6-5 )
C
C TEF1 = GAMMA1
C TEF2 = GAMMA2 , AND SO ON.
C FFT1 = DELTA1
C FFT2 = DELTA2 , AND SO ON.
C
C NEUTRON ABSORPTION ( SEE TABLE 6-6 )
C
C Q530 = FRACTION OF FISSION NEUTRONS ( NEVER SPEND TIME IN
C THE REFLECTOR ) WHICH ARE ABSORBED BY THE CORE
C AS GROUP 1
C Q531 = FRACTION OF FISSION NEUTRONS ( NEVER SPEND TIME IN
C THE REFLECTOR ) WHICH ARE ABSORBED BY THE CORE
C AS GROUP 2
C Q532 = FRACTION OF FISSION NEUTRONS ( NEVER SPEND TIME IN
C THE REFLECTOR ) WHICH ARE ABSORBED BY THE CORE
C AS GROUP 3
C Q533 = FRACTION OF FISSION NEUTRONS ( NEVER SPEND TIME IN
C THE REFLECTOR ) WHICH ARE ABSORBED BY THE CORE
C AS GROUP 4
C Q540 = FRACTION OF FISSION NEUTRONS ( SPEND TIME IN THE
C REFLECTOR ) WHICH ARE ABSORBED BY THE CORE
C AS GROUP 1
C Q541 = FRACTION OF FISSION NEUTRONS ( SPEND TIME IN THE
C REFLECTOR ) WHICH ARE ABSORBED BY THE CORE
C AS GROUP 2
C Q542 = FRACTION OF FISSION NEUTRONS ( SPEND TIME IN THE
C REFLECTOR ) WHICH ARE ABSORBED BY THE CORE
C AS GROUP 3
C Q543 = FRACTION OF FISSION NEUTRONS ( SPEND TIME IN THE
C REFLECTOR ) WHICH ARE ABSORBED BY THE CORE
C AS GROUP 4
C
C EFFECTIVE NEUTRON MULTIPLICATION FACTOR ( SEE TABLE 6-7 )
C
C EKEF = EFFECTIVE NEUTRON MULTIPLICATION FACTOR
C *****
C
C WRITE(02,200)R,DF,A
C WRITE(02,202)RA11,RA12,RA13,RA14
C WRITE(02,204)RA22,RA23,RA24
C WRITE(02,206)RA33,RA34,RA44
C WRITE(02,208)TEF1,TEF2,TEF3,TEF4
C WRITE(02,210)FFT1,FFT2,FFT3,FFT4
C WRITE(02,212)Q530,Q531,Q532,Q533
C WRITE(02,214)Q540,Q541,Q542,Q543
C WRITE(02,216)EKEF
200 FORMAT(3X,'R =',E12.5,1X,'DF =',E12.5,1X,'A =',E12.5)
202 FORMAT(3X,'RA11=',E12.5,1X,'RA12=',E12.5,1X,'RA13=',E12.5,
C #1X,'RA14=',E12.5)
204 FORMAT(3X,'RA22=',E12.5,1X,'RA23=',E12.5,1X,'RA24=',E12.5)
206 FORMAT(3X,'RA33=',E12.5,1X,'RA34=',E12.5,1X,'RA44=',E12.5)
208 FORMAT(3X,'TEF1=',E12.5,1X,'TEF2=',E12.5,1X,'TEF3=',E12.5,
C #1X,'TEF4=',E12.5)
210 FORMAT(3X,'FFT1=',E12.5,1X,'FFT2=',E12.5,1X,'FFT3=',E12.5,
C #1X,'FFT4=',E12.5)
212 FORMAT(3X,'Q530=',E12.5,1X,'Q531=',E12.5,1X,'Q532=',E12.5,
C #1X,'Q533=',E12.5)
214 FORMAT(3X,'Q540=',E12.5,1X,'Q541=',E12.5,1X,'Q542=',E12.5,
C #1X,'Q543=',E12.5)
216 FORMAT(3X,'EFFECTIVE MULTIPLICATION FACTOR',9X,'=',E12.5)
C STOP
C END

```


Source Listing of ALB3

```

C *****
C TRANSPORT - THEORY REFLECTOR ALBEDOS
C ( P3 - PLANE SYMMETRY / FOUR-GROUP ENERGY )
C *****
C
C PROGRAM ALB3
C NTIN=01
C NTOUT=02
C OPEN(NTIN, FILE='ALB3.IN', STATUS='OLD')
C OPEN(NTOUT, FILE='ALB3.OUT', STATUS='UNKNOWN')
C
C *****
C INPUT DATA
C
C B = ANGULAR COEFFICIENT FOR AN ANISOTROPIC SOURCE
C ON THE REFLECTOR-CORE INTERFACE FOR REFLECTOR
C ALBEDO CALCULATION
C
C REFLECTOR GROUP PARAMETERS ( SEE TABLE 6-3 )
C
C DR1 = DIFFUSION COEFFICIENT FOR GROUP 1
C DR2 = DIFFUSION COEFFICIENT FOR GROUP 2 , AND SO ON.
C ENR1 = MACROSCOPIC (n,2n) CROSS SECTION FOR GROUP 1
C ERA1 = MACROSCOPIC ABSORPTION CROSS SECTION FOR GROUP 1
C ERT2 = MACROSCOPIC TOTAL CROSS SECTION FOR GROUP 2
C ERT3 = MACROSCOPIC TOTAL CROSS SECTION FOR GROUP 3 ,
C AND SO ON.
C ERS11= MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C GROUP 1 TO GROUP 1
C ERS13= MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C GROUP 1 TO GROUP 3 , AND SO ON.
C *****
C
C READ(01,10)B
C READ(01,11)DR1,DR2,DR3,DR4,ENR1,ERA1,ERT2,ERT3,ERT4,ERS11,
C #ERS12,ERS13,ERS14,ERS22,ERS23,ERS24,ERS33,ERS34,ERS44
10 FORMAT(E17.7)
11 FORMAT(4E17.7)
C ERT1=ERA1+ERS12+ERS13+ERS14-ENR1
C ERT2=ERT2-ERS22
C ERT3=ERT3-ERS33
C ERT4=ERT4-ERS44
C ERT1=ERA1+ENR1+ERS11+ERS12+ERS13+ERS14
C
C *****
C REFLECTOR ALBEDO CALCULATION
C *****
C
C HH=1.+B/2.
C HI=(4.+3.*B)/12.
C X30=35.*ERT1*ERT1+55.*ERT1*ERR1
C X31=SQRT(1225.*ERT1**4+3025.*ERT1*ERT1*ERR1*ERR1+70.*ERR1*
C #ERT1**3)
C X32=(X30+X31)/18.
C XN1=SQRT(X32)
C X33=(X30-X31)/18.
C XN2=SQRT(X33)
C X34=35.*ERT2*ERT2+55.*ERT2*ERR2
C X35=SQRT(1225.*ERT2**4+3025.*ERT2*ERT2*ERR2*ERR2+70.*ERR2*
C #ERT2**3)
C X36=(X34+X35)/18.
C XN1=SQRT(X36)
C X37=(X34-X35)/18.
C XN2=SQRT(X37)
C X38=35.*ERT3*ERT3+55.*ERT3*ERR3
C X39=SQRT(1225.*ERT3**4+3025.*ERT3*ERT3*ERR3*ERR3+70.*ERR3*
C #ERT3**3)
C X40=(X38+X39)/18.
C XN1=SQRT(X40)
C X41=(X38-X39)/18.
C XN2=SQRT(X41)
C X42=35.*ERT4*ERT4+55.*ERT4*ERR4
C X43=SQRT(1225.*ERT4**4+3025.*ERT4*ERT4*ERR4*ERR4+70.*ERR4*
C #ERT4**3)

```

```

X44=(X42+X43)/18.
XQ1=SQRT(X44)
X45=(X42-X43)/18.
XQ2=SQRT(X45)
X50=ERR1/XN1
X51=ERR1/XN2
X52=(3.*ERT1*X50-XN1)/(2.*XN1)
X53=(3.*ERT1*X51-XN2)/(2.*XN2)
X54=(3.*XN1*X52)/(7.*ERT1)
X55=(3.*XN2*X53)/(7.*ERT1)
X56=10.+24.*X50+25.*X52+16.*X54
X57=10.+24.*X51+25.*X53+16.*X55
X58=8.*X50+5.*X52+4.
X59=8.*X51+5.*X53+4.
X60=(80.*HI*X59-16.*HH*X57)/(X56*X59-X57*X58)
X61=(16.*HH*X56-80.*HI*X58)/(X56*X59-X57*X58)
X62=((0.5-X50+5.*X52/8.)*X60+(0.5-X51+5.*X53/8.)*X61)/2.
RA11=X62/HH
SQ1=35.*ERT2*ERT2+55.*ERT2*ERR2
SQ2=105.*ERR2*ERT2**3
SQ3=35.*ERT3*ERT3+55.*ERT3*ERR3
SQ4=105.*ERR3*ERT3**3
SQ5=35.*ERT4*ERT4+55.*ERT4*ERR4
SQ6=105.*ERR4*ERT4**3
X100=9.*XN1**4-SQ1*XN1**2+SQ2
X101=9.*XN2**4-SQ1*XN2**2+SQ2
X102=9.*XN1**4-SQ3*XN1**2+SQ4
X103=9.*XN2**4-SQ3*XN2**2+SQ4
X104=9.*XN1**4-SQ3*XN1**2+SQ4
X105=9.*XN2**4-SQ3*XN2**2+SQ4
X106=9.*XN1**4-SQ5*XN1**2+SQ6
X107=9.*XN2**4-SQ5*XN2**2+SQ6
X108=9.*XN1**4-SQ5*XN1**2+SQ6
X109=9.*XN2**4-SQ5*XN2**2+SQ6
X110=9.*XP1**4-SQ5*XP1**2+SQ6
X111=9.*XP2**4-SQ5*XP2**2+SQ6
X63=(35.*XN1*ERT2**2-9.*XN1**3)*ERS12*X60
X64=X63/X100
X65=(35.*XN2*ERT2**2-9.*XN2**3)*ERS12*X61
X66=X65/X101
X67=14.*ERT2*ERS12*X60*XN1*XN1/X100
X68=14.*ERT2*ERS12*X61*XN2*XN2/X101
X69=(ERS12*X60+XN1*X64)/ERR2
X70=(ERS12*X61+XN2*X66)/ERR2
X71=3.*XN1*X67/(7.*ERT2)
X72=3.*XN2*X68/(7.*ERT2)
X73=(3.*ERT2-XN1*XN1/ERR2)/(2.*XN1)
X74=(3.*ERT2-XN2*XN2/ERR2)/(2.*XN2)
X75=XN1/ERR2
X76=XN2/ERR2
X77=3.*XN1*X73/(7.*ERT2)
X78=3.*XN2*X74/(7.*ERT2)
X79=4.*X75+8.+5.*X73
X80=4.*X76+8.+5.*X74
X81=-8.*(X64+X66)-5.*(X67+X68)-4.*(X69+X70)
X82=10.*X75+24.+25.*X73+16.*X77
X83=10.*X76+24.+25.*X74+16.*X78
X84=-24.*(X64+X66)-25.*(X67+X68)-10.*(X69+X70)-16.*(X71+
#X72)
X85=X79*X83-X80*X82
X86=(X81*X83-X80*X84)/X85
X87=(X79*X84-X81*X82)/X85
X88=((X75*X86+X76*X87+X69+X70)/2.-(X86+X87+X64+X66)+5.*(
#X73*X86+X74*X87+X67+X68)/8.)/2.
RA12=X88/HH
X89=XN1*X86/ERR2
X90=XN2*X87/ERR2
X91=(-9.*XN1**3+35.*XN1*ERT3**2)*(ERS13*X60+ERS23*X69)
X92=X91/X102
X93=(-9.*XN2**3+35.*XN2*ERT3**2)*(ERS13*X61+ERS23*X70)
X94=X93/X103
X95=(-9.*XN1**3+35.*XN1*ERT3**2)*ERS23*X89/X104
X96=(-9.*XN2**3+35.*XN2*ERT3**2)*ERS23*X90/X105
Y10=14.*ERT3*(ERS13*X60+ERS23*X69)*XN1**2/X102
Y11=14.*ERT3*(ERS13*X61+ERS23*X70)*XN2**2/X103
Y12=14.*ERT3*ERS23*X89*XN1**2/X104
Y13=14.*ERT3*ERS23*X90*XN2**2/X105
Y14=(ERS13*X60+ERS23*X69+XN1*X92)/ERR3
Y15=(ERS13*X61+ERS23*X70+XN2*X94)/ERR3

```

```

Y16=(ERS23*X89+X1*X95)/ERR3
Y17=(ERS23*X90+X2*X96)/ERR3
Y18=3.*Y10*X1/(7.*ERT3)
Y19=3.*Y11*X2/(7.*ERT3)
Y20=3.*Y12*X1/(7.*ERT3)
Y21=3.*Y13*X2/(7.*ERT3)
Y22=XP1/ERR3
Y23=XP2/ERR3
Y24=(3.*ERT3-XP1*XP1/ERR3)/(2.*XP1)
Y25=(3.*ERT3-XP2*XP2/ERR3)/(2.*XP2)
Y26=3.*XP1*Y24/(7.*ERT3)
Y27=3.*XP2*Y25/(7.*ERT3)
Y28=4.*Y22+8.*Y5
Y29=4.*Y23+8.*Y25
Y30=-4.*(Y14+Y15+Y16+Y17)-8.*(X92+X94+X95+X96)-5.*(Y10+
#Y11+Y12+Y13)
Y31=10.*Y22+24.*Y24+16.*Y26
Y32=10.*Y23+24.*Y25+16.*Y27
Y33=-10.*(Y14+Y15+Y16+Y17)-24.*(X92+X94+X95+X96)-25.*(Y10+
#Y11+Y12+Y13)-16.*(Y18+Y19+Y20+Y21)
Y34=Y28*Y32-Y29*Y31
Y35=(Y30*Y32-Y29*Y33)/Y34
Y36=(Y28*Y33-Y30*Y31)/Y34
Y37=(Y22*Y35+Y23*Y36+Y14+Y15+Y16+Y17)/2.-(Y35+Y36+X92+X94+
#X95+X96)+5.*(Y24*Y35+Y25*Y36+Y10+Y11+Y12+Y13)/8./2.
RA13=Y37/HH
Y38=XP1*Y35/ERR3
Y39=XP2*Y36/ERR3
Y40=(-9.*X1**3+35.*XN1*ERT4**2)*(ERS34*Y14+ERS24*X69+
#ERS14*X60)
Y41=Y40/X106
Y42=(-9.*XN2**3+35.*XN2*ERT4**2)*(ERS34*Y15+ERS24*X70+
#ERS14*X61)
Y43=Y42/X107
Y44=(-9.*X1**3+35.*XN1*ERT4**2)*(ERS34*Y16+ERS24*X89)
Y45=Y44/X108
Y46=(-9.*XN2**3+35.*XN2*ERT4**2)*(ERS34*Y17+ERS24*X90)
Y47=Y46/X109
Y48=(-9.*XP1**3+35.*XP1*ERT4**2)*ERS34*Y38/X110
Y49=(-9.*XP2**3+35.*XP2*ERT4**2)*ERS34*Y39/X111
Y50=14.*ERT4*(ERS34*Y14+ERS24*X69+ERS14*X60)*XN1**2
Y51=Y50/X106
Y52=14.*ERT4*(ERS34*Y15+ERS24*X70+ERS14*X61)*XN2**2
Y53=Y52/X107
Y54=14.*ERT4*(ERS34*Y16+ERS24*X89)*XN1**2
Y55=Y54/X108
Y56=14.*ERT4*(ERS34*Y17+ERS24*X90)*XN2**2
Y57=Y56/X109
Y58=14.*ERT4*ERS34*Y38*XP1**2/X110
Y59=14.*ERT4*ERS34*Y39*XP2**2/X111
Y60=(ERS34*Y14+ERS24*X69+ERS14*X60+Y41*XN1)/ERR4
Y61=(ERS34*Y15+ERS24*X70+ERS14*X61+Y43*XN2)/ERR4
Y62=(ERS34*Y16+ERS24*X89+Y45*XN1)/ERR4
Y63=(ERS34*Y17+ERS24*X90+Y47*XN2)/ERR4
Y64=(ERS34*Y38+Y48*XP1)/ERR4
Y65=(ERS34*Y39+Y49*XP2)/ERR4
Y65=3.*Y51*XN1/(7.*ERT4)
Y66=3.*Y53*XN2/(7.*ERT4)
Y67=3.*Y55*XN1/(7.*ERT4)
Y68=3.*Y57*XN2/(7.*ERT4)
Y69=3.*Y58*XP1/(7.*ERT4)
Y70=3.*Y59*XP2/(7.*ERT4)
Y71=XQ1/ERR4
Y72=XQ2/ERR4
Y73=(3.*ERT4-XQ1*XQ1/ERR4)/(2.*XQ1)
Y74=(3.*ERT4-XQ2*XQ2/ERR4)/(2.*XQ2)
Y75=3.*XQ1*Y73/(7.*ERT4)
Y76=3.*XQ2*Y74/(7.*ERT4)
Y78=4.*Y71+8.*Y73
Y79=4.*Y72+8.*Y74
Y80=-8.*(Y41+Y43+Y45+Y47+Y48+Y49)-5.*(Y51+Y53+Y55+Y57+
#Y58+Y59)-4.*(Y60+Y61+Y62+Y63+Y64+Y65)
Y81=10.*Y71+24.*Y73+16.*Y75
Y82=10.*Y72+24.*Y74+16.*Y76
Y83=-24.*(Y41+Y43+Y45+Y47+Y48+Y49)-25.*(Y51+Y53+Y55+Y57+
#Y58+Y59)-10.*(Y60+Y61+Y62+Y63+Y64+Y65)-16.*(Y75+Y76+Y77+
#Y68+Y69+Y70)
Y84=Y78*Y82-Y79*Y81
Y85=(Y80*Y82-Y79*Y83)/Y84

```

```

Y86=(Y78*Y83-Y80*Y81)/Y84
Y87=((Y71*Y85+Y72*Y86+Y60+Y61+Y62+Y63+Y64+Y65)/2.-(Y85+
#Y86+Y41+Y43+Y45+Y47+Y48+Y49)+5.*(Y73*Y85+Y74*Y86+Y51+Y53+
#Y55+Y57+Y58+Y59)/8.)/2.
RA14=Y87/HH
W10=ERR2/XM1
W11=ERR2/XM2
W12=(3.*ERT2*W10-XM1)/(2.*XM1)
W13=(3.*ERT2*W11-XM2)/(2.*XM2)
W14=3.*XM1*W12/(7.*ERT2)
W15=3.*XM2*W13/(7.*ERT2)
W16=10.+24.*W10+25.*W12+16.*W14
W17=10.+24.*W11+25.*W13+16.*W15
W18=8.*W10+5.*W12+4.
W19=8.*W11+5.*W13+4.
W20=W16*W19-W17*W18
W21=16.*(5.*W19*HI-W17*HH)/W20
W22=16.*(W16*HH-5.*W18*HI)/W20
W23=((0.5-W10+5.*W12/8.)*W21+(0.5-W11+5.*W13/8.)*W22)/2.
RA22=W23/HH
W24=(35.*XM1*ERT3**2-9.*XM1**3)*ERS23*W21/X104
W25=(35.*XM2*ERT3**2-9.*XM2**3)*ERS23*W22/X105
W26=14.*ERT3*ERS23*W21*XM1**2/X104
W27=14.*ERT3*ERS23*W22*XM2**2/X105
W28=(ERS23*W21+XM1*W24)/ERR3
W29=(ERS23*W22+XM2*W25)/ERR3
W30=3.*XM1*W26/(7.*ERT3)
W31=3.*XM2*W27/(7.*ERT3)
W32=(3.*ERT3*XP1*XP1/ERR3)/(2.*XP1)
W33=(3.*ERT3*XP2*XP2/ERR3)/(2.*XP2)
W34=XP1/ERR3
W35=XP2/ERR3
W36=3.*XP1*W32/(7.*ERT3)
W37=3.*XP2*W33/(7.*ERT3)
W38=4.*W34+8.+5.*W32
W39=4.*W35+8.+5.*W33
W40=-8.*(W24+W25)-5.*(W26+W27)-4.*(W28+W29)
W41=10.*W34+24.+25.*W32+16.*W36
W42=10.*W35+24.+25.*W33+16.*W37
W43=-24.*(W24+W25)-25.*(W26+W27)-10.*(W28+W29)-16.*(
#W30+W31)
W44=W38*W42-W39*W41
W45=(W40*W42-W39*W43)/W44
W46=(W38*W43-W40*W41)/W44
W47=((W34*W45+W35*W46+W28+W29)/2.-(W45+W46+W24+W25)+5.*(
#W32*W45+W33*W46+W26+W27)/8.)/2.
RA23=W47/HH
W48=W34*W45
W49=W35*W46
W50=(-9.*XM1**3+35.*XM1*ERT4**2)*(ERS24*W21+ERS34*W28)/X108
W51=(-9.*XM2**3+35.*XM2*ERT4**2)*(ERS24*W22+ERS34*W29)/X109
W52=(-9.*XP1**3+35.*XP1*ERT4**2)*ERS34*W48/X110
W53=(-9.*XP2**3+35.*XP2*ERT4**2)*ERS34*W49/X111
W54=14.*ERT4*XM1**2*(ERS24*W21+ERS34*W28)/X108
W55=14.*ERT4*XM2**2*(ERS24*W22+ERS34*W29)/X109
W56=14.*ERT4*ERS34*W48*XP1**2/X110
W57=14.*ERT4*ERS34*W49*XP2**2/X111
W58=(ERS24*W21+ERS34*W28+XM1*W50)/ERR4
W59=(ERS24*W22+ERS34*W29+XM2*W51)/ERR4
W60=(ERS34*W48+XP1*W52)/ERR4
W61=(ERS34*W49+XP2*W53)/ERR4
W62=3.*XM1*W54/(7.*ERT4)
W63=3.*XM2*W55/(7.*ERT4)
W64=3.*XP1*W56/(7.*ERT4)
W65=3.*XP2*W57/(7.*ERT4)
W66=(3.*ERT4*XQ1*XQ1/ERR4)/(2.*XQ1)
W67=(3.*ERT4*XQ2*XQ2/ERR4)/(2.*XQ2)
W68=XQ1/ERR4
W69=XQ2/ERR4
W70=3.*XQ1*W66/(7.*ERT4)
W71=3.*XQ2*W67/(7.*ERT4)
W72=4.*W68+8.+5.*W66
W73=4.*W69+8.+5.*W67
W74=-8.*(W50+W51+W52+W53)-5.*(W54+W55+W56+W57)-4.*(
#W58+W59+W60+W61)
W75=10.*W68+24.+25.*W66+16.*W70
W76=10.*W69+24.+25.*W67+16.*W71
W77=-24.*(W50+W51+W52+W53)-25.*(W54+W55+W56+W57)-10.*(
#W58+W59+W60+W61)-16.*(W62+W63+W64+W65)

```

```

W78=W72*W76-W73*W75
W79=(W74*W76-W73*W77)/W78
W80=(W72*W77-W74*W75)/W78
W81=((W68*W79+W69*W80+W58+W59+W60+W61)/2.-(W79+W80+W50+
#W51+W52+W53)+5.*(W66*W79+W67*W80+W54+W55+W56+W57)/8.)/2.
RA24=W81/HH
Z10=ERR3/XP1
Z11=ERR3/XP2
Z12=(3.*ERT3*Z10-XP1)/(2.*XP1)
Z13=(3.*ERT3*Z11-XP2)/(2.*XP2)
Z14=3.*XP1*Z12/(7.*ERT3)
Z15=3.*XP2*Z13/(7.*ERT3)
Z16=10.+24.*Z10+25.*Z12+16.*Z14
Z17=10.+24.*Z11+25.*Z13+16.*Z15
Z18=8.*Z10+5.*Z12+4.
Z19=8.*Z11+5.*Z13+4.
Z20=Z16*Z19-Z17*Z18
Z21=16.*(5.*HI*Z19-HH*Z17)/Z20
Z22=16.*(HH*Z16-5.*HI*Z18)/Z20
Z23=((0.5-Z10+5.*Z12/8.)*Z21+(0.5-Z11+5.*Z13/8.)*Z22)/2.
RA33=Z23/HH
Z24=(35.*XP1*ERT4**2-9.*XP1**3)*ERS34*Z21/X110
Z25=(35.*XP2*ERT4**2-9.*XP2**3)*ERS34*Z22/X111
Z26=14.*ERT4*ERS34*Z21*XP1**2/X110
Z27=14.*ERT4*ERS34*Z22*XP2**2/X111
Z28=(ERS34*Z21+XP1*Z24)/ERR4
Z29=(ERS34*Z22+XP2*Z25)/ERR4
Z30=3.*XP1*Z26/(7.*ERT4)
Z31=3.*XP2*Z27/(7.*ERT4)
Z32=XQ1/ERR4
Z33=XQ2/ERR4
Z34=(3.*ERT4-XQ1*XQ1/ERR4)/(2.*XQ1)
Z35=(3.*ERT4-XQ2*XQ2/ERR4)/(2.*XQ2)
Z36=3.*XQ1*Z34/(7.*ERT4)
Z37=3.*XQ2*Z35/(7.*ERT4)
Z38=4.*Z32+8.+5.*Z34
Z39=4.*Z33+8.+5.*Z35
Z40=8.*(Z24+Z25)-5.*(Z26+Z27)-4.*(Z28+Z29)
Z41=10.*(Z32+24.+25.*Z34+16.*Z36
Z42=10.*(Z33+24.+25.*Z35+16.*Z37
Z43=-24.*(Z24+Z25)-25.*(Z26+Z27)-10.*(Z28+Z29)-16.*(
#Z30+Z31)
Z44=Z38*Z42-Z39*Z41
Z45=(Z40*Z42-Z39*Z43)/Z44
Z46=(Z38*Z43-Z40*Z41)/Z44
Z47=((Z32*Z45+Z33*Z46+Z28+Z29)/2.-(Z45+Z46+Z24+Z25)+5.*(
#Z34*Z45+Z35*Z46+Z26+Z27)/8.)/2.
RA34=Z47/HH
Z48=ERR4/XQ1
Z49=ERR4/XQ2
Z50=(3.*ERT4*Z48-XQ1)/(2.*XQ1)
Z51=(3.*ERT4*Z49-XQ2)/(2.*XQ2)
Z52=3.*XQ1*Z50/(7.*ERT4)
Z53=3.*XQ2*Z51/(7.*ERT4)
Z54=10.+24.*Z48+25.*Z50+16.*Z52
Z55=10.+24.*Z49+25.*Z51+16.*Z53
Z56=8.*Z48+5.*Z50+4.
Z57=8.*Z49+5.*Z51+4.
Z58=Z54*Z57-Z55*Z56
Z59=16.*(5.*HI*Z57-HH*Z55)/Z58
Z60=16.*(HH*Z54-5.*HI*Z56)/Z58
Z61=((0.5-Z48+5.*Z50/8.)*Z59+(0.5-Z49+5.*Z51/8.)*Z60)/2.
RA44=Z61/HH

```

```

*****
OUTPUT RESULTS

```

```

REFLECTOR ALBEDOS (SEE TABLE )

```

```

RA11 = PROBABILITY THAT A NEUTRON FROM GROUP 1 BE
      REFLECTED WITH GROUP 1 ENERGY
RA12 = PROBABILITY THAT A NEUTRON FROM GROUP 1 BE
      REFLECTED WITH GROUP 2 ENERGY , AND SO ON.
*****

```

```

WRITE (02,100)B
WRITE (02,102)RA11,RA12,RA13,RA14
WRITE (02,104)RA22,RA23,RA24
WRITE (02,106)RA33,RA34,RA44

```

```

100 FORMAT(3X,'B  =',E12.5)
102 FORMAT(3X,'RA11=',E12.5,1X,'RA12=',E12.5,1X,'RA13=',E12.5,
1X,'RA14=',E12.5)
104 FORMAT(3X,'RA22=',E12.5,1X,'RA23=',E12.5,1X,'RA24=',E12.5)
106 FORMAT(3X,'RA33=',E12.5,1X,'RA34=',E12.5,1X,'RA44=',E12.5)
STOP
END

```

Source Listing of ALB4

```

C
C *****
C FOUR-GROUP ALBEDOS, NONLEAKAGE PROBABILITIES AND EFFECTIVE
C NEUTRON MULTIPLICATION FACTOR FOR A SMALL OPTICAL-PATH-
C LENGTH CORE REACTOR ( THREE-REGION GAS CORE REACTOR )
C *****
C
C PROGRAM ALB4
C NTIN=01
C NTOUT=02
C OPEN(NTIN, FILE='ALB4.IN', STATUS='OLD')
C OPEN(NTOUT, FILE='ALB4.OUT', STATUS='UNKNOWN')
C
C *****
C INPUT DATA
C
C RP = OUTER HYDROGEN RADIUS
C RC = FUEL RADIUS
C PDF = HYDROGEN DENSITY FACTOR
C
C OSCILLATIONS OF THE FUEL-HYDROGEN BOUNDARY
C RDEL = DIFFERENCE BETWEEN THE WAVE MAXIMUM AMPLITUDE
C AND THE UNPERTURBED FUEL RADIUS ( RC )
C
C NONUNIFORM ATOMIC DENSITY DISTRIBUTION IN THE FUEL
C REGION ( SEE TABLE 7-9 )
C AF = ANGULAR COEFFICIENT OF A LINEAR ATOMIC DENSITY
C DISTRIBUTION
C
C FUEL GROUP PARAMETERS ( SEE TABLE 6-2 )
C
C EN1 = MACROSCOPIC (n,2n) CROSS SECTION FOR GROUP 1
C EA1 = MACROSCOPIC ABSORPTION CROSS SECTION FOR GROUP 1
C EA2 = MACROSCOPIC ABSORPTION CROSS SECTION FOR GROUP 2 ,
C AND SO ON.
C FV1 = ( AVERAGE NUMBER OF FISSION NEUTRONS RELEASED IN A
C FISSION REACTION ) * ( MACROSCOPIC FISSION CROSS
C SECTION ) FOR GROUP 1
C FV2 = ( AVERAGE NUMBER OF FISSION NEUTRONS RELEASED IN A
C FISSION REACTION ) * ( MACROSCOPIC FISSION CROSS
C SECTION ) FOR GROUP 2 , AND SO ON.
C ET1 = MACROSCOPIC TOTAL CROSS SECTION FOR GROUP 1
C ET2 = MACROSCOPIC TOTAL CROSS SECTION FOR GROUP 2 ,
C AND SO ON.
C ES11 = MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C GROUP 1 TO GROUP 1
C ES12 = MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C GROUP 1 TO GROUP 2 , AND SO ON.
C CHI1 = PROBABILITY THAT A FISSION NEUTRON WILL BE BORN
C WITH AN ENERGY IN GROUP 1
C CHI2 = PROBABILITY THAT A FISSION NEUTRON WILL BE BORN
C WITH AN ENERGY IN GROUP 2
C
C HYDROGEN GROUP PARAMETERS ( SEE TABLE 7-2 )
C
C EPA1 = MACROSCOPIC ABSORPTION CROSS SECTION FOR GROUP 1
C EPA2 = MACROSCOPIC ABSORPTION CROSS SECTION FOR GROUP 2 ,
C AND SO ON.
C EPS11 = MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C GROUP 1 TO GROUP 1
C EPS12 = MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C GROUP 1 TO GROUP 2 , AND SO ON.
C
C REFLECTOR GROUP PARAMETERS ( SEE TABLE 6-3 )
C

```

```

C      DR1 = DIFFUSION COEFFICIENT FOR GROUP 1
C      DR2 = DIFFUSION COEFFICIENT FOR GROUP 2 , AND SO ON.
C      ERN1 = MACROSCOPIC (n,2n) CROSS SECTION FOR GROUP 1
C      ERA1 = MACROSCOPIC ABSORPTION CROSS SECTION FOR GROUP 1
C      ERA2 = MACROSCOPIC ABSORPTION CROSS SECTION FOR GROUP 2 ,
C            AND SO ON.
C      ERS12= MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C            GROUP 1 TO GROUP 2
C      ERS13= MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C            GROUP 1 TO GROUP 3 , AND SO ON.
C      *****
C
10     READ(01,10)RP,RC
11     READ(01,11)PDF,RDEL,AF
12     READ(01,12)ERN1,EA1,EA2,EA3,EA4,FV1,FV2,FV3,FV4,ET1,ET2,ET3,
      #ET4,ES11,ES12,ES13,ES14,ES22,ES23,ES24,ES33,ES34,ES44,CHI1,
      #CHI2,EPA1,EPA2,EPA3,EPA4,EPS11,EPS12,EPS13,EPS14,EPS22,
      #EPS23,EPS24,EPS33,EPS34,EPS44,DR1,DR2,DR3,DR4,ERN1,ERA1,
      #ERA2,ERA3,ERA4,ERS12,ERS13,ERS14,ERS23,ERS24,ERS34
10     FORMAT(2E17.7)
11     FORMAT(3E17.7)
12     FORMAT(4E17.7)
      ERR1=ERA1+ERS12+ERS13+ERS14-ERN1
      ERR2=ERA2+ERS23+ERS24
      ERR3=ERA3+ERS34
      ERR4=ERA4
C
C      *****
C      REFLECTOR ALBEDO CALCULATION
C      *****
C
      RK1=SQRT(ERR1/DR1)
      RK2=SQRT(ERR2/DR2)
      RK3=SQRT(ERR3/DR3)
      RK4=SQRT(ERR4/DR4)
      Z1=EXP(-RK1*RP)/RP
      Z2=EXP(-RK2*RP)/RP
      Z3=EXP(-RK3*RP)/RP
      Z4=EXP(-RK4*RP)/RP
      G34=RK1+1./RP
      G35=RK2+1./RP
      G36=RK3+1./RP
      G37=RK4+1./RP
      RA11=(1.-2.*DR1*G34)/(1.+2.*DR1*G34)
      RA22=(1.-2.*DR2*G35)/(1.+2.*DR2*G35)
      RA33=(1.-2.*DR3*G36)/(1.+2.*DR3*G36)
      RA44=(1.-2.*DR4*G37)/(1.+2.*DR4*G37)
      G38=4./(Z1*(1.+2.*DR1*G34))
      G39=G38*ERS12/(DR2*(RK2*RK2-RK1*RK1))
      G40=G39*Z1*(1.+2.*DR2*G34)/(Z2*(1.+2.*DR2*G35))
      G41=(G38*ERS13+G39*ERS23)/(DR3*(RK3*RK3-RK1*RK1))
      G42=G40*ERS23/(DR3*(RK3*RK3-RK2*RK2))
      G43=(G41*Z1*(1.+2.*DR3*G34)+G42*Z2*(1.+2.*DR3*G35))/(Z3*(
      #1.+2.*DR3*G36))
      G44=(G38*ERS14+G39*ERS24+G41*ERS34)/(DR4*(RK4*RK4-RK1*RK1))
      G45=(G40*ERS24+G42*ERS34)/(DR4*(RK4*RK4-RK2*RK2))
      G46=G43*ERS34/(DR4*(RK4*RK4-RK3*RK3))
      G47=(G44*Z1*(1.+2.*DR4*G34)+G45*Z2*(1.+2.*DR4*G35)+G46*Z3*
      #1.(1.+2.*DR4*G36))/(Z4*(1.+2.*DR4*G37))
      RA12=(G40*Z2+G39*Z1)/2.
      RA13=(G43*Z3+G41*Z1+G42*Z2)/2.
      RA14=(G47*Z4+G44*Z1+G45*Z2+G46*Z3)/2.
      G48=4./(Z2*(1.+2.*DR2*G35))
      G49=G48*ERS23/(DR3*(RK3*RK3-RK2*RK2))
      G50=G49*Z2*(1.+2.*DR3*G35)/(Z3*(1.+2.*DR3*G36))
      G51=(G48*ERS24+G49*ERS34)/(DR4*(RK4*RK4-RK2*RK2))
      G52=G50*ERS34/(DR4*(RK4*RK4-RK3*RK3))
      G53=(G51*Z2*(1.+2.*DR4*G35)+G52*Z3*(1.+2.*DR4*G36))/(Z4*(
      #1.+2.*DR4*G37))
      RA23=(G50*Z3+G49*Z2)/2.
      RA24=(G53*Z4+G51*Z2+G52*Z3)/2.
      G54=4./(Z3*(1.+2.*DR3*G36))
      G55=G54*ERS34/(DR4*(RK4*RK4-RK3*RK3))
      G56=G55*Z3*(1.+2.*DR4*G36)/(Z4*(1.+2.*DR4*G37))
      RA34=(G56*Z4+G55*Z3)/2.
C
C      *****
C      CALCULATION OF FUEL AND HYDROGEN TRANSMITTANCES
C      *****

```

C

```

EPT1=EPA1+EPS11+EPS12+EPS13+EPS14
EPT2=EPA2+EPS22+EPS23+EPS24
EPT3=EPA3+EPS33+EPS34
EPT4=EPA4+EPS44
EPA1=PDF*EPA1
EPA2=PDF*EPA2
EPA3=PDF*EPA3
EPA4=PDF*EPA4
EPS11=PDF*EPS11
EPS12=PDF*EPS12
EPS13=PDF*EPS13
EPS14=PDF*EPS14
EPS22=PDF*EPS22
EPS23=PDF*EPS23
EPS24=PDF*EPS24
EPS33=PDF*EPS33
EPS34=PDF*EPS34
EPS44=PDF*EPS44
EPT1=PDF*EPT1
EPT2=PDF*EPT2
EPT3=PDF*EPT3
EPT4=PDF*EPT4
EA1=EA1+EN1
FR=1.0
WO=ASIN(FR*RC/RP)
A=0.0
N=90
H=(WO-A)/FLOAT(N)
NN=N-1
AFF=(1.+AF*RC/2.)/(1.+0.75*AF*RC)
FF=AFF/(FR*FR*FR)
FP=(RP**3-RC**3)/(RP**3-(FR*RC)**3)
SAP=SQRT(RP*RP-FR*FR*RC*RC)/RP
SAC=1.-SAP
SX1=FR*RC/RP
XX=0.0
DO 40 K=1,NN
  X1=A+FLOAT(K)*H
  Q1=2.*FF*ET1*SQRT(FR*FR*RC*RC-RP*RP*SIN(X1)*SIN(X1))
  X=EXP(-Q1)*SIN(X1)
  XX=XX+X
40 CONTINUE
EF1=H*(SX1/2.+XX)/SAC
SY1=FR*RC/RP
YY=0.0
DO 42 J=1,NN
  Y1=A+FLOAT(J)*H
  Q2=2.*FF*ET2*SQRT(FR*FR*RC*RC-RP*RP*SIN(Y1)*SIN(Y1))
  Y=EXP(-Q2)*SIN(Y1)
  YY=YY+Y
42 CONTINUE
EF2=H*(SY1/2.+YY)/SAC
SW1=FR*RC/RP
WW=0.0
DO 44 I=1,NN
  W1=A+FLOAT(I)*H
  Q3=2.*FF*ET3*SQRT(FR*FR*RC*RC-RP*RP*SIN(W1)*SIN(W1))
  W=EXP(-Q3)*SIN(W1)
  WW=WW+W
44 CONTINUE
EF3=H*(SW1/2.+WW)/SAC
SZ1=FR*RC/RP
ZZ=0.0
DO 46 L=1,NN
  Z5=A+FLOAT(L)*H
  Q4=2.*FF*ET4*SQRT(FR*FR*RC*RC-RP*RP*SIN(Z5)*SIN(Z5))
  Z=EXP(-Q4)*SIN(Z5)
  ZZ=ZZ+Z
46 CONTINUE
EF4=H*(SZ1/2.+ZZ)/SAC
C4=RP*SAP
TX1=EXP(-EPT1*FP*C4)*FR*RC/RP
XX1=0.0
DO 48 KK=1,NN
  X11=A+FLOAT(KK)*H
  H2=FP*EPT1*(RP*COS(X11)-SQRT(FR*FR*RC*RC-RP*RP*SIN(X11)
    *2))
  X2=EXP(-H2)*SIN(X11)

```



```

XX1=XX1+X2
48  CONTINUE
   EPP1=H*(TX1/2.+XX1)/SAC
   TK2=EXP(-EPT2*FP*C4)*FR*RC/RP
   XX2=0.0
   DO 50 JJ=1,NN
   X12=A+FLOAT(JJ)*H
   H3=FP*EPT2*(RP*COS(X12)-SQRT(FR*FR*RC*RC-RP*RP)*SIN(X12)
   & **2)
   X3=EXP(-H3)*SIN(X12)
   XX2=XX2+X3
50  CONTINUE
   EPP2=H*(TX2/2.+XX2)/SAC
   TK3=EXP(-EPT3*FP*C4)*FR*RC/RP
   XX3=0.0
   DO 52 LL=1,NN
   X13=A+FLOAT(LL)*H
   H4=FP*EPT3*(RP*COS(X13)-SQRT(FR*FR*RC*RC-RP*RP)*SIN(X13)
   & **2)
   X4=EXP(-H4)*SIN(X13)
   XX3=XX3+X4
52  CONTINUE
   EPP3=H*(TX3/2.+XX3)/SAC
   TK4=EXP(-EPT4*FP*C4)*FR*RC/RP
   XX4=0.0
   DO 54 II=1,NN
   X14=A+FLOAT(II)*H
   H5=FP*EPT4*(RP*COS(X14)-SQRT(FR*FR*RC*RC-RP*RP)*SIN(X14)
   & **2)
   X5=EXP(-H5)*SIN(X14)
   XX4=XX4+X5
54  CONTINUE
   EPP4=H*(TX4/2.+XX4)/SAC
   W11=FP*EPT1*2.*RP*SAP
   OPP1=(1.-EXP(-W11))/W11
   W12=FP*EPT2*2.*RP*SAP
   OPP2=(1.-EXP(-W12))/W12
   W13=FP*EPT3*2.*RP*SAP
   OPP3=(1.-EXP(-W13))/W13
   W14=FP*EPT4*2.*RP*SAP
   OPP4=(1.-EXP(-W14))/W14
   H11=1.-(SAC*EPP4*EPP4*EF4+SAP*OPP4)*RA44
   H12=1.-(SAC*EPP3*EPP3*EF3+SAP*OPP3)*RA33
   H13=1.-(SAC*EPP2*EPP2*EF2+SAP*OPP2)*RA22
   H14=1.-(SAC*EPP1*EPP1*EF1+SAP*OPP1)*RA11
   RCC=FR*RC
   RO=( (RP**3+RCC**3)/2. )**(1./3.)
   IF(AF)601,600,601
600  RCC=RCC/(2.** (1./3))
   GO TO 602
601  PT=3.14159
   A1=4./(3.*AF)
   A4=-2.*RCC*RCC*RCC*(1./3.+AF*RCC/4.)/AF
   CA2=-4.*A4
   CA3=-A1*A1*A4
   P1=CA2/3.
   P2=-CA3/2.
   P3=P1**3+P2*P2
   P4=P2*(ABS(P3))**0.5
   P5=P2-(ABS(P3))**0.5
   P6=(ABS(P4))** (1./3.)
   P7=(ABS(P5))** (1./3.)
   IF(P3)610,652,652
610  TETA=ACOS(P2/(-P1**3)**0.5)
   RQ1=2.*(-P1)**0.5*COS(TETA/3.)
   RQ2=2.*(-P1)**0.5*COS(TETA/3.+2.*PI/3.)
   RQ3=2.*(-P1)**0.5*COS(TETA/3.+4.*PI/3.)
   DSR1=A1*A1+4.*RQ1
   ESR1=RQ1*RQ1-4.*A4
   IF(DSR1)625,620,620
620  CONTINUE
   IF(ESR1)625,622,622
622  XR1=RQ1
   GO TO 664
625  DSR2=A1*A1+4.*RQ2
   ESR2=RQ2*RQ2-4.*A4
   IF(DSR2)635,630,630
630  CONTINUE
   IF(ESR2)635,632,632

```

```

632 XR1=RQ2
    GO TO 664
635 DSR3=A1*A1+4.*RQ3
    ESR3=RQ3*RQ3-4.*A4
    IF(DSR3)999,640,640
640 CONTINUE
    IF(ESR3)999,642,642
642 XR1=RQ3
    GO TO 664
652 CONTINUE
    IF(P4)654,656,656
654 P6=P6
    GO TO 657
656 P6=P6
657 CONTINUE
    IF(P5)658,660,660
658 P7=P7
    GO TO 661
660 P7=P7
661 CONTINUE
    XR1=P6+P7
664 DIS=A1*A1+4.*XR1
    RIS=XR1*XR1-4.*A4
    IF(DIS)999,666,666
666 CONTINUE
    IF(RIS)999,667,667
667 CONTINUE
    P8=(A1-SQRT(DIS))/2.
    P9=(XR1+SQRT(RIS))/2.
    P10=(A1+SQRT(DIS))/2.
    P11=(XR1-SQRT(RIS))/2.
    DIS1=P8*P8-4.*P9
    P12=(ABS(DIS1))*0.5
    IF(DIS1)681,670,670
670 ZZ1=(-P8+P12)/2.
    ZZ2=(-P8-P12)/2.
    IF(ZZ1)673,673,671
671 CONTINUE
    IF(ZZ1-RCC)672,673,673
672 ROC=ZZ1
    GO TO 602
673 CONTINUE
    IF(ZZ2)681,681,674
674 CONTINUE
    IF(ZZ2-RCC)676,681,681
676 ROC=ZZ2
    GO TO 602
681 DIS2=P10*P10-4.*P11
    P13=(ABS(DIS2))*0.5
    IF(DIS2)999,682,682
682 ZZ3=(-P10+P13)/2.
    ZZ4=(-P10-P13)/2.
    IF(ZZ3)685,685,683
683 CONTINUE
    IF(ZZ3-RCC)684,685,685
684 ROC=ZZ3
    GO TO 602
685 CONTINUE
    IF(ZZ4)999,999,686
686 CONTINUE
    IF(ZZ4-RCC)687,999,999
687 ROC=ZZ4
    GO TO 602
602 CONTINUE
    BB=2.*ASIN(1.)
    NK=150
    HK1=NK-1
    R1=RO
    R2=RO
    R3=RO
    R4=RO
    U1=ASIN(FR*RC/R1)
    HH1=(BB-U1)/FLOAT(NK)
    SU1=FR*RC/R1
    CU1=SQRT(1.-SU1*SU1)
    SG1=R1*CU1+SQRT(RP*RP-R1*R1*SU1*SU1)
    FX1=EXP(-FP*EFT1*SG1)*SU1
    FP1=0.0
    DO 200 K9=1,NK1

```

```

X30=U1+FLOAT(X9)*HH1
X31=-SIN(X30)
X32=R1*COS(X30)+SQRT(RP*RP-R1*R1*X31*X31)
X33=EXP(-FP*EPT1*X32)*X31
PP1=PP1+X33
200 CONTINUE
XPP1=HH1*(PX1/2.+PP1)/(1.+CU1)
U2=ASIN(FR*RC/R2)
HH2=(BB-U2)/FLOAT(NK)
SU2=FR*RC/R2
CU2=SQRT(1.-SU2*SU2)
SG2=R2*CU2+SQRT(RP*RP-R2*R2*SU2*SU2)
PX2=EXP(-FP*EPT2*SG2)*SU2
PP2=0.0
DO 201 K10=1,NK1
X35=U2+FLOAT(K10)*HH2
X36=-SIN(X35)
X37=R2*COS(X35)+SQRT(RP*RP-R2*R2*X36*X36)
X38=EXP(-FP*EPT2*X37)*X36
PP2=PP2+X38
201 CONTINUE
XPP2=HH2*(PX2/2.+PP2)/(1.+CU2)
U3=ASIN(FR*RC/R3)
HH3=(BB-U3)/FLOAT(NK)
SU3=FR*RC/R3
CU3=SQRT(1.-SU3*SU3)
SG3=R3*CU3+SQRT(RP*RP-R3*R3*SU3*SU3)
PX3=EXP(-FP*EPT3*SG3)*SU3
PP3=0.0
DO 202 K11=1,NK1
X40=U3+FLOAT(K11)*HH3
X41=-SIN(X40)
X42=R3*COS(X40)+SQRT(RP*RP-R3*R3*X41*X41)
X43=EXP(-FP*EPT3*X42)*X41
PP3=PP3+X43
202 CONTINUE
XPP3=HH3*(PX3/2.+PP3)/(1.+CU3)
U4=ASIN(FR*RC/R4)
HH4=(BB-U4)/FLOAT(NK)
SU4=FR*RC/R4
CU4=SQRT(1.-SU4*SU4)
SG4=R4*CU4+SQRT(RP*RP-R4*R4*SU4*SU4)
PX4=EXP(-FP*EPT4*SG4)*SU4
PP4=0.0
DO 203 K12=1,NK1
X45=U4+FLOAT(K12)*HH4
X46=-SIN(X45)
X47=R4*COS(X45)+SQRT(RP*RP-R4*R4*X46*X46)
X48=EXP(-FP*EPT4*X47)*X46
PP4=PP4+X48
203 CONTINUE
XPP4=HH4*(PX4/2.+PP4)/(1.+CU4)
HL1=U1/FLOAT(NK)
SQ1=R1*CU1
PX5=EXP(-FP*EPT1*SQ1)*SU1
PP5=0.0
DO 204 K13=1,NK1
X50=A+FLOAT(K13)*HL1
X51=-SIN(X50)
X52=R1*COS(X50)-SQRT(RCC*RCC-R1*R1*X51*X51)
X53=EXP(-FP*EPT1*X52)*X51
PP5=PP5+X53
204 CONTINUE
YP11=HL1*(PX5/2.+PP5)/(1.-CU1)
HL2=U2/FLOAT(NK)
SQ2=R2*CU2
PX6=EXP(-FP*EPT2*SQ2)*SU2
PP6=0.0
DO 205 K14=1,NK1
X55=A+FLOAT(K14)*HL2
X56=-SIN(X55)
X57=R2*COS(X55)-SQRT(RCC*RCC-R2*R2*X56*X56)
X58=EXP(-FP*EPT2*X57)*X56
PP6=PP6+X58
205 CONTINUE
YP12=HL2*(PX6/2.+PP6)/(1.-CU2)
HL3=U3/FLOAT(NK)
SQ3=R3*CU3
PX7=EXP(-FP*EPT3*SQ3)*SU3

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PP7=0.0
DO 206 K15=1,NK1
X60=A+FLOAT(K15)*HL3
X61=SIN(X60)
X62=R3*COS(X60)-SQRT(RCC*RCC-R3*R3*X61*X61)
X63=EXP(-FP*EPT3*X62)*X61
PP7=PP7+X63
206 CONTINUE
YP13=HL3*(PX7/2.+PP7)/(1.-CU3)
HL4=U4/FLOAT(NK)
SQ4=R4*CU4
PX8=EXP(-FP*EPT4*SQ4)*SU4
PP8=0.0
DO 207 K16=1,NK1
X65=A+FLOAT(K16)*HL4
X66=SIN(X65)
X67=R4*COS(X65)-SQRT(RCC*RCC-R4*R4*X66*X66)
X68=EXP(-FP*EPT4*X67)*X66
PP8=PP8+X68
207 CONTINUE
YP14=HL4*(PX8/2.+PP8)/(1.-CU4)
SMNO=SQRT(RP*RP-RCC*RCC)
PX9=EXP(-FP*EPT1*SMNO)*SU1
PP9=0.0
DO 208 K17=1,NK1
X70=A+FLOAT(K17)*HL1
X71=SIN(X70)
X72=SQRT(RP*RP-R1*R1*X71*X71)-SQRT(RCC*RCC-R1*R1*X71*X71)
X73=EXP(-FP*EPT1*X72)*X71
PP9=PP9+X73
208 CONTINUE
YP21=HL1*(PX9/2.+PP9)/(1.-CU1)
PX10=EXP(-FP*EPT2*SMNO)*SU2
PP10=0.0
DO 209 K18=1,NK1
X75=A+FLOAT(K18)*HL2
X76=SIN(X75)
X77=SQRT(RP*RP-R2*R2*X76*X76)-SQRT(RCC*RCC-R2*R2*X76*X76)
X78=EXP(-FP*EPT2*X77)*X76
PP10=PP10+X78
209 CONTINUE
YP22=HL2*(PX10/2.+PP10)/(1.-CU2)
PX11=EXP(-FP*EPT3*SMNO)*SU3
PP11=0.0
DO 210 K19=1,NK1
X80=A+FLOAT(K19)*HL3
X81=SIN(X80)
X82=SQRT(RP*RP-R3*R3*X81*X81)-SQRT(RCC*RCC-R3*R3*X81*X81)
X83=EXP(-FP*EPT3*X82)*X81
PP11=PP11+X83
210 CONTINUE
YP23=HL3*(PX11/2.+PP11)/(1.-CU3)
PX12=EXP(-FP*EPT4*SMNO)*SU4
PP12=0.0
DO 211 K20=1,NK1
X85=A+FLOAT(K20)*HL4
X86=SIN(X85)
X87=SQRT(RP*RP-R4*R4*X86*X86)-SQRT(RCC*RCC-R4*R4*X86*X86)
X88=EXP(-FP*EPT4*X87)*X86
PP12=PP12+X88
211 CONTINUE
YP24=HL4*(PX12/2.+PP12)/(1.-CU4)
PX13=SU1
PP13=0.0
DO 212 K21=1,NK1
X90=A+FLOAT(K21)*HL1
X91=SIN(X90)
X92=FP*EPT1*2.*SQRT(RCC*RCC-R1*R1*X91*X91)
X93=EXP(+X92)*X91
PP13=PP13+X93
212 CONTINUE
YCC1=HL1*(PX13/2.+PP13)/(1.-CU1)
PX14=SU2
PP14=0.0
DO 213 K22=1,NK1
X95=A+FLOAT(K22)*HL2
X96=SIN(X95)
X97=FP*EPT2*2.*SQRT(RCC*RCC-R2*R2*X96*X96)
X98=EXP(+X97)*X96

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PP14=PP14+X98
213 CONTINUE
YCC2=HL2*(PX14/2.+PP14)/(1.-CU2)
PX15=SU3
PP15=0.0
DO 214 K23=1,NK1
X100=A+FLOAT(K23)*HL3
X101=SIN(X100)
X102=-FF*ET3*2.*SQRT(RCC*RCC-R3*R3*X101*X101)
X103=EXP(+X102)*X101
PP15=PP15+X103
214 CONTINUE
YCC3=HL3*(PX15/2.+PP15)/(1.-CU3)
PX16=SU4
PP16=0.0
DO 215 K24=1,NK1
X105=A+FLOAT(K24)*HL4
X106=SIN(X105)
X107=-FF*ET4*2.*SQRT(RCC*RCC-R4*R4*X106*X106)
X108=EXP(+X107)*X106
PP16=PP16+X108
215 CONTINUE
YCC4=HL4*(PX16/2.+PP16)/(1.-CU4)
I1=180
HI=2.*ASIN(1.)/FLOAT(I1)
I2=I1-1
XI1=0.0
DO 220 K1=1,I2
Y10=FLOAT(K1)*HI
Y11=SIN(Y10)
Y12=-FF*ET1*(ROC*COS(Y10)+SQRT(RCC*RCC-ROC*ROC*Y11*Y11))
Y13=EXP(Y12)*Y11
XI1=XI1+Y13
220 CONTINUE
TC1=HI*XI1/2.
XI2=0.0
DO 222 K2=1,I2
Y15=FLOAT(K2)*HI
Y16=SIN(Y15)
Y17=-FF*ET2*(ROC*COS(Y15)+SQRT(RCC*RCC-ROC*ROC*Y16*Y16))
Y18=EXP(Y17)*Y16
XI2=XI2+Y18
222 CONTINUE
TC2=HI*XI2/2.
XI3=0.0
DO 224 K3=1,I2
Y20=FLOAT(K3)*HI
Y21=SIN(Y20)
Y22=-FF*ET3*(ROC*COS(Y20)+SQRT(RCC*RCC-ROC*ROC*Y21*Y21))
Y23=EXP(Y22)*Y21
XI3=XI3+Y23
224 CONTINUE
TC3=HI*XI3/2.
XI4=0.0
DO 226 K4=1,I2
Y25=FLOAT(K4)*HI
Y26=SIN(Y25)
Y27=-FF*ET4*(ROC*COS(Y25)+SQRT(RCC*RCC-ROC*ROC*Y26*Y26))
Y28=EXP(Y27)*Y26
XI4=XI4+Y28
226 CONTINUE
TC4=HI*XI4/2.
XI5=0.0
DO 228 K5=1,I2
Y30=FLOAT(K5)*HI
Y31=SIN(Y30)
Y32=SQRT(RP*RP-ROC*ROC*Y31*Y31)-SQRT(RCC**2-ROC**2*Y31**2)
Y33=EXP(-FF*ET1*Y32)*Y31
XI5=XI5+Y33
228 CONTINUE
TP1=HI*XI5/2.
XI6=0.0
DO 230 K6=1,I2
Y35=FLOAT(K6)*HI
Y36=SIN(Y35)
Y37=SQRT(RP*RP-ROC*ROC*Y36*Y36)-SQRT(RCC**2-ROC**2*Y36**2)
Y38=EXP(-FF*ET2*Y37)*Y36
XI6=XI6+Y38
230 CONTINUE

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TP2=HI*XI6/2.
XI7=0.0
DO 232 K7=1,I2
Y40=FLOAT(K7)*HI
Y41=SIN(Y40)
Y42=SQRT(RP*RP-ROC*ROC*Y41*Y41)-SQRT(RCC**2-ROC**2*Y41**2)
Y43=EXP(-FP*EPT3*Y42)*Y41
XI7=XI7+Y43
232 CONTINUE
TP3=HI*XI7/2.
XI8=0.0
DO 234 K8=1,I2
Y45=FLOAT(K8)*HI
Y46=SIN(Y45)
Y47=SQRT(RP*RP-ROC*ROC*Y46*Y46)-SQRT(RCC**2-ROC**2*Y46**2)
Y48=EXP(-FP*EPT4*Y47)*Y46
XI8=XI8+Y48
234 CONTINUE
TP4=HI*XI8/2.
GC1=(1.-CU1)/2.
GP1=1.-GC1
GC2=(1.-CU2)/2.
GP2=1.-GC2
GC3=(1.-CU3)/2.
GP3=1.-GC3
GC4=(1.-CU4)/2.
GP4=1.-GC4
236 IF(RDEL) 500,500,236
RCL=RCC+RDEL
RCS=RCC-RDEL
RH=RO
PI=3.14159
A7=(RCC+RCS)/(2.*RH)
A8=ASIN(A7)
A9=COS(A8)
A10=1.-A9
GP1=A10/2.
A11=RCC*COS(PI/8.)
B11=RCC*SIN(PI/8.)
A12=RCS*SIN(PI/4.)
B12=A12
A13=RCC*COS(3.*PI/8.)
B13=RCC*SIN(3.*PI/8.)
A14=(RH+A11)**2+B11*B11
A15=SQRT(A14)
A16=(RH+A12)**2+B12*B12
A17=SQRT(A16)
A18=(RH+A13)**2+B13*B13
A19=SQRT(A18)
A20=SQRT(RH*RH+RCL*RCL)
A21=ASIN(B11/A15)
A22=ASIN(B12/A17)
A23=ASIN(B13/A19)
A24=ASIN(RCL/A20)
A25=B11/A15
A26=B12/A17
A27=B13/A19
A28=RCL/A20
A30=COS(A21)
A31=COS(A22)
A32=COS(A23)
A33=COS(A24)
A35=RH*A30+SQRT(RP*RP-RH*RH*A25*A25)
A36=RH*A31+SQRT(RP*RP-RH*RH*A26*A26)
A37=RH*A32+SQRT(RP*RP-RH*RH*A27*A27)
A38=RH*A33+SQRT(RP*RP-RH*RH*A28*A28)
A39=RH*A9+SQRT(RP*RP-RH*RH*A7*A7)
A40=RH*A30-SQRT(RCL*RCL-RH*RH*A25*A25)
A41=RH*A31-SQRT(RCL*RCL-RH*RH*A26*A26)
A42=RH*A32-SQRT(RCL*RCL-RH*RH*A27*A27)
A43=(RCL+RCC)/2.
A44=RH*A33-SQRT(A43*A43-RH*RH*A28*A28)
AA44=SQRT(B11*B11+(RH-A11)*(RH-A11))
A45=EXP(-EPT1*A40)*A25
A46=EXP(-EPT1*A41)*A26
A47=EXP(-EPT1*A42)*A27
A48=EXP(-EPT1*A44)*A28
A49=EXP(-EPT1*AA44)*A7
A50=EXP(-EPT2*A40)*A25

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A51=EXP(-EPT2*A41)*A26
A52=EXP(-EPT2*A42)*A27
A53=EXP(-EPT2*A44)*A28
A54=EXP(-EPT2*AA44)*A7
A55=EXP(-EPT3*A40)*A25
A56=EXP(-EPT3*A41)*A26
A57=EXP(-EPT3*A42)*A27
A58=EXP(-EPT3*A44)*A28
A59=EXP(-EPT3*AA44)*A7
A60=EXP(-EPT4*A40)*A25
A61=EXP(-EPT4*A41)*A26
A62=EXP(-EPT4*A42)*A27
A63=EXP(-EPT4*A44)*A28
A64=EXP(-EPT4*AA44)*A7
A65=(A45*A21+(A45+A46)*(A22-A21)+(A46+A47)*(A23-A22)+(
#A47+A48)*(A24-A23)+(A48+A49)*(A8-A24))/2.
TIH1=A65/A10
A66=(A50*A21+(A50+A51)*(A22-A21)+(A51+A52)*(A23-A22)+(
#A52+A53)*(A24-A23)+(A53+A54)*(A8-A24))/2.
TIH2=A66/A10
A67=(A55*A21+(A55+A56)*(A22-A21)+(A56+A57)*(A23-A22)+(
#A57+A58)*(A24-A23)+(A58+A59)*(A8-A24))/2.
TIH3=A67/A10
A68=(A60*A21+(A60+A61)*(A22-A21)+(A61+A62)*(A23-A22)+(
#A62+A63)*(A24-A23)+(A63+A64)*(A8-A24))/2.
TIH4=A68/A10
A70=A15-A40
A71=A17-A41
A72=A19-A42
A73=A20-A44
A75=EXP(-ET1*A70)*A25
A76=EXP(-ET1*A71)*A26
A77=EXP(-ET1*A72)*A27
A78=EXP(-ET1*A73)*A28
A80=EXP(-ET2*A70)*A25
A81=EXP(-ET2*A71)*A26
A82=EXP(-ET2*A72)*A27
A83=EXP(-ET2*A73)*A28
A85=EXP(-ET3*A70)*A25
A86=EXP(-ET3*A71)*A26
A87=EXP(-ET3*A72)*A27
A88=EXP(-ET3*A73)*A28
A90=EXP(-ET4*A70)*A25
A91=EXP(-ET4*A71)*A26
A92=EXP(-ET4*A72)*A27
A93=EXP(-ET4*A73)*A28
A95=(A75*A21+(A75+A76)*(A22-A21)+(A76+A77)*(A23-A22)+(
#A77+A78)*(A24-A23)+(A78+A79)*(A8-A24))/2.
TIH1=A95/A10
A96=(A80*A21+(A80+A81)*(A22-A21)+(A81+A82)*(A23-A22)+(
#A82+A83)*(A24-A23)+(A83+A79)*(A8-A24))/2.
TIH2=A96/A10
A97=(A85*A21+(A85+A86)*(A22-A21)+(A86+A87)*(A23-A22)+(
#A87+A88)*(A24-A23)+(A88+A79)*(A8-A24))/2.
TIH3=A97/A10
A98=(A90*A21+(A90+A91)*(A22-A21)+(A91+A92)*(A23-A22)+(
#A92+A93)*(A24-A23)+(A93+A79)*(A8-A24))/2.
TIH4=A98/A10
A100=A35-A15
A101=A36-A17
A102=A37-A19
A103=A38-A20
A104=A39-AA44
A106=EXP(-EPT1*A100)*A25
A107=EXP(-EPT1*A101)*A26
A108=EXP(-EPT1*A102)*A27
A109=EXP(-EPT1*A103)*A28
A110=EXP(-EPT1*A104)*A7
A112=EXP(-EPT2*A100)*A25
A113=EXP(-EPT2*A101)*A26
A114=EXP(-EPT2*A102)*A27
A115=EXP(-EPT2*A103)*A28
A116=EXP(-EPT2*A104)*A7
A118=EXP(-EPT3*A100)*A25
A119=EXP(-EPT3*A101)*A26
A120=EXP(-EPT3*A102)*A27
A121=EXP(-EPT3*A103)*A28
A122=EXP(-EPT3*A104)*A7
A124=EXP(-EPT4*A100)*A25

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A125=EXP(-EPT4*A101)*A26
A126=EXP(-EPT4*A102)*A27
A127=EXP(-EPT4*A103)*A28
A128=EXP(-EPT4*A104)*A7
A130=(A106*A21+(A106+A107)*(A22-A21)+(A107+A108)*(A23-A22)
+ (A108+A109)*(A24-A23)+(A109+A110)*(A8-A24))/2.
TZH1=A130/A10
A131=(A112*A21+(A112+A113)*(A22-A21)+(A113+A114)*(A23-A22)
+ (A114+A115)*(A24-A23)+(A115+A116)*(A8-A24))/2.
TZH2=A131/A10
A132=(A118*A21+(A118+A119)*(A22-A21)+(A119+A120)*(A23-A22)
+ (A120+A121)*(A24-A23)+(A121+A122)*(A8-A24))/2.
TZH3=A132/A10
A133=(A124*A21+(A124+A125)*(A22-A21)+(A125+A126)*(A23-A22)
+ (A126+A127)*(A24-A23)+(A127+A128)*(A8-A24))/2.
TZH4=A133/A10
A138=1.+A9
NNK=150
NNK1=NNK-1
HY1=(BB-A8)/FLOAT(NNK)
A139=EXP(-EPT1*A39)*A7
A140=EXP(-EPT2*A39)*A7
A141=EXP(-EPT3*A39)*A7
A142=EXP(-EPT4*A39)*A7
A143=0.0
A144=0.0
A145=0.0
A146=0.0
DO 304 K30=1, NNK1
A148=A8+FLOAT(K30)*HY1
A149=SIN(A148)
A150=COS(A148)
A151=RH*A150+SQRT(RP*RP-RH*RH*A149*A149)
A152=EXP(-EPT1*A151)*A149
A153=EXP(-EPT2*A151)*A149
A154=EXP(-EPT3*A151)*A149
A155=EXP(-EPT4*A151)*A149
A143=A143+A152
A144=A144+A153
A145=A145+A154
A146=A146+A155
304 CONTINUE
THH1=HY1*(A143+A139/2.)/A138
THH2=HY1*(A144+A140/2.)/A138
THH3=HY1*(A145+A141/2.)/A138
THH4=HY1*(A146+A142/2.)/A138
A200=-RCL*RCL/RH
A201=SQRT(-A200*(RH+A200))*RCC/RCL
A202=SQRT(A201*A201+(RH+A200)*(RH+A200))
A203=ASIN(A201/A202)
A204=SIN(A203)
A205=COS(A203)
A206=RH*A205+SQRT(RP*RP-RH*RH*A204*A204)
A207=A206-A202
A208=1.-A205
A209=1.+A205
GF2=A208/2.
A210=EXP(-EPT1*A202)*A204
A211=EXP(-EPT2*A202)*A204
A212=EXP(-EPT3*A202)*A204
A213=EXP(-EPT4*A202)*A204
A215=0.0
A216=0.0
A217=0.0
A218=0.0
HY2=A203/FLOAT(NNK)
DO 308 K33=1, NNK1
A220=FLOAT(K33)*HY2
A221=SIN(A220)
A222=COS(A220)
A223=A221/A222
A224=RH*A223
A225=A223**2+(RCC/RCL)**2
A226=(A224-RCC)*(A224+RCC)
A227=(2.*A223*A224)**2-4.*A225*A226
A228=SQRT(A227)
A229=(-2.*A223*A224-A228)/(2.*A225)
A230=(RH+A229)/A222
A231=EXP(-EPT1*A230)*A221

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A232=EXP(-EPT2*A230)*A221
A233=EXP(-EPT3*A230)*A221
A234=EXP(-EPT4*A230)*A221
A215=A215+A231
A216=A216+A232
A217=A217+A233
A218=A218+A234
308 CONTINUE
E1H1=HY2*(A215+A210/2.)/A208
E1H2=HY2*(A216+A211/2.)/A208
E1H3=HY2*(A217+A212/2.)/A208
E1H4=HY2*(A218+A213/2.)/A208
A288=0.0
A289=0.0
A290=0.0
A291=0.0
DO 316 K37=1,NNK1
A292=FLOAT(K37)*HY2
A293=SIN(A292)
A294=COS(A292)
A295=A293/A294
A296=RH*A295
A297=A295**2+(RCC/RCL)**2
A298=(A296-RCC)*(A296+RCC)
A299=(2.*A295*A296)**2-4.*A297*A298
A300=SQRT(A299)
A301=A300/A297
A302=A301/A294
A303=EXP(-EPT1*A302)*A293
A304=EXP(-EPT2*A302)*A293
A305=EXP(-EPT3*A302)*A293
A306=EXP(-EPT4*A302)*A293
A288=A288+A303
A289=A289+A304
A290=A290+A305
A291=A291+A306
316 CONTINUE
E1F1=HY2*(A288+A204/2.)/A208
E1F2=HY2*(A289+A204/2.)/A208
E1F3=HY2*(A290+A204/2.)/A208
E1F4=HY2*(A291+A204/2.)/A208
A364=EXP(-EPT1*A207)*A204
A365=EXP(-EPT2*A207)*A204
A366=EXP(-EPT3*A207)*A204
A367=EXP(-EPT4*A207)*A204
A368=0.0
A369=0.0
A370=0.0
DO 324 K38=1,NNK1
A372=FLOAT(K38)*HY2
A373=SIN(A372)
A374=COS(A372)
A375=A373/A374
A376=RH*A375
A377=A375**2+(RCC/RCL)**2
A378=(A376-RCC)*(A376+RCC)
A379=(2.*A375*A376)**2-4.*A377*A378
A380=SQRT(A379)
A381=(-2.*A375*A376+A380)/(2.*A377)
A382=(RH+A381)/A374
A383=RH*A374+SQRT(RP*RP-RH*RH*A373*A373)
A384=A383-A382
A385=EXP(-EPT1*A384)*A373
A386=EXP(-EPT2*A384)*A373
A387=EXP(-EPT3*A384)*A373
A388=EXP(-EPT4*A384)*A373
A368=A368+A385
A369=A369+A386
A370=A370+A387
A371=A371+A388
324 CONTINUE
E2H1=HY2*(A368+A364/2.)/A208
E2H2=HY2*(A369+A365/2.)/A208
E2H3=HY2*(A370+A366/2.)/A208
E2H4=HY2*(A371+A367/2.)/A208
A400=EXP(-EPT1*A206)*A204
A401=EXP(-EPT2*A206)*A204
A402=EXP(-EPT3*A206)*A204

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A403=EXP(-EPT4*A206)*A204
A404=0.0
A405=0.0
A406=0.0
A407=0.0
HY3=(PI-A203)/FLOAT(NNK)
DO 326 K39=1,NNK1
A408=A203+FLOAT(K39)*HY3
A409=SIN(A408)
A410=COS(A408)
A411=RH*A410+SQRT(RP*RP-RH*RH*A409*A409)
A412=EXP(-EPT1*A411)*A409
A413=EXP(-EPT2*A411)*A409
A414=EXP(-EPT3*A411)*A409
A415=EXP(-EPT4*A411)*A409
A404=A404+A412
A405=A405+A413
A406=A406+A414
A407=A407+A415
326 CONTINUE
EHH1=HY3*(A404+A400/2.)/A209
EHH2=HY3*(A405+A401/2.)/A209
EHH3=HY3*(A406+A402/2.)/A209
EHH4=HY3*(A407+A403/2.)/A209
A500=-RCC*RCC/RH
A501=SQRT(-A500*(RH+A500))*RCL/RCC
A502=SQRT(A501*A501+(RH+A500)*(RH+A500))
A503=ASIN(A501/A502)
A504=SIN(A503)
A505=COS(A503)
A506=RH*A505+SQRT(RP*RP-RH*RH*A504*A504)
A507=A506-A502
A508=1.-A505
A509=1.+A505
GF3=A508/2.
A510=EXP(-EPT1*A502)*A504
A511=EXP(-EPT2*A502)*A504
A512=EXP(-EPT3*A502)*A504
A513=EXP(-EPT4*A502)*A504
A515=0.0
A516=0.0
A517=0.0
A518=0.0
HY4=A503/FLOAT(NNK)
DO 330 K40=1,NNK1
A520=FLOAT(K40)*HY4
A521=SIN(A520)
A522=COS(A520)
A523=A521/A522
A524=RH*A523
A525=A523**2+(RCL/RCC)**2
A526=(A524-RCL)*(A524+RCL)
A527=(2.*A523*A524)**2-4.*A525*A526
A528=SQRT(A527)
A529=(-2.*A523*A524-A528)/(2.*A525)
A530=(RH+A529)/A522
A531=EXP(-EPT1*A530)*A521
A532=EXP(-EPT2*A530)*A521
A533=EXP(-EPT3*A530)*A521
A534=EXP(-EPT4*A530)*A521
A515=A515+A531
A516=A516+A532
A517=A517+A533
A518=A518+A534
330 CONTINUE
V1H1=HY4*(A515+A510/2.)/A508
V1H2=HY4*(A516+A511/2.)/A508
V1H3=HY4*(A517+A512/2.)/A508
V1H4=HY4*(A518+A513/2.)/A508
A540=0.0
A541=0.0
A542=0.0
A543=0.0
DO 332 K41=1,NNK1
A545=FLOAT(K41)*HY4
A546=SIN(A545)
A547=COS(A545)
A548=A546/A547
A549=RH*A548

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A550=A548**2+(RCL/RCC)**2
A551=(A549-RCL)*(A549+RCL)
A552=(2.*A548*A549)**2-4.*A550*A551
A553=SQR(A552)
A554=A553/A550
A555=A554/A547
A556=EXP(-EPT1*A555)*A546
A557=EXP(-EPT2*A555)*A546
A558=EXP(-EPT3*A555)*A546
A559=EXP(-EPT4*A555)*A546
A540=A540+A556
A541=A541+A557
A542=A542+A558
A543=A543+A559
332 CONTINUE
V1F1=HY4*(A540+A504/2.)/A508
V1F2=HY4*(A541+A504/2.)/A508
V1F3=HY4*(A542+A504/2.)/A508
V1F4=HY4*(A543+A504/2.)/A508
A564=EXP(-EPT1*A507)*A504
A565=EXP(-EPT2*A507)*A504
A566=EXP(-EPT3*A507)*A504
A567=EXP(-EPT4*A507)*A504
A568=0.0
A569=0.0
A570=0.0
A571=0.0
DO 334 K42=1,NNK1
A572=FLOAT(K42)*HY4
A573=SIN(A572)
A574=COS(A572)
A575=A573/A574
A576=RH*A575
A577=A575**2+(RCL/RCC)**2
A578=(A576-RCL)*(A576+RCL)
A579=(2.*A575*A576)**2-4.*A577*A578
A580=SQR(A579)
A581=(-2.*A575*A576+A580)/(2.*A577)
A582=(RH+A581)/A574
A583=RH*A574+SQR(RP*RP-RH*RH*A573*A573)
A584=A583-A582
A585=EXP(-EPT1*A584)*A573
A586=EXP(-EPT2*A584)*A573
A587=EXP(-EPT3*A584)*A573
A588=EXP(-EPT4*A584)*A573
A568=A568+A585
A569=A569+A586
A570=A570+A587
A571=A571+A588
334 CONTINUE
V2H1=HY4*(A568+A564/2.)/A508
V2H2=HY4*(A569+A565/2.)/A508
V2H3=HY4*(A570+A566/2.)/A508
V2H4=HY4*(A571+A567/2.)/A508
A600=EXP(-EPT1*A506)*A504
A601=EXP(-EPT2*A506)*A504
A602=EXP(-EPT3*A506)*A504
A603=EXP(-EPT4*A506)*A504
A604=0.0
A605=0.0
A606=0.0
A607=0.0
HY5=(PI-A503)/FLOAT(NNK1)
DO 336 K43=1,NNK1
A608=A503+FLOAT(K43)*HY5
A609=SIN(A608)
A610=COS(A608)
A611=RH*A610+SQR(RP*RP-RH*RH*A609*A609)
A612=EXP(-EPT1*A611)*A609
A613=EXP(-EPT2*A611)*A609
A614=EXP(-EPT3*A611)*A609
A615=EXP(-EPT4*A611)*A609
A604=A604+A612
A605=A605+A613
A606=A606+A614
A607=A607+A615
336 CONTINUE
VHH1=HY5*(A604+A600/2.)/A509
VHH2=HY5*(A605+A601/2.)/A509

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VHH3=HY5*(A606+A602/2.)/A509
VHH4=HY5*(A607+A603/2.)/A509
A700=-RCC*RCC/RH
A701=SQRT(-A700*(RH+A700))*RCS/RCC
A702=SQRT(A701*A701+(RH+A700)*(RH+A700))
A703=ASIN(A701/A702)
A704=SIN(A703)
A705=COS(A703)
A706=RH*A705+SQRT(RP*RP-RH*RH*A704*A704)
A707=A706-A702
A708=1.-A705
A709=1.+A705
GP4=A708/2.
A710=EXP(-EPT1*A702)*A704
A711=EXP(-EPT2*A702)*A704
A712=EXP(-EPT3*A702)*A704
A713=EXP(-EPT4*A702)*A704
A715=0.0
A716=0.0
A717=0.0
A718=0.0
HY6=A703/FLOAT(NNK)
DO 338 K4=1,NNK1
A720=FLOAT(K44)*HY6
A721=SIN(A720)
A722=COS(A720)
A723=A721/A722
A724=RH*A723
A725=A723**2+(RCS/RCC)**2
A726=(A724-RCS)*(A724+RCS)
A727=(2.*A723*A724)**2-4.*A725*A726
A728=SQRT(A727)
A729=(-2.*A723*A724-A728)/(2.*A725)
A730=(RH+A729)/A722
A731=EXP(-EPT1*A730)*A721
A732=EXP(-EPT2*A730)*A721
A733=EXP(-EPT3*A730)*A721
A734=EXP(-EPT4*A730)*A721
A715=A715+A731
A716=A716+A732
A717=A717+A733
A718=A718+A734
338 CONTINUE
F1H1=HY6*(A715+A710/2.)/A708
F1H2=HY6*(A716+A711/2.)/A708
F1H3=HY6*(A717+A712/2.)/A708
F1H4=HY6*(A718+A713/2.)/A708
A740=0.0
A741=0.0
A742=0.0
A743=0.0
DO 340 K45=1,NNK1
A745=FLOAT(K45)*HY6
A746=SIN(A745)
A747=COS(A745)
A748=A746/A747
A749=RH*A748
A750=A748**2+(RCS/RCC)**2
A751=(A749-RCS)*(A749+RCS)
A752=(2.*A748*A749)**2-4.*A750*A751
A753=SQRT(A752)
A754=A753/A750
A755=A754/A747
A756=EXP(-ET1*A755)*A746
A757=EXP(-ET2*A755)*A746
A758=EXP(-ET3*A755)*A746
A759=EXP(-ET4*A755)*A746
A740=A740+A756
A741=A741+A757
A742=A742+A758
A743=A743+A759
340 CONTINUE
F1F1=HY6*(A740+A704/2.)/A708
F1F2=HY6*(A741+A704/2.)/A708
F1F3=HY6*(A742+A704/2.)/A708
F1F4=HY6*(A743+A704/2.)/A708
A764=EXP(-EPT1*A707)*A704
A765=EXP(-EPT2*A707)*A704
A766=EXP(-EPT3*A707)*A704

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A767=EXP(-EPT4*A707)*A704
A768=0.0
A769=0.0
A770=0.0
A771=0.0
DO 342 K46=1,NNK1
A772=FLOAT(K46)*HY6
A773=SIN(A772)
A774=COS(A772)
A775=A773/A774
A776=RH*A775
A777=A775**2+(RCS/RCC)**2
A778=(A776-RCS)*(A776+RCS)
A779=(2.*A775*A776)**2-4.*A777*A778
A780=SQRT(A779)
A781=(-2.*A775*A776+A780)/(2.*A777)
A782=(RH+A781)/A774
A783=RH*A774+SQRT(RP*RP-RH*RH*A773*A773)
A784=A783-A782
A785=EXP(-EPT1*A784)*A773
A786=EXP(-EPT2*A784)*A773
A787=EXP(-EPT3*A784)*A773
A788=EXP(-EPT4*A784)*A773
A768=A768+A785
A769=A769+A786
A770=A770+A787
A771=A771+A788
342 CONTINUE
FZH1=HY6*(A768+A764/2.)/A708
FZH2=HY6*(A769+A765/2.)/A708
FZH3=HY6*(A770+A766/2.)/A708
FZH4=HY6*(A771+A767/2.)/A708
A800=EXP(-EPT1*A706)*A704
A801=EXP(-EPT2*A706)*A704
A802=EXP(-EPT3*A706)*A704
A803=EXP(-EPT4*A706)*A704
A804=0.0
A805=0.0
A806=0.0
A807=0.0
HY7=(PI-A703)/FLOAT(NNK1)
DO 344 K47=1,NNK1
A808=A703+FLOAT(K47)*HY7
A809=SIN(A808)
A810=COS(A808)
A811=RH*A810+SQRT(RP*RP-RH*RH*A809*A809)
A812=EXP(-EPT1*A811)*A809
A813=EXP(-EPT2*A811)*A809
A814=EXP(-EPT3*A811)*A809
A815=EXP(-EPT4*A811)*A809
A804=A804+A812
A805=A805+A813
A806=A806+A814
A807=A807+A815
344 CONTINUE
FHH1=HY7*(A804+A800/2.)/A709
FHH2=HY7*(A805+A801/2.)/A709
FHH3=HY7*(A806+A802/2.)/A709
FHH4=HY7*(A807+A803/2.)/A709
B20=RCC*COS(PI/8.)
B21=RCC*SIN(PI/8.)
B22=RCL*COS(PI/4.)
B23=RCL*SIN(PI/4.)
B24=RCC*COS(3.*PI/8.)
B25=RCC*SIN(3.*PI/8.)
B26=RCS
B27=-RCC*COS(3.*PI/8.)
B28=RCC*SIN(3.*PI/8.)
B29=-RCL*COS(PI/4.)
B30=RCL*SIN(PI/4.)
B31=(RH+B20)**2+B21*B21
B32=SQRT(B31)
B33=(RH+B22)**2+B23*B23
B34=SQRT(B33)
B35=(RH+B24)**2+B25*B25
B36=SQRT(B35)
B37=SQRT(RH*RH+B26*B26)
B38=(RH+B37)**2+B28*B28
B39=SQRT(B38)

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B40=(RH*B29)**2+B30*B30
B41=SQRT(B40)
B42=ASIN(B21/B32)
B43=ASIN(B23/B34)
B44=ASIN(B25/B36)
B45=ASIN(B26/B37)
B46=ASIN(B28/B39)
B47=ASIN(B30/B41)
B48=B21/B32
B49=B23/B34
B50=B25/B36
B51=B26/B37
B52=B28/B39
B53=B30/B41
B54=COS(B42)
B55=COS(B43)
B56=COS(B44)
B57=COS(B45)
B58=COS(B46)
B59=COS(B47)
B60=1.-B59
B61=(RCC+RCS)/2.
B62=(RCC+RCL)/2.
B63=RH*B54-SQRT(B61*B61-RH*RH*B48*B48)
B64=RH*B55-SQRT(B61*B61-RH*RH*B49*B49)
B65=RH*B56-SQRT(B61*B61-RH*RH*B50*B50)
B66=RH*B57-SQRT(B62*B62-RH*RH*B51*B51)
B67=RH*B58-SQRT(B62*B62-RH*RH*B52*B52)
B68=RH*B54+SQRT(RP*RP-RH*RH*B48*B48)
B69=RH*B55+SQRT(RP*RP-RH*RH*B49*B49)
B70=RH*B56+SQRT(RP*RP-RH*RH*B50*B50)
B71=RH*B57+SQRT(RP*RP-RH*RH*B51*B51)
B72=RH*B58+SQRT(RP*RP-RH*RH*B52*B52)
B73=RH*B59+SQRT(RP*RP-RH*RH*B53*B53)
B75=EXP(-EPT1*B63)*B48
B76=EXP(-EPT1*B64)*B49
B77=EXP(-EPT1*B65)*B50
B78=EXP(-EPT1*B66)*B51
B79=EXP(-EPT1*B67)*B52
B80=EXP(-EPT1*B41)*B53
B81=EXP(-EPT2*B63)*B48
B82=EXP(-EPT2*B64)*B49
B83=EXP(-EPT2*B65)*B50
B84=EXP(-EPT2*B66)*B51
B85=EXP(-EPT2*B67)*B52
B86=EXP(-EPT2*B41)*B53
B87=EXP(-EPT3*B63)*B48
B88=EXP(-EPT3*B64)*B49
B89=EXP(-EPT3*B65)*B50
B90=EXP(-EPT3*B66)*B51
B91=EXP(-EPT3*B67)*B52
B92=EXP(-EPT3*B41)*B53
B93=EXP(-EPT4*B63)*B48
B94=EXP(-EPT4*B64)*B49
B95=EXP(-EPT4*B65)*B50
B96=EXP(-EPT4*B66)*B51
B97=EXP(-EPT4*B67)*B52
B98=EXP(-EPT4*B41)*B53
B100=(B75*B42+(B75+B76)*(B43-B42)+(B76+B77)*(B44-B43)+(
#B77+B78)*(B45-B44)+(B78+B79)*(B46-B45)+(B79+B80)*(B47-
#B46))/2.
Z1H1=B100/B60
B101=(B81*B42+(B81+B82)*(B43-B42)+(B82+B83)*(B44-B43)+(
#B83+B84)*(B45-B44)+(B84+B85)*(B46-B45)+(B85+B86)*(B47-
#B46))/2.
Z1H2=B101/B60
B102=(B87*B42+(B87+B88)*(B43-B42)+(B88+B89)*(B44-B43)+(
#B89+B90)*(B45-B44)+(B90+B91)*(B46-B45)+(B91+B92)*(B47-
#B46))/2.
Z1H3=B102/B60
B103=(B93*B42+(B93+B94)*(B43-B42)+(B94+B95)*(B44-B43)+(
#B95+B96)*(B45-B44)+(B96+B97)*(B46-B45)+(B97+B98)*(B47-
#B46))/2.
Z1H4=B103/B60
GF5=B60/2.
B105=B32-B63
B106=B34-B64
B107=B36-B65
B108=B37-B66

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B109=B39-B67
B110=EXP(-ET1*B105)*B48
B111=EXP(-ET1*B106)*B49
B112=EXP(-ET1*B107)*B50
B113=EXP(-ET1*B108)*B51
B114=EXP(-ET1*B109)*B52
B115=EXP(-ET2*B105)*B48
B116=EXP(-ET2*B106)*B49
B117=EXP(-ET2*B107)*B50
B118=EXP(-ET2*B108)*B51
B119=EXP(-ET2*B109)*B52
B120=EXP(-ET3*B105)*B48
B121=EXP(-ET3*B106)*B49
B122=EXP(-ET3*B107)*B50
B123=EXP(-ET3*B108)*B51
B124=EXP(-ET3*B109)*B52
B125=EXP(-ET4*B105)*B48
B126=EXP(-ET4*B106)*B49
B127=EXP(-ET4*B107)*B50
B128=EXP(-ET4*B108)*B51
B129=EXP(-ET4*B109)*B52
B130=(B110*B42+(B110+B111)*(B43-B42)+(B111+B112)*(B44-B43
)+(B112+B113)*(B45-B44)+(B113+B114)*(B46-B45)+(B114+B53)*
(B47-B46))/2.
Z1F1=B130/B60
B131=(B115*B42+(B115+B116)*(B43-B42)+(B116+B117)*(B44-
B43)+(B117+B118)*(B45-B44)+(B118+B119)*(B46-B45)+(
B119+B53)*(B47-B46))/2.
Z1F2=B131/B60
B132=(B120*B42+(B120+B121)*(B43-B42)+(B121+B122)*(B44-
B43)+(B122+B123)*(B45-B44)+(B123+B124)*(B46-B45)+(
B124+B53)*(B47-B46))/2.
Z1F3=B132/B60
B133=(B125*B42+(B125+B126)*(B43-B42)+(B126+B127)*(B44-
B43)+(B127+B128)*(B45-B44)+(B128+B129)*(B46-B45)+(
B129+B53)*(B47-B46))/2.
Z1F4=B133/B60
B135=B68-B32
B136=B69-B34
B137=B70-B36
B138=B71-B37
B139=B72-B39
B140=B73-B41
B141=EXP(-EPT1*B135)*B48
B142=EXP(-EPT1*B136)*B49
B143=EXP(-EPT1*B137)*B50
B144=EXP(-EPT1*B138)*B51
B145=EXP(-EPT1*B139)*B52
B146=EXP(-EPT1*B140)*B53
B147=EXP(-EPT2*B135)*B48
B148=EXP(-EPT2*B136)*B49
B149=EXP(-EPT2*B137)*B50
B150=EXP(-EPT2*B138)*B51
B151=EXP(-EPT2*B139)*B52
B152=EXP(-EPT2*B140)*B53
B154=EXP(-EPT3*B135)*B48
B155=EXP(-EPT3*B136)*B49
B156=EXP(-EPT3*B137)*B50
B157=EXP(-EPT3*B138)*B51
B158=EXP(-EPT3*B139)*B52
B159=EXP(-EPT3*B140)*B53
B160=EXP(-EPT4*B135)*B48
B161=EXP(-EPT4*B136)*B49
B162=EXP(-EPT4*B137)*B50
B163=EXP(-EPT4*B138)*B51
B164=EXP(-EPT4*B139)*B52
B165=EXP(-EPT4*B140)*B53
B166=(B141*B42+(B141+B142)*(B43-B42)+(B142+B143)*(B44-
B43)+(B143+B144)*(B45-B44)+(B144+B145)*(B46-B45)+(
B145+B146)*(B47-B46))/2.
Z2H1=B166/B60
B167=(B147*B42+(B147+B148)*(B43-B42)+(B148+B149)*(B44-
B43)+(B149+B150)*(B45-B44)+(B150+B151)*(B46-B45)+(
B151+B152)*(B47-B46))/2.
Z2H2=B167/B60
B168=(B154*B42+(B154+B155)*(B43-B42)+(B155+B156)*(B44-
B43)+(B156+B157)*(B45-B44)+(B157+B158)*(B46-B45)+(
B158+B159)*(B47-B46))/2.
Z2H3=B168/B60

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B169=(B160*B42+(B160+B161)*(B43-B42)+(B161+B162)*(B44-
#B43)+(B162+B163)*(B45-B44)+(B163+B164)*(B46-B45)+(
#B164+B165)*(B47-B46))/2.
Z2H4=B169/B60
B170=1.+B59
HJ1=(PI-B47)/FLOAT(NNK)
B171=EXP(-EPT1*B73)*B53
B172=EXP(-EPT2*B73)*B53
B173=EXP(-EPT3*B73)*B53
B174=EXP(-EPT4*B73)*B53
B175=0.0
B176=0.0
B177=0.0
B178=0.0
DO 346 K50=1,NNK1
B179=B47+FLOAT(K50)*HJ1
B180=SIN(B179)
B181=COS(B179)
B182=RH*B181+SQRT(RP*RP-RH*RH*B180*B180)
B183=EXP(-EPT1*B182)*B180
B184=EXP(-EPT2*B182)*B180
B185=EXP(-EPT3*B182)*B180
B186=EXP(-EPT4*B182)*B180
B175=B175+B183
B176=B176+B184
B177=B177+B185
B178=B178+B186
346 CONTINUE
ZHH1=HJ1*(B175+B171/2.)/B170
ZHH2=HJ1*(B176+B172/2.)/B170
ZHH3=HJ1*(B177+B173/2.)/B170
ZHH4=HJ1*(B178+B174/2.)/B170
B200=-RCS*RCS/RH
B201=SQRT(-B200*(RH+B200))*RCC/RCS
B202=(RH+B200)**2+B201*B201
B203=SQRT(B202)
B204=B201/B203
B205=(RH+B200)/B203
B206=ASIN(B204)
B207=1.-B205
B208=1.+B205
GF6=B207/2.
B209=RH*B205+SQRT(RP*RP-RH*RH*B204*B204)
B210=B209-B203
B211=EXP(-EPT1*B203)*B204
B212=EXP(-EPT2*B203)*B204
B213=EXP(-EPT3*B203)*B204
B214=EXP(-EPT4*B203)*B204
B215=0.0
B216=0.0
B217=0.0
B218=0.0
B220=0.0
B221=0.0
B222=0.0
B223=0.0
B225=EXP(-EPT1*B210)*B204
B226=EXP(-EPT2*B210)*B204
B227=EXP(-EPT3*B210)*B204
B228=EXP(-EPT4*B210)*B204
B229=0.0
B230=0.0
B231=0.0
B232=0.0
B240=EXP(-EPT1*B209)*B204
B241=EXP(-EPT2*B209)*B204
B242=EXP(-EPT3*B209)*B204
B243=EXP(-EPT4*B209)*B204
B244=0.0
B245=0.0
B246=0.0
B247=0.0
HJ2=B206/FLOAT(NNK)
DO 348 K51=1,NNK1
B300=FLOAT(K51)*HJ2
B301=SIN(B300)
B302=COS(B300)
B303=B301/B302
B304=RH*B303

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B305=B303**2+(RCC/RCS)**2
B306=(B304+RCC)*(B304-RCC)
B307=(2.*B303*B304)**2-4.*B305*B306
B308=SQRT(B307)
B309=(-2.*B303*B304-B308)/(2.*B305)
B310=(RH+B309)/B302
B311=(-2.*B303*B304+B308)/(2.*B305)
B312=(RH+B311)/B302
B313=RH*B302+SQRT(RP*RP-RH*RH*B301*B301)
B315=EXP(-EPT1*B310)*B301
B316=EXP(-EPT2*B310)*B301
B317=EXP(-EPT3*B310)*B301
B318=EXP(-EPT4*B310)*B301
B215=B215+B315
B216=B216+B316
B217=B217+B317
B218=B218+B318
B320=B312-B310
B321=EXP(-ET1*B320)*B301
B322=EXP(-ET2*B320)*B301
B323=EXP(-ET3*B320)*B301
B324=EXP(-ET4*B320)*B301
B220=B220+B321
B221=B221+B322
B222=B222+B323
B223=B223+B324
B330=B313-B312
B331=EXP(-EPT1*B330)*B301
B332=EXP(-EPT2*B330)*B301
B333=EXP(-EPT3*B330)*B301
B334=EXP(-EPT4*B330)*B301
B229=B229+B331
B230=B230+B332
B231=B231+B333
B232=B232+B334
348 CONTINUE
G1H1=HJ2*(B215+B211/2.)/B207
G1H2=HJ2*(B216+B212/2.)/B207
G1H3=HJ2*(B217+B213/2.)/B207
G1H4=HJ2*(B218+B214/2.)/B207
G1F1=HJ2*(B220+B204/2.)/B207
G1F2=HJ2*(B221+B204/2.)/B207
G1F3=HJ2*(B222+B204/2.)/B207
G1F4=HJ2*(B223+B204/2.)/B207
G2H1=HJ2*(B229+B225/2.)/B207
G2H2=HJ2*(B230+B226/2.)/B207
G2H3=HJ2*(B231+B227/2.)/B207
G2H4=HJ2*(B232+B228/2.)/B207
HJ3=(PI-B206)/FLOAT(NNK)
DO 350 K52=1,NNK1
B340=B206+FLOAT(K52)*HJ3
B341=SIN(B340)
B342=COS(B340)
B343=RH*B342+SQRT(RP*RP-RH*RH*B341*B341)
B344=EXP(-EPT1*B343)*B341
B345=EXP(-EPT2*B343)*B341
B346=EXP(-EPT3*B343)*B341
B347=EXP(-EPT4*B343)*B341
B244=B244+B344
B245=B245+B345
B246=B246+B346
B247=B247+B347
350 CONTINUE
GHH1=HJ3*(B244+B240/2.)/B208
GHH2=HJ3*(B245+B241/2.)/B208
GHH3=HJ3*(B246+B242/2.)/B208
GHH4=HJ3*(B247+B243/2.)/B208
GCC=(GF1+GF2+GF5+GF6)*0.2+(GF3+GF4)*0.1
GFP=1.-GCC
HYP11=(T1H1+E1H1+Z1H1+G1H1)*0.2+(V1H1+F1H1)*0.1
HYP12=(T1H2+E1H2+Z1H2+G1H2)*0.2+(V1H2+F1H2)*0.1
HYP13=(T1H3+E1H3+Z1H3+G1H3)*0.2+(V1H3+F1H3)*0.1
HYP14=(T1H4+E1H4+Z1H4+G1H4)*0.2+(V1H4+F1H4)*0.1
HYCC1=(T1F1+E1F1+Z1F1+G1F1)*0.2+(V1F1+F1F1)*0.1
HYCC2=(T1F2+E1F2+Z1F2+G1F2)*0.2+(V1F2+F1F2)*0.1
HYCC3=(T1F3+E1F3+Z1F3+G1F3)*0.2+(V1F3+F1F3)*0.1
HYCC4=(T1F4+E1F4+Z1F4+G1F4)*0.2+(V1F4+F1F4)*0.1
HYP21=(T2H1+E2H1+Z2H1+G2H1)*0.2+(V2H1+F2H1)*0.1
HYP22=(T2H2+E2H2+Z2H2+G2H2)*0.2+(V2H2+F2H2)*0.1

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HYP23=(T2H3+E2H3+Z2H3+G2H3)*0.2+(V2H3+F2H3)*0.1
HYP24=(T2H4+E2H4+Z2H4+G2H4)*0.2+(V2H4+F2H4)*0.1
HXPP1=(THH1+EHH1+ZHH1+GHH1)*0.2+(VHH1+FHH1)*0.1
HXPP2=(THH2+EHH2+ZHH2+GHH2)*0.2+(VHH2+FHH2)*0.1
HXPP3=(THH3+EHH3+ZHH3+GHH3)*0.2+(VHH3+FHH3)*0.1
HXPP4=(THH4+EHH4+ZHH4+GHH4)*0.2+(VHH4+FHH4)*0.1
GC1=GCC
GC2=GCC
GC3=GCC
GC4=GCC
GP1=GPP
GP2=GPP
GP3=GPP
GP4=GPP
YP11=HYP11
YP12=HYP12
YP13=HYP13
YP14=HYP14
YCC1=HYCC1
YCC2=HYCC2
YCC3=HYCC3
YCC4=HYCC4
YP21=HYP21
YP22=HYP22
YP23=HYP23
YP24=HYP24
XPP1=HXPP1
XPP2=HXPP2
XPP3=HXPP3
XPP4=HXPP4
500 CONTINUE
C
C *****
C NEUTRON ABSORPTION BY THE FUEL AND HYDROGEN REGIONS
C *****
C
H100=SAC*EPP4*(1.-EF4)/H11
H101=H100
H103=(SAP*(1.-OPP4)+SAC*(1.-EPP4)+SAC*EPP4*EF4*(1.-EPP4))/
#H11
H104=H103*EPA4/EPT4
H105=H103*EPS44/EPT4
H106=1.-TC4
H107=H106
H109=TC4*(1.-TP4)
H110=H109*EPA4/EPT4
H111=H109*EPS44/EPT4
H112=TC4*TP4*RA44
H113=H107+H112*H101
H115=H110+H112*H104
H116=H111+H112*H105
H117=(GP4*XPP4+GC4*YP14*YCC4*YP24)*RA44
H118=GC4*YP14*(1.-YCC4)
H119=H118
H121=GP4*(1.-XPP4)+GC4*(1.-YP14)+GC4*YP14*YCC4*(1.-YP24)
H122=H121*EPA4/EPT4
H123=H121*EPS44/EPT4
H124=H119+H117*H101
H126=H122+H117*H104
H127=H123+H117*H105
H128=H124/(1.-H127)
H130=H126/(1.-H127)
H131=H113+H116*H128
H133=H115+H116*H130
Q11=H131
Q10=H133
Q9=H128
Q8=H130
Q12=H104+H105*Q8
Q13=H101+H105*Q9
Q20=(SAC*EPP3*EPP3*EF3+SAP*OPP3)*RA34/H12
QQ20=SAC*EPP3*(1.-EF3)/H12
Q21=QQ20*EA3/ET3
Q22=QQ20*ES33/ET3
Q23=QQ20*ES34/ET3
Q24=(SAP*(1.-OPP3)+SAC*(1.-EPP3)+SAC*EPP3*EF3*(1.-EPP3))/
#H12
Q25=Q24*EPA3/EPT3
Q26=Q24*EPS33/EPT3

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Q27=Q24*EPS34/EPT3
 Q28=(GP3*XP3+GC3*YP13*YCC3*YP23)*RA34
 Q29=(GP3*XP3+GC3*YP13*YCC3*YP23)*RA33
 QQ29=GC3*YP13*(1.-YCC3)
 Q30=QQ29*EA3/ET3
 Q31=QQ29*ES33/ET3
 Q32=QQ29*ES34/ET3
 Q33=GP3*(1.-XP3)+GC3*(1.-YP13)+GC3*YP13*YCC3*(1.-YP23)
 Q34=Q33*EPA3/EPT3
 Q35=Q33*EPS33/EPT3
 Q36=Q33*EPS34/EPT3
 Q37=Q20*Q13+Q23*Q11+Q27*Q9
 Q38=Q20*Q12+Q23*Q10+Q27*Q8
 Q39=Q28*Q13+Q32*Q11+Q36*Q9+Q29*Q37
 Q40=Q28*Q12+Q32*Q10+Q36*Q8+Q29*Q38
 Q42=Q30+Q29*Q21
 Q43=Q34+Q29*Q25
 Q44=Q35+Q29*Q26
 Q45=Q31+Q29*Q22
 Q46=TC3*TP3*RA34
 Q47=TC3*TP3*RA33
 QQ47=1.-TC3
 Q48=QQ47*EA3/ET3
 Q49=QQ47*ES33/ET3
 Q50=QQ47*ES34/ET3
 QQ50=TC3*(1.-TP3)
 Q51=QQ50*EPA3/EPT3
 Q52=QQ50*EPS33/EPT3
 Q53=QQ50*EPS34/EPT3
 Q54=Q46*Q13+Q50*Q11+Q53*Q9+Q47*Q37
 Q55=Q46*Q12+Q50*Q10+Q53*Q8+Q47*Q38
 Q56=Q48+Q47*Q21
 Q57=Q51+Q47*Q25
 Q58=Q49+Q47*Q22
 Q59=Q52+Q47*Q26
 Q390=1.-Q44
 Q60=Q54+Q59*Q39/Q390
 Q61=Q55+Q59*Q40/Q390
 Q62=Q56+Q59*Q42/Q390
 Q63=Q57+Q59*Q43/Q390
 Q64=Q58+Q59*Q45/Q390
 Q391=1.-Q64
 Q65=Q60/Q391
 Q66=Q61/Q391
 Q67=Q62/Q391
 Q68=Q63/Q391
 Q69=(Q39+Q45*Q65)/Q390
 Q70=(Q40+Q45*Q66)/Q390
 Q71=(Q42+Q45*Q67)/Q390
 Q72=(Q43+Q45*Q68)/Q390
 Q73=Q37+Q22*Q65+Q26*Q69
 Q74=Q38+Q22*Q66+Q26*Q70
 Q75=Q21+Q22*Q67+Q26*Q71
 Q76=Q25+Q22*Q68+Q26*Q72
 Q100=(SAC*EPP2*EPP2*EF2+SAP*OPP2)*RA24/H13
 Q101=(SAC*EPP2*EPP2*EF2+SAP*OPP2)*RA23/H13
 QQ101=SAC*EPP2*(1.-EF2)/H13
 Q102=QQ101*EA2/ET2
 Q103=QQ101*ES22/ET2
 Q104=QQ101*ES23/ET2
 Q105=QQ101*ES24/ET2
 Q106=(SAP*(1.-OPP2)+SAC*(1.-EPP2)+SAC*EPP2*EF2*(1.-EPP2))/H13
 Q107=Q106*EPA2/EPT2
 Q108=Q106*EPS22/EPT2
 Q109=Q106*EPS23/EPT2
 Q110=Q106*EPS24/EPT2
 Q111=Q100*Q13+Q105*Q11+Q110*Q9+Q101*Q73+Q104*Q65+Q109*Q69
 Q112=Q100*Q12+Q105*Q10+Q110*Q8+Q101*Q74+Q104*Q66+Q109*Q70
 Q113=Q101*Q75+Q104*Q67+Q109*Q71
 Q114=Q101*Q76+Q104*Q68+Q109*Q72
 Q115=TC2*TP2*RA24
 Q116=TC2*TP2*RA23
 Q117=TC2*TP2*RA22
 QQ117=1.-TC2
 Q118=QQ117*EA2/ET2
 Q119=QQ117*ES22/ET2
 Q120=QQ117*ES23/ET2
 Q121=QQ117*ES24/ET2

QQ122=TC2*(1.-TP2)
 Q122=QQ122*EPA2/EPT2
 Q123=QQ122*EPS22/EPT2
 Q124=QQ122*EPS23/EPT2
 Q125=QQ122*EPS24/EPT2
 Q126=(GP2*XPP2+GC2*YP12*YCC2*YP22)*RA24
 Q127=(GP2*XPP2+GC2*YP12*YCC2*YP22)*RA23
 Q128=(GP2*XPP2+GC2*YP12*YCC2*YP22)*RA22
 QQ129=GC2*YP12*(1.-YCC2)
 Q129=QQ129*EA2/ET2
 Q130=QQ129*ES22/ET2
 Q131=QQ129*ES23/ET2
 Q132=QQ129*ES24/ET2
 Q133=GP2*(1.-XPP2)+GC2*(1.-YP12)+GC2*YP12*YCC2*(1.-YP22)
 Q134=Q133*EPA2/EPT2
 Q135=Q133*EPS22/EPT2
 Q136=Q133*EPS23/EPT2
 Q137=Q133*EPS24/EPT2
 Q138=Q126*Q13+Q132*Q11+Q137*Q9+Q127*Q73+Q131*Q65+Q136*Q69+
 #Q128*Q111
 Q139=Q126*Q12+Q132*Q10+Q137*Q8+Q127*Q74+Q131*Q66+Q136*Q70+
 #Q128*Q112
 Q140=Q127*Q75+Q131*Q67+Q136*Q71+Q128*Q113
 Q141=Q127*Q76+Q131*Q68+Q136*Q72+Q128*Q114
 Q142=Q129+Q128*Q102
 Q143=Q130+Q128*Q103
 Q144=Q134+Q128*Q107
 Q145=Q135+Q128*Q108
 Q146=Q115*Q13+Q121*Q11+Q125*Q9+Q116*Q73+Q120*Q65+Q124*Q69+
 #Q117*Q111
 Q147=Q115*Q12+Q121*Q10+Q125*Q8+Q116*Q74+Q120*Q66+Q124*Q70+
 #Q117*Q112
 Q148=Q116*Q75+Q120*Q67+Q124*Q71+Q117*Q113
 Q149=Q116*Q76+Q120*Q68+Q124*Q72+Q117*Q114
 Q392=(Q123+Q117*Q108)/(1.-Q145)
 Q150=Q146+Q138*Q392
 Q151=Q147+Q139*Q392
 Q152=Q148+Q140*Q392
 Q153=Q149+Q141*Q392
 Q154=Q118+Q117*Q102+Q142*Q392
 Q155=Q122+Q117*Q107+Q144*Q392
 Q156=Q119+Q117*Q103+Q143*Q392
 Q393=1.-Q156
 Q157=Q150/Q393
 Q158=Q151/Q393
 Q159=Q152/Q393
 Q160=Q153/Q393
 Q161=Q154/Q393
 Q162=Q155/Q393
 Q163=Q143/(1.-Q145)
 Q394=1.-Q145
 Q164=Q138/Q394+Q163*Q157
 Q165=Q139/Q394+Q163*Q158
 Q166=Q140/Q394+Q163*Q159
 Q167=Q141/Q394+Q163*Q160
 Q168=Q142/Q394+Q163*Q161
 Q169=Q144/Q394+Q163*Q162
 Q170=Q111+Q103*Q157+Q108*Q164
 Q171=Q112+Q103*Q158+Q108*Q165
 Q172=Q113+Q103*Q159+Q108*Q166
 Q173=Q114+Q103*Q160+Q108*Q167
 Q174=Q102+Q103*Q161+Q108*Q168
 Q175=Q107+Q103*Q162+Q108*Q169
 Q190=SAC*EPP1*EPP1*EF1+SAP*OPP1
 Q200=Q190*RA14/H14
 Q201=Q190*RA13/H14
 Q202=Q190*RA12/H14
 QQ203=SAC*EPP1*(1.-EF1)/H14
 Q203=QQ203*EA1/ET1
 Q204=QQ203*ES11/ET1
 Q205=QQ203*ES12/ET1
 Q206=QQ203*ES13/ET1
 Q207=QQ203*ES14/ET1
 Q208=(SAP*(1.-OPP1)+SAC*(1.-EPP1))+SAC*EPP1*EF1*(1.-EPP1)/
 #H14
 Q209=Q208*EPA1/EPT1
 Q210=Q208*EPS11/EPT1
 Q211=Q208*EPS12/EPT1
 Q212=Q208*EPS13/EPT1

Q213=Q208*EFS14/EPT1
 Q214=Q200*Q13+Q207*Q11+Q213*Q9+Q201*Q73+Q206*Q65+Q212*Q69+
 #Q202*Q170+Q205*Q157+Q211*Q164
 Q215=Q200*Q12+Q207*Q10+Q213*Q8+Q201*Q74+Q206*Q66+Q212*Q70+
 #Q202*Q171+Q205*Q158+Q211*Q165
 Q216=Q201*Q75+Q206*Q67+Q212*Q71+Q202*Q172+Q205*Q159+Q211*
 #Q166
 Q217=Q201*Q76+Q206*Q68+Q212*Q72+Q202*Q173+Q205*Q160+Q211*
 #Q167
 Q218=Q202*Q174+Q205*Q161+Q211*Q168
 Q219=Q202*Q175+Q205*Q162+Q211*Q169
 Q220=TC1*TP1*RA14
 Q221=TC1*TP1*RA13
 Q222=TC1*TP1*RA12
 Q223=TC1*TP1*RA11
 Q224=1.-TC1
 Q224=Q224*EA1/ET1
 Q225=Q224*ES11/ET1
 Q226=Q224*ES12/ET1
 Q227=Q224*ES13/ET1
 Q228=Q224*ES14/ET1
 Q229=TC1*(1.-TP1)
 Q229=Q229*EPA1/EPT1
 Q230=Q229*EFS11/EPT1
 Q231=Q229*EFS12/EPT1
 Q232=Q229*EFS13/EPT1
 Q233=Q229*EFS14/EPT1
 Q234=Q220*Q13+Q228*Q11+Q233*Q9+Q221*Q73+Q227*Q65+Q232*Q69+
 #Q222*Q170+Q226*Q157+Q231*Q164+Q223*Q214
 Q235=Q220*Q12+Q228*Q10+Q233*Q8+Q221*Q74+Q227*Q66+Q232*Q70+
 #Q222*Q171+Q226*Q158+Q231*Q165+Q223*Q215
 Q236=Q221*Q75+Q227*Q67+Q232*Q71+Q222*Q172+Q226*Q159+Q231*
 #Q166+Q223*Q216
 Q237=Q221*Q76+Q227*Q68+Q232*Q72+Q222*Q173+Q226*Q160+Q231*
 #Q167+Q223*Q217
 Q238=Q222*Q174+Q226*Q161+Q231*Q168+Q223*Q218
 Q239=Q222*Q175+Q226*Q162+Q231*Q169+Q223*Q219
 Q240=Q223*Q203+Q224
 Q241=Q223*Q204+Q225
 Q242=Q223*Q209+Q229
 Q243=Q223*Q210+Q230
 Q191=GP1*XP11+GC1*YP11*YCC1*YP21
 Q244=Q191*RA14
 Q245=Q191*RA13
 Q246=Q191*RA12
 Q247=Q191*RA11
 Q248=GC1*YP11*(1.-YCC1)
 Q248=Q248*EA1/ET1
 Q249=Q248*ES11/ET1
 Q250=Q248*ES12/ET1
 Q251=Q248*ES13/ET1
 Q252=Q248*ES14/ET1
 Q253=GP1*(1.-XP11)+GC1*(1.-YP11)+GC1*YP11*YCC1*(1.-YP21)
 Q254=Q253*EPA1/EPT1
 Q255=Q253*EFS11/EPT1
 Q256=Q253*EFS12/EPT1
 Q257=Q253*EFS13/EPT1
 Q258=Q253*EFS14/EPT1
 Q259=Q244*Q13+Q252*Q11+Q258*Q9+Q245*Q73+Q251*Q65+Q257*Q69+
 #Q246*Q170+Q250*Q157+Q256*Q164+Q247*Q214
 Q260=Q244*Q12+Q252*Q10+Q258*Q8+Q245*Q74+Q251*Q66+Q257*Q70+
 #Q246*Q171+Q250*Q158+Q256*Q165+Q247*Q215
 Q261=Q245*Q75+Q251*Q67+Q257*Q71+Q246*Q172+Q250*Q159+Q256*
 #Q166+Q247*Q216
 Q262=Q245*Q76+Q251*Q68+Q257*Q72+Q246*Q173+Q250*Q160+Q256*
 #Q167+Q247*Q217
 Q263=Q246*Q174+Q250*Q161+Q256*Q168+Q247*Q218
 Q264=Q246*Q175+Q250*Q162+Q256*Q169+Q247*Q219
 Q265=Q247*Q203+Q248
 Q266=Q247*Q204+Q249
 Q267=Q247*Q209+Q254
 Q268=Q247*Q210+Q255
 Q400=1.-Q268
 Q269=Q259/Q400
 Q270=Q260/Q400
 Q271=Q261/Q400
 Q272=Q262/Q400
 Q273=Q263/Q400
 Q274=Q264/Q400

Q275=Q265/Q400
 Q276=Q266/Q400
 Q277=Q267/Q400
 Q278=Q234+Q243*Q269
 Q279=Q235+Q243*Q270
 Q280=Q236+Q243*Q271
 Q281=Q237+Q243*Q272
 Q282=Q238+Q243*Q273
 Q283=Q239+Q243*Q274
 Q284=Q240+Q243*Q275
 Q285=Q241+Q243*Q276
 Q286=Q242+Q243*Q277
 Q401=1.-Q285
 Q287=Q278/Q401
 Q288=Q279/Q401
 Q289=Q280/Q401
 Q290=Q281/Q401
 Q291=Q282/Q401
 Q292=Q283/Q401
 Q293=Q284/Q401
 Q294=Q286/Q401
 Q295=Q269+Q276*Q287
 Q296=Q270+Q276*Q288
 Q297=Q271+Q276*Q289
 Q298=Q272+Q276*Q290
 Q299=Q273+Q276*Q291
 Q300=Q274+Q276*Q292
 Q301=Q275+Q276*Q293
 Q302=Q277+Q276*Q294
 Q303=Q214+Q204*Q287+Q210*Q295
 Q304=Q215+Q204*Q288+Q210*Q296
 Q305=Q216+Q204*Q289+Q210*Q297
 Q306=Q217+Q204*Q290+Q210*Q298
 Q307=Q218+Q204*Q291+Q210*Q299
 Q308=Q219+Q204*Q292+Q210*Q300
 Q309=Q203+Q204*Q293+Q210*Q301
 Q310=Q209+Q204*Q294+Q210*Q302
 Q500=(1.-TC4)*EA4/ET4
 Q501=(1.-TC4)*ES44/ET4
 Q502=Q500/(1.-Q501)
 Q503=(1.-TC3)*EA3/ET3
 Q504=(1.-TC3)*ES33/ET3
 Q505=(1.-TC3)*ES34/ET3
 Q506=Q503/(1.-Q504)
 Q507=Q505/(1.-Q504)
 Q508=Q507*Q502
 Q509=(1.-TC2)*EA2/ET2
 Q510=(1.-TC2)*ES22/ET2
 Q511=(1.-TC2)*ES23/ET2
 Q512=(1.-TC2)*ES24/ET2
 Q513=Q509/(1.-Q510)
 Q514=Q511/(1.-Q510)
 Q515=Q512/(1.-Q510)
 Q516=Q515*Q502+Q514*Q508
 Q517=Q514*Q506
 Q518=(1.-TC1)*EA1/ET1
 Q519=(1.-TC1)*ES11/ET1
 Q520=(1.-TC1)*ES12/ET1
 Q521=(1.-TC1)*ES13/ET1
 Q522=(1.-TC1)*ES14/ET1
 Q523=Q518/(1.-Q519)
 Q524=Q520/(1.-Q519)
 Q525=Q521/(1.-Q519)
 Q526=Q522/(1.-Q519)
 Q527=Q524*Q513
 Q528=Q524*Q517+Q525*Q506
 Q529=Q524*Q516+Q525*Q508+Q526*Q502
 Q530=CHI1*Q523
 Q531=CHI1*Q527+CHI2*Q513
 Q532=CHI1*Q528+CHI2*Q517
 Q533=CHI1*Q529+CHI2*Q516
 Q320=CHI1*Q293
 Q321=CHI1*Q291+CHI2*Q161
 Q322=CHI1*Q289+CHI2*Q159
 Q323=CHI1*Q287+CHI2*Q157
 Q534=Q320-Q530
 Q535=Q321-Q531
 Q536=Q322-Q532
 Q537=Q323-Q533

Q324=Q320+Q321+Q322+Q323
 Q538=Q530+Q531+Q532+Q533
 Q539=Q324+Q538
 Q330=CHI1*Q294
 Q331=CHI1*Q292+CHI2*Q162
 Q332=CHI1*Q290+CHI2*Q160
 Q333=CHI1*Q288+CHI2*Q158
 Q334=Q330+Q331+Q332+Q333

 EFFECTIVE NEUTRON MULTIPLICATION FACTOR CALCULATION

EKEF=Q320*(FV1+2.*EN1)/EA1+Q321*FV2/EA2+Q322*FV3/EA3+
 #Q323*FV4/EA4

 OUTPUT RESULTS

REFLECTOR ALBEDOS (SEE TABLE 7-3)

RA11 = PROBABILITY THAT A NEUTRON FROM GROUP 1 BE REFLECTED
 WITH GROUP 1 ENERGY
 RA12 = PROBABILITY THAT A NEUTRON FROM GROUP 1 BE REFLECTED
 WITH GROUP 2 ENERGY , AND SO ON.

FUEL AND HYDROGEN TRANSMITTANCES (SEE TABLE 7-4)

TC1 = ZETA 1 (SEE EQ. 7-1)
 TC2 = ZETA 2 , AND SO ON (SEE EQ. 7-1).
 TP1 = EPSILON 1 (SEE EQ. 7-2)
 TP2 = EPSILON 2 , AND SO ON (SEE EQ. 7-2).
 YP11 = SA-HYDROGEN-PATH XI FOR GROUP 1 (SEE EQ. 7-3)
 YP12 = SA-HYDROGEN-PATH XI FOR GROUP 2 (SEE EQ. 7-3)
 YP13 = SA-HYDROGEN-PATH XI FOR GROUP 3 (SEE EQ. 7-3)
 YP14 = SA-HYDROGEN-PATH XI FOR GROUP 4 (SEE EQ. 7-3)
 YCC1 = AB-FUEL-PATH XI FOR GROUP 1 (SEE EQ. 7-4)
 YCC2 = AB-FUEL-PATH XI FOR GROUP 2 (SEE EQ. 7-4)
 YCC3 = AB-FUEL-PATH XI FOR GROUP 3 (SEE EQ. 7-4)
 YCC4 = AB-FUEL-PATH XI FOR GROUP 4 (SEE EQ. 7-4)
 YP21 = BC-HYDROGEN-PATH XI FOR GROUP 1 (SEE EQ. 7-5)
 YP22 = BC-HYDROGEN-PATH XI FOR GROUP 2 (SEE EQ. 7-5)
 YP23 = BC-HYDROGEN-PATH XI FOR GROUP 3 (SEE EQ. 7-5)
 YP24 = BC-HYDROGEN-PATH XI FOR GROUP 4 (SEE EQ. 7-5)
 XPP1 = SQ-HYDROGEN-PATH XI FOR GROUP 1 (SEE EQ. 7-6)
 XPP2 = SQ-HYDROGEN-PATH XI FOR GROUP 2 (SEE EQ. 7-6)
 XPP3 = SQ-HYDROGEN-PATH XI FOR GROUP 3 (SEE EQ. 7-6)
 XPP4 = SQ-HYDROGEN-PATH XI FOR GROUP 4 (SEE EQ. 7-6)
 EPP1 = SF-HYDROGEN-PATH ETA FOR GROUP 1 (SEE EQ. 7-7)
 EPP2 = SF-HYDROGEN-PATH ETA FOR GROUP 2 , AND SO ON
 (SEE EQ. 7-7).
 EF1 = FG-FUEL-PATH ETA FOR GROUP 1 (SEE EQ. 7-8)
 EF2 = FG-FUEL-PATH ETA FOR GROUP 2 , AND SO ON
 (SEE EQ. 7-8).
 OPP1 = SL-HYDROGEN-PATH FOR GROUP 1 (SEE EQ. 7-10)
 OPP2 = SL-HYDROGEN-PATH FOR GROUP 2 , AND SO ON
 (SEE EQ. 7-10).
 NOTE : SF-HYDROGEN-PATH TRANSMITTANCES EQUAL TO
 GH-HYDROGEN-PATH TRANSMITTANCES (SEE EQ. 7-9 AND
 EQ. 7-7).

NEUTRON ABSORPTION BY THE FUEL REGION

Q530 = FRACTION OF FISSION NEUTRONS (NEVER TRAVERSE THE
 FUEL BOUNDARY) WHICH ARE ABSORBED AS GROUP 1
 Q531 = FRACTION OF FISSION NEUTRONS (NEVER TRAVERSE THE
 FUEL BOUNDARY) WHICH ARE ABSORBED AS GROUP 2
 Q532 = FRACTION OF FISSION NEUTRONS (NEVER TRAVERSE THE
 FUEL BOUNDARY) WHICH ARE ABSORBED AS GROUP 3
 Q533 = FRACTION OF FISSION NEUTRONS (NEVER TRAVERSE THE
 FUEL BOUNDARY) WHICH ARE ABSORBED AS GROUP 4
 Q534 = FRACTION OF FISSION NEUTRONS (TRAVERSE THE FUEL
 BOUNDARY) WHICH ARE ABSORBED AS GROUP 1
 Q535 = FRACTION OF FISSION NEUTRONS (TRAVERSE THE FUEL
 BOUNDARY) WHICH ARE ABSORBED AS GROUP 2
 Q536 = FRACTION OF FISSION NEUTRONS (TRAVERSE THE FUEL
 BOUNDARY) WHICH ARE ABSORBED AS GROUP 3
 Q537 = FRACTION OF FISSION NEUTRONS (TRAVERSE THE FUEL

```

C          BOUNDARY ) WHICH ARE ABSORBED AS GROUP 4
C
C          NEUTRON ABSORPTION BY THE HYDROGEN REGION
C
C          Q330 = FRACTION OF FISSION NEUTRONS WHICH ARE ABSORBED
C              AS GROUP 1
C          Q331 = FRACTION OF FISSION NEUTRONS WHICH ARE ABSORBED
C              AS GROUP 2
C          Q332 = FRACTION OF FISSION NEUTRONS WHICH ARE ABSORBED
C              AS GROUP 3
C          Q333 = FRACTION OF FISSION NEUTRONS WHICH ARE ABSORBED
C              AS GROUP 4
C
C          EFFECTIVE NEUTRON MULTIPLICATION FACTOR
C
C          EKEF = EFFECTIVE NEUTRON MULTIPLICATION FACTOR
C          *****
C
C          WRITE(02,100)RP,RC
C          WRITE(02,101)PDF,RDEL,AF
C          WRITE(02,102)RA11,RA12,RA13,RA14
C          WRITE(02,103)RA22,RA23,RA24
C          WRITE(02,104)RA33,RA34,RA44
C          WRITE(02,105)TC1,TC2,TC3,TC4
C          WRITE(02,106)TP1,TP2,TP3,TP4
C          WRITE(02,107)YP11,YP12,YP13,YP14
C          WRITE(02,108)YCC1,YCC2,YCC3,YCC4
C          WRITE(02,109)YP21,YP22,YP23,YP24
C          WRITE(02,110)XPP1,XPP2,XPP3,XPP4
C          WRITE(02,111)EPP1,EPP2,EPP3,EPP4
C          WRITE(02,112)EP1,EP2,EP3,EP4
C          WRITE(02,113)OPP1,OPP2,OPP3,OPP4
C          WRITE(02,114)Q530,Q531,Q532,Q533
C          WRITE(02,115)Q534,Q535,Q536,Q537
C          WRITE(02,116)Q330,Q331,Q332,Q333
C          WRITE(02,117)EKEF
100  FORMAT(3X,'RP  =',E12.5,1X,'RC  =',E12.5)
101  FORMAT(3X,'PDF  =',E12.5,1X,'RDEL=',E12.5,1X,'AF  =',E12.5)
102  FORMAT(3X,'RA11=',E12.5,1X,'RA12=',E12.5,1X,'RA13=',E12.5,
      #1X,'RA14=',E12.5)
103  FORMAT(3X,'RA22=',E12.5,1X,'RA23=',E12.5,1X,'RA24=',E12.5)
104  FORMAT(3X,'RA33=',E12.5,1X,'RA34=',E12.5,1X,'RA44=',E12.5)
105  FORMAT(3X,'TC1  =',E12.5,1X,'TC2  =',E12.5,1X,'TC3  =',E12.5,
      #1X,'TC4  =',E12.5)
106  FORMAT(3X,'TP1  =',E12.5,1X,'TP2  =',E12.5,1X,'TP3  =',E12.5,
      #1X,'TP4  =',E12.5)
107  FORMAT(3X,'YP11=',E12.5,1X,'YP12=',E12.5,1X,'YP13=',E12.5,
      #1X,'YP14=',E12.5)
108  FORMAT(3X,'YCC1=',E12.5,1X,'YCC2=',E12.5,1X,'YCC3=',E12.5,
      #1X,'YCC4=',E12.5)
109  FORMAT(3X,'YP21=',E12.5,1X,'YP22=',E12.5,1X,'YP23=',E12.5,
      #1X,'YP24=',E12.5)
110  FORMAT(3X,'XPP1=',E12.5,1X,'XPP2=',E12.5,1X,'XPP3=',E12.5,
      #1X,'XPP4=',E12.5)
111  FORMAT(3X,'EPP1=',E12.5,1X,'EPP2=',E12.5,1X,'EPP3=',E12.5,
      #1X,'EPP4=',E12.5)
112  FORMAT(3X,'EP1  =',E12.5,1X,'EP2  =',E12.5,1X,'EP3  =',E12.5,
      #1X,'EP4  =',E12.5)
113  FORMAT(3X,'OPP1=',E12.5,1X,'OPP2=',E12.5,1X,'OPP3=',E12.5,
      #1X,'OPP4=',E12.5)
114  FORMAT(3X,'Q530=',E12.5,1X,'Q531=',E12.5,1X,'Q532=',E12.5,
      #1X,'Q533=',E12.5)
115  FORMAT(3X,'Q534=',E12.5,1X,'Q535=',E12.5,1X,'Q536=',E12.5,
      #1X,'Q537=',E12.5)
116  FORMAT(3X,'Q330=',E12.5,1X,'Q331=',E12.5,1X,'Q332=',E12.5,
      #1X,'Q333=',E12.5)
117  FORMAT(3X,'EFFECTIVE MULTIPLICATION FACTOR',9X,'=',E12.5)
999  STOP
      END

```


Source Listing of ALB5

```

C *****
C REFLECTOR ALBEDOS FOR FINITE REFLECTOR
C ( P1 - SPHERICAL GEOMETRY / FOUR - GROUP ENERGY )
C *****
C
C PROGRAM ALB5
C NTIN=01
C NTOUT=02
C OPEN(NTIN, FILE='ALB5.IN', STATUS='OLD')
C OPEN(NTOUT, FILE='ALB5.OUT', STATUS='UNKNOWN')
C
C *****
C INPUT DATA
C
C R = INNER RADIUS OF THE REFLECTOR
C T = THICKNESS OF THE REFLECTOR
C
C REFLECTOR GROUP PARAMETERS ( SEE TABLE 6-3 )
C
C DR1 = DIFFUSION COEFFICIENT FOR GROUP 1
C DR2 = DIFFUSION COEFFICIENT FOR GROUP 2 , AND SO ON.
C ERN1 = MACROSCOPIC (n,2n) CROSS SECTION FOR GROUP 1
C ERA1 = MACROSCOPIC ABSORPTION CROSS SECTION FOR GROUP 1
C ERA2 = MACROSCOPIC ABSORPTION CROSS SECTION FOR GROUP 2 ,
C AND SO ON.
C ERS12= MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C GROUP 1 TO GROUP 2
C ERS13= MACROSCOPIC GROUP-TRANSFER CROSS SECTION FROM
C GROUP 1 TO GROUP 3 , AND SO ON.
C *****
C
C READ(01,10)R,T
C READ(01,11)DR1,DR2,DR3,DR4,ERN1,ERA1,ERA2,ERA3,ERA4,
C ERS12,ERS13,ERS14,ERS23,ERS24,ERS34
10 FORMAT(2E17.7)
11 FORMAT(4E17.7)
C ERR1=ERA1+ERS12+ERS13+ERS14-ERN1
C ERR2=ERA2+ERS23+ERS24
C ERR3=ERA3+ERS34
C ERR4=ERA4
C RK1=SQRT(ERR1/DR1)
C RK2=SQRT(ERR2/DR2)
C RK3=SQRT(ERR3/DR3)
C RK4=SQRT(ERR4/DR4)
C
C *****
C REFLECTOR ALBEDO CALCULATION
C *****
C
C H=R+T
C X1=EXP(-RK1*R)
C X2=EXP(-RK2*R)
C X3=EXP(-RK3*R)
C X4=EXP(-RK4*R)
C X5=EXP(+RK1*R)
C X6=EXP(+RK2*R)
C X7=EXP(+RK3*R)
C X8=EXP(+RK4*R)
C X9=EXP(-RK1*H)
C X10=EXP(-RK2*H)
C X11=EXP(-RK3*H)
C X12=EXP(-RK4*H)
C X13=EXP(+RK1*H)
C X14=EXP(+RK2*H)
C X15=EXP(+RK3*H)
C X16=EXP(+RK4*H)
C X17=DR1*(RK1+1./H)/2.
C X18=DR1*(RK1-1./H)/2.
C X19=DR1*(RK1+1./R)/2.
C X20=DR1*(RK1-1./R)/2.
C X21=DR2*(RK1+RK1-RK2+RK2)
C X22=2.*DR2*(RK1+1./R)
C X23=2.*DR2*(RK1-1./R)
C X24=2.*DR2*(RK2+1./R)

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X25=2.*DR2*(RK2-1./R)
X26=2.*DR2*(RK1+1./H)
X27=2.*DR2*(RK1-1./H)
X28=2.*DR2*(RK2+1./H)
X29=2.*DR2*(RK2-1./H)
X30=DR3*(RK1+RK1-RK3+RK3)
X31=DR3*(RK2+RK2-RK3+RK3)
X32=2.*DR3*(RK1+1./R)
X33=2.*DR3*(RK1-1./R)
X34=2.*DR3*(RK2+1./R)
X35=2.*DR3*(RK2-1./R)
X36=2.*DR3*(RK3+1./R)
X37=2.*DR3*(RK3-1./R)
X38=2.*DR3*(RK1+1./H)
X39=2.*DR3*(RK1-1./H)
X40=2.*DR3*(RK2+1./H)
X41=2.*DR3*(RK2-1./H)
X42=2.*DR3*(RK3+1./H)
X43=2.*DR3*(RK3-1./H)
X44=DR4*(RK1+RK1-RK4+RK4)
X45=DR4*(RK2+RK2-RK4+RK4)
X46=DR4*(RK3+RK3-RK4+RK4)
X47=2.*DR4*(RK1+1./R)
X48=2.*DR4*(RK1-1./R)
X49=2.*DR4*(RK2+1./R)
X50=2.*DR4*(RK2-1./R)
X51=2.*DR4*(RK3+1./R)
X52=2.*DR4*(RK3-1./R)
X55=2.*DR4*(RK4+1./R)
X56=2.*DR4*(RK4-1./R)
X57=2.*DR4*(RK1+1./H)
X58=2.*DR4*(RK1-1./H)
X59=2.*DR4*(RK2+1./H)
X60=2.*DR4*(RK2-1./H)
X61=2.*DR4*(RK3+1./H)
X62=2.*DR4*(RK3-1./H)
X63=2.*DR4*(RK4+1./H)
X64=2.*DR4*(RK4-1./H)
Y11=X9*(.25-X17)/H
Y12=X13*(.25+X18)/H
Y13=X1*(.25+X19)/R
Y14=X5*(.25-X20)/R
YD1=Y11*Y14-Y12*Y13
Y15=Y12/YD1
Y16=Y11/YD1
RA11=Y15*X1*(.25-X19)/R+Y16*X5*(.25+X20)/R
Y17=ERS12*Y15/X21
Y18=ERS12*Y16/X21
Y19=Y17*X1*(1.+X22)-Y18*X5*(1.-X23)
Y20=X2*(1.+X24)
Y21=X6*(1.-X25)
Y22=Y17*X9*(1.-X26)-Y18*X13*(1.+X27)
Y23=X10*(1.-X28)
Y24=X14*(1.+X29)
YD2=Y20*Y24-Y21*Y23
Y25=(Y19*Y24-Y22*Y21)/YD2
Y26=(Y20*Y22-Y19*Y23)/YD2
Y27=(Y25*X2+Y26*X6+Y17*X1+Y18*X5)/R
RA12=Y27/2.
Y28=(ERS13*Y15+ERS23*Y17)/X30
Y29=(ERS13*Y16+ERS23*Y18)/X30
Y30=ERS23*Y25/X31
Y31=ERS23*Y26/X31
Y32=Y28*X1*(1.+X32)-Y29*X5*(1.-X33)-Y30*X2*(1.+X34)-Y31*
#X6*(1.-X35)
Y33=X3*(1.+X36)
Y34=X7*(1.-X37)
Y35=Y28*X9*(1.-X38)-Y29*X13*(1.+X39)-Y30*X10*(1.-X40)-
#Y31*X14*(1.+X41)
Y36=X11*(1.-X42)
Y37=X15*(1.+X43)
YD3=Y33*Y37-Y34*Y36
Y38=(Y32*Y37-Y35*Y34)/YD3
Y39=(Y33*Y35-Y32*Y36)/YD3
Y40=(Y38*X3+Y39*X7+Y28*X1+Y29*X5+Y30*X2+Y31*X6)/R
RA13=Y40/2.
Y41=(ERS14*Y15+ERS24*Y17+ERS34*Y28)/X44
Y42=(ERS14*Y16+ERS24*Y18+ERS34*Y29)/X44
Y43=(ERS24*Y25+ERS34*Y30)/X45

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Y44=-(ERS24*Y26+ERS34*Y31)/X45
Y45=-ERS34*Y38/X46
Y46=-ERS34*Y39/X46
Y47=-Y41*X1*(1.+X47)-Y42*X5*(1.-X48)-Y43*X2*(1.+X49)-Y44*
#X6*(1.-X50)-(Y45*X3*(1.+X51)-Y46*X7*(1.-X52)
Y48=X4*(1.+X55)
Y49=X8*(1.-X56)
Y50=-Y41*X9*(1.-X57)-Y42*X13*(1.+X58)-Y43*X10*(1.-X59)-
#Y44*X14*(1.+X60)-Y45*X11*(1.-X61)-Y46*X15*(1.+X56)
Y51=X12*(1.-X63)
Y52=X16*(1.+X64)
YD4=Y48*Y52-Y49*Y51
Y53=(Y47*Y52-Y49*Y50)/YD4
Y54=(Y48*Y50-Y47*Y51)/YD4
Y55=(Y53*X4+Y54*X8+Y41*X1+Y42*X5+Y43*X2+Y44*X6+Y45*X3+Y46*
#X7)/R
RA14=Y55/2.
Y56=X10*(.25-X28/4.)/H
Y57=X14*(.25+X29/4.)/H
Y58=X2*(.25+X24/4.)/R
Y59=X6*(.25-X25/4.)/R
YD5=Y56*Y59-Y57*Y58
Y60=-Y57/YD5
Y61=Y56/YD5
RA22=Y60*X2*(.25-X24/4.)/R+Y61*X6*(.25+X25/4.)/R
Y62=-ERS23*Y60/X31
Y63=-ERS23*Y61/X31
Y64=-Y62*X2*(1.+X34)-Y63*X6*(1.-X35)
Y65=X3*(1.+X36)
Y66=X7*(1.-X37)
Y67=-Y62*X10*(1.-X40)-Y63*X14*(1.+X41)
Y68=X11*(1.-X42)
Y69=X15*(1.+X43)
YD6=Y65*Y69-Y66*Y68
Y70=(Y64*Y69-Y66*Y67)/YD6
Y71=(Y65*Y67-Y64*Y68)/YD6
Y72=(Y70*X3+Y71*X7+Y62*X2+Y63*X6)/R
RA23=Y72/2.
Y73=-(ERS24*Y60+ERS34*Y62)/X45
Y74=-(ERS24*Y61+ERS34*Y63)/X45
Y75=-ERS34*Y70/X46
Y76=-ERS34*Y71/X46
Y77=-Y73*X2*(1.+X49)-Y74*X6*(1.-X50)-Y75*X3*(1.+X51)-Y76*
#X7*(1.-X52)
Y78=X4*(1.+X55)
Y79=X8*(1.-X56)
Y80=-Y73*X10*(1.-X59)-Y74*X14*(1.+X60)-Y75*X11*(1.-X61)-
#Y76*X15*(1.+X62)
Y81=X12*(1.-X63)
Y82=X16*(1.+X64)
YD7=Y78*Y82-Y79*Y81
Y83=(Y77*Y82-Y79*Y80)/YD7
Y84=(Y78*Y80-Y77*Y81)/YD7
Y85=(Y83*X4+Y84*X8+Y73*X2+Y74*X6+Y75*X3+Y76*X7)/R
RA24=Y85/2.
Y86=X11*(.25-X42/4.)/H
Y87=X15*(.25+X43/4.)/H
Y88=X3*(.25+X36/4.)/R
Y89=X7*(.25-X37/4.)/R
YD8=Y86*Y89-Y87*Y88
Y90=-Y87/YD8
Y91=Y86/YD8
RA33=Y90*X3*(.25-X36/4.)/R+Y91*X7*(.25+X37/4.)/R
Y92=-ERS34*Y90/X46
Y93=-ERS34*Y91/X46
Y94=-Y92*X3*(1.+X51)-Y93*X7*(1.-X52)
Y95=X4*(1.+X55)
Y96=X8*(1.-X56)
Y97=-Y92*X11*(1.-X61)-Y93*X15*(1.+X62)
Y98=X12*(1.-X63)
Y99=X16*(1.+X64)
YD9=Y95*Y99-Y96*Y98
Y100=(Y94*Y99-Y96*Y97)/YD9
Y101=(Y95*Y97-Y94*Y98)/YD9
Y102=(Y100*X4+Y101*X8+Y92*X3+Y93*X7)/R
RA34=Y102/2.
Y103=X12*(.25-X36/4.)/H
Y104=X16*(.25+X64/4.)/H
Y105=X4*(.25+X55/4.)/R

```

```
C
Y106=X8*(.25-X56/4.) /R
YD10=Y103+Y106-Y104= Y105
Y107=-Y104/YD10
Y108=Y103/YD10
RA44=Y107*X4*(.25-X55/4.) /R+Y108*X8*(.25+X56/4.) /R
*****
OUTPUT RESULTS
REFLECTOR ALBEDOS
RA11 = PROBABILITY THAT A NEUTRON FROM GROUP 1 BE
      REFLECTED WITH GROUP 1 ENERGY
RA12 = PROBABILITY THAT A NEUTRON FROM GROUP 1 BE
      REFLECTED WITH GROUP 2 ENERGY , AND SO ON.
*****
WRITE(02,100)R,T
WRITE(02,101)RA11,RA12,RA13,RA14
WRITE(02,102)RA22,RA23,RA24
WRITE(02,103)RA33,RA34,RA44
100 FORMAT(3X,'R'=' ,E12.5,1X,'T'=' ,E12.5)
101 FORMAT(3X,'RA11=',E12.5,1X,'RA12=',E12.5,1X,'RA13=',E12.5,
    #1X,'RA14=',E12.5)
102 FORMAT(3X,'RA22=',E12.5,1X,'RA23=',E12.5,1X,'RA24=',E12.5)
103 FORMAT(3X,'RA33=',E12.5,1X,'RA34=',E12.5,1X,'RA44=',E12.5)
STOP
END
```

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BIOGRAPHICAL SKETCH

Ronaldo Glicerio Cabral was born on July 8, 1946, in Rio de Janeiro, Brazil. In December of 1968, Ronaldo Glicerio Cabral graduated as an army officer at the Academia Militar de Agulhas Negras (Agulhas Negras Military Academy) in Rio de Janeiro, Brazil. In December of 1976, he received the degree of Bachelor of Science in electrical engineering from the Instituto Militar de Engenharia (Military Engineering Institute), in Rio de Janeiro, Brazil. He received the degree of Master of Science in nuclear engineering in December of 1981, from the same Institute.

Ronaldo Glicerio Cabral worked on many projects for the Brazilian Army and taught courses in the Nuclear Engineering Department in the Military Engineering Institute.

In January of 1989, he was admitted to the graduate program of the Nuclear Engineering Department at the University of Florida to pursue a Ph.D. degree, which has been sponsored by the Brazilian Army.

Ronaldo Glicerio Cabral has been a member of Alpha Nu Sigma, the Nuclear Engineering Honor Society, since 1990.

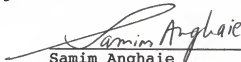
He is married to Aureliana Azambuja Costa Cabral. Their son, Ronaldo Glicerio Cabral Filho, is 16 years old.

I certify that I have read this study and that in my opinion it conforms to acceptable standards of scholarly presentation and is fully adequate, in scope and quality, as a dissertation for the degree of Doctor of Philosophy.



Alan M. Jacobs, Chairman
Professor of Nuclear
Engineering Sciences

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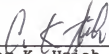
Samim Anghaie
Associate Professor of Nuclear
Engineering Sciences

I certify that I have read this study and that in my opinion it conforms to acceptable standards of scholarly presentation and is fully adequate, in scope and quality, as a dissertation for the degree of Doctor of Philosophy.




Edward T. Dugan
Associate Professor of Nuclear
Engineering Sciences

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Chung K. Hsieh
Professor of Mechanical
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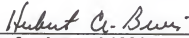
I certify that I have read this study and that in my opinion it conforms to acceptable standards of scholarly presentation and is fully adequate, in scope and quality, as a dissertation for the degree of Doctor of Philosophy.



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This dissertation was submitted to the Graduate Faculty of the College of Engineering and to the Graduate School and was accepted as partial fulfillment of the requirements for the degree of Doctor of Philosophy.

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